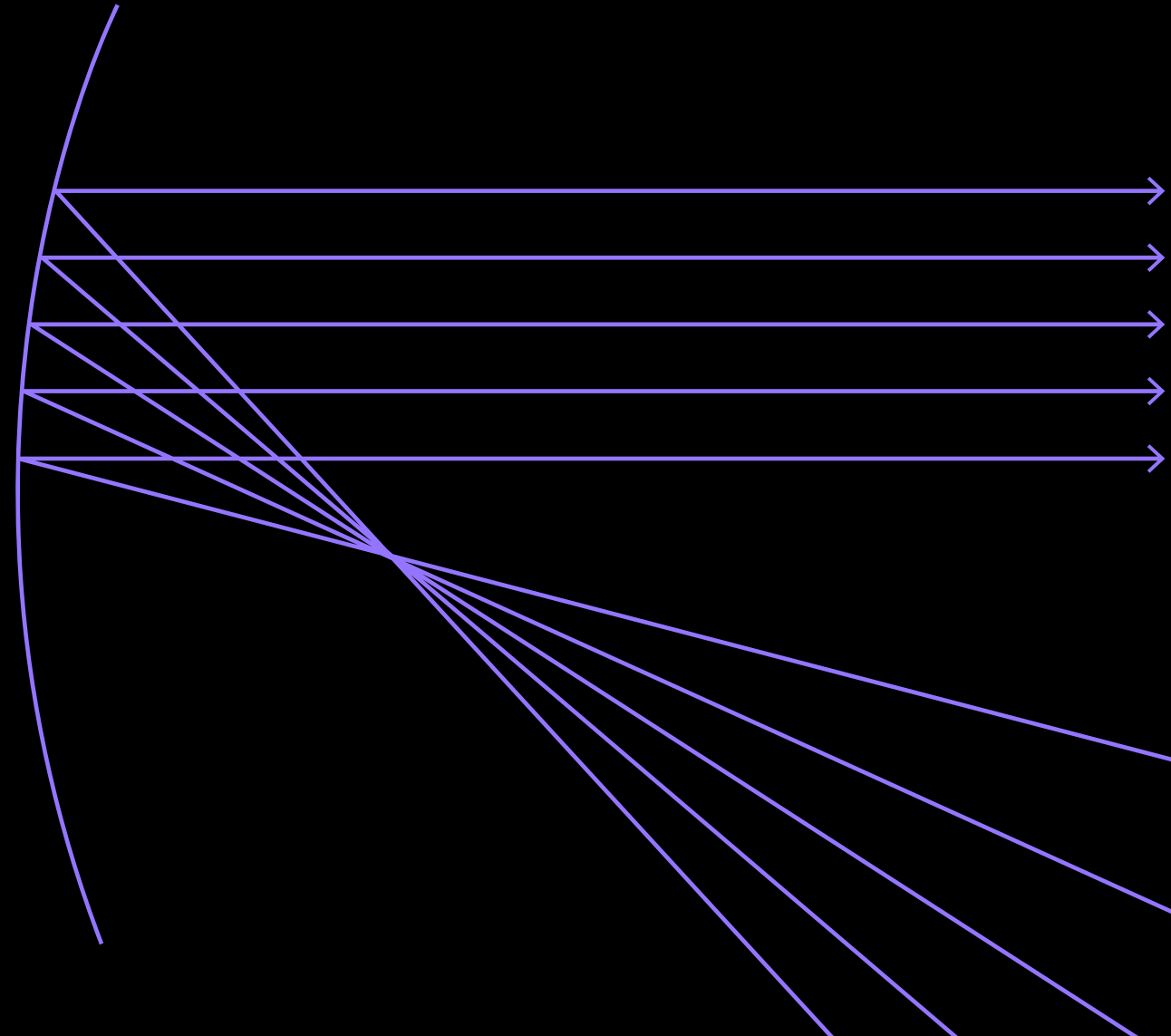


All in a Spin: Rotating Trapped Microspheres

Featuring Kishan Dholakia University of Adelaide, University of St. Andrews

09 February 2023



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A Quick Zoom Tutorial

- Submit a question by clicking on “Q&A”
- Like a question that’s been submitted?
Click the “thumbs up” icon to vote for it.
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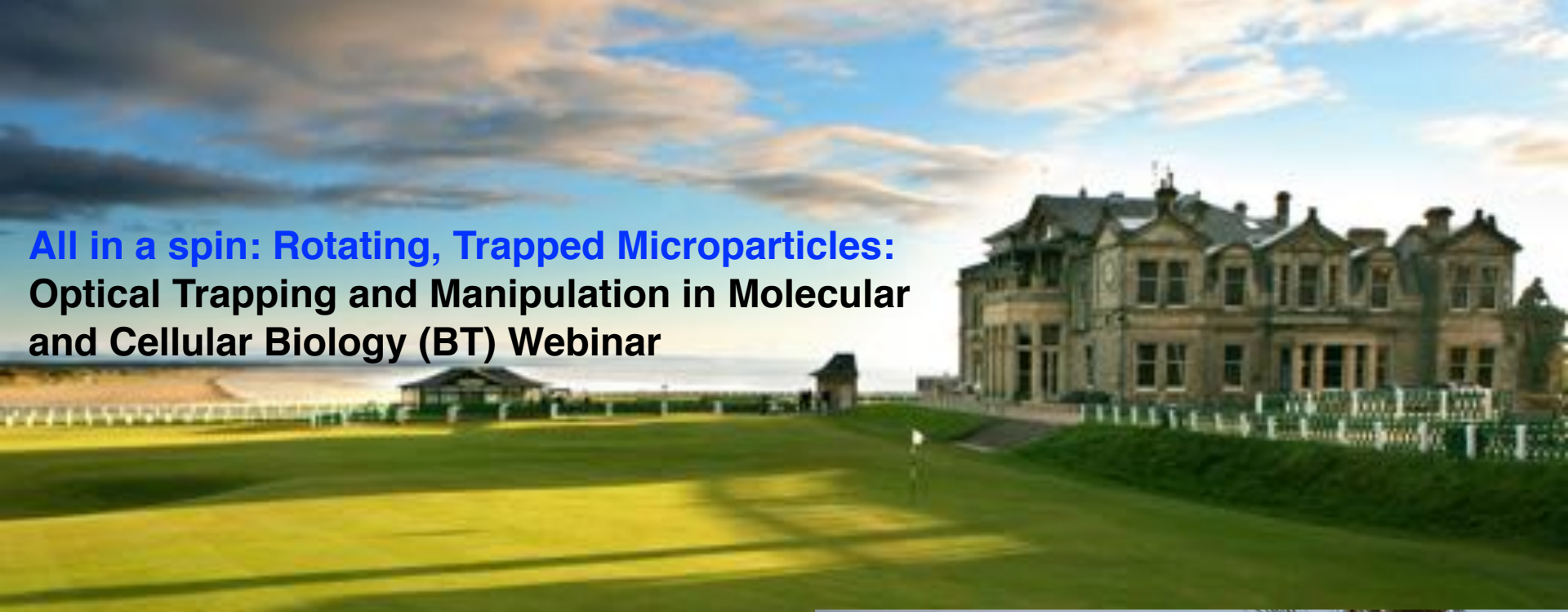
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**All in a spin: Rotating, Trapped Microparticles:
Optical Trapping and Manipulation in Molecular
and Cellular Biology (BT) Webinar**

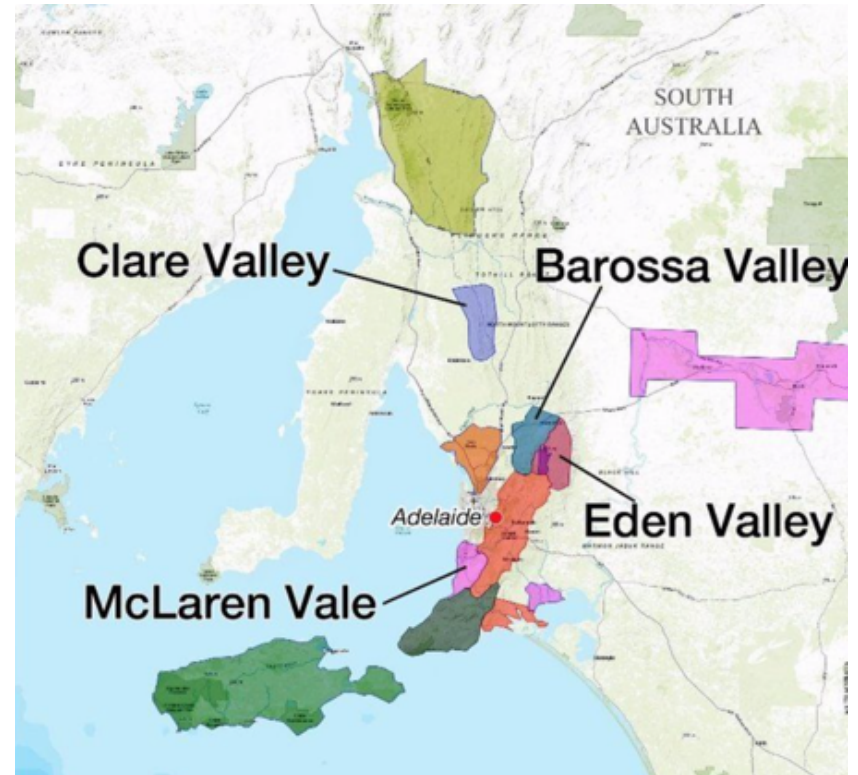
Kishan Dholakia
***Centre of Light for Life, University of
Adelaide, Australia and
School of Physics and Astronomy
University of St Andrews, Scotland***

 @OpticManip

<http://opticalmanipulationgroup.wp.st-andrews.ac.uk>

OPTICA | Formerly
OSA

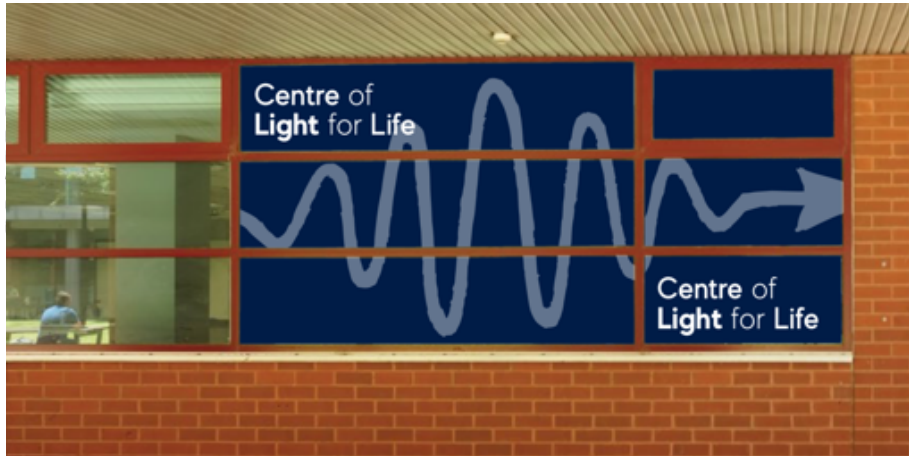






Secured over A\$8M in funding since 2021

500m² lab and office space for 20 people



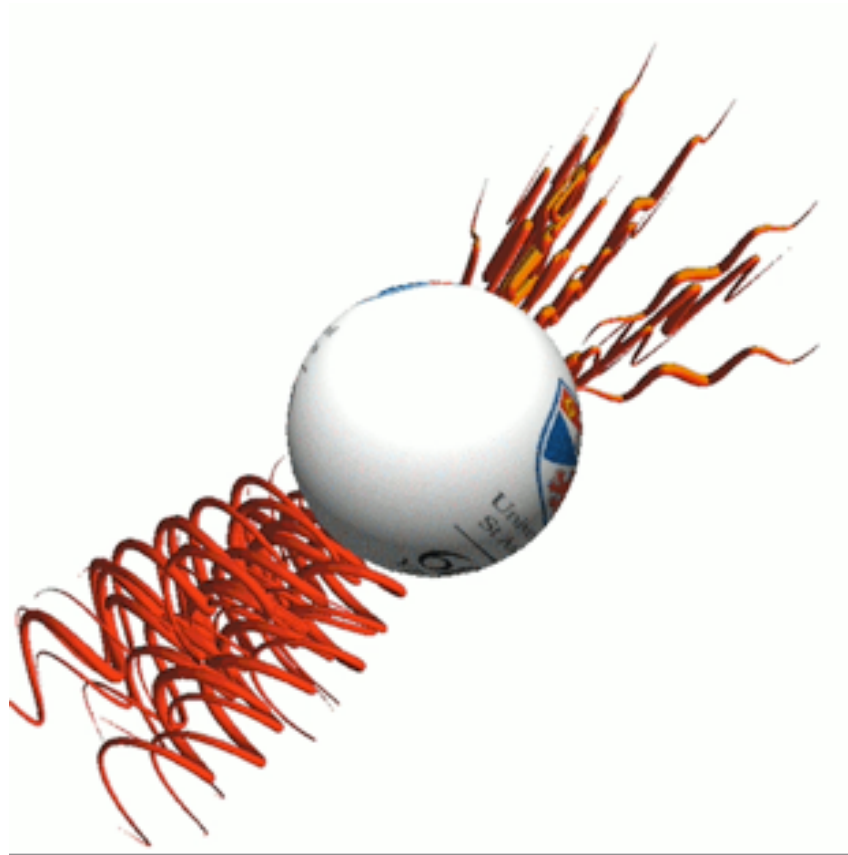
TOPICS

**IMAGING AT DEPTH
MANIPULATION AND LEVITATED OPTOMECHANICS
QUANTUM INSPIRED BIOPHOTONICS
HEALTH: IVF IMPROVEMENT, EARLY DIAGNOSIS**

Group of 18 (Jan 2023).

Several PDRA and PhD positions (incl joint with St Andrews) - please get in touch!

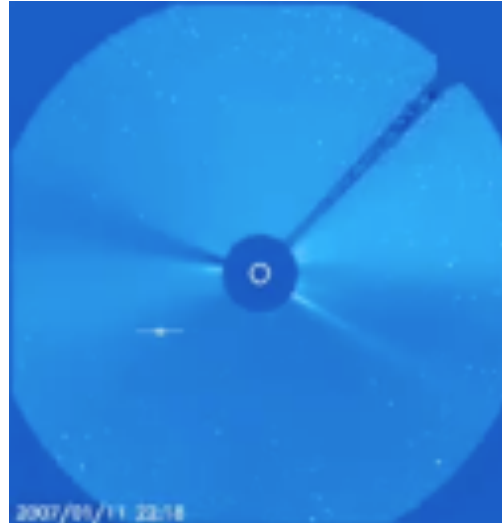
Rotation in optical traps



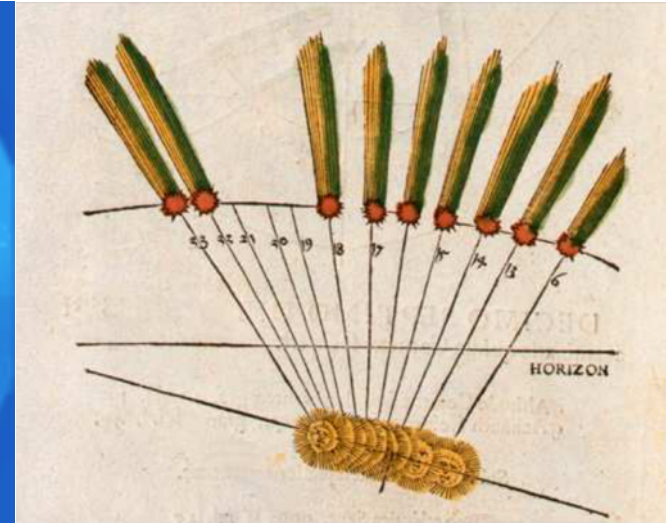
The Tail of a Comet



Pen & Ink Rendition of Engraved Frontispiece from "The Tome of Hevelius," 1668.



<http://sohowww.nascom.nasa.gov/hotshots/>



Peter Arpian « Astronomicum Caesareum » (1577)

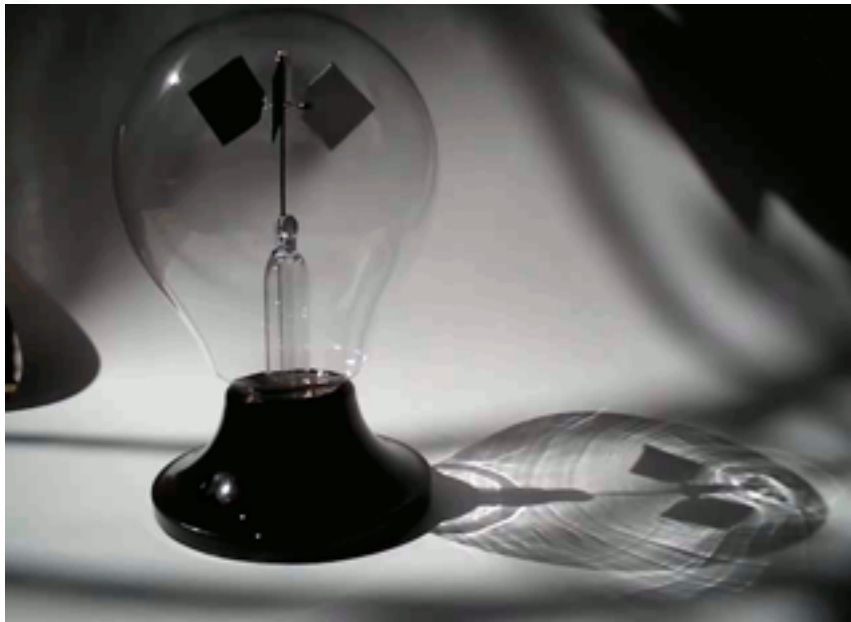
Kepler (1619): 'The direct rays of the Sun strike upon it [the comet], penetrate its substance, draw away with them a portion of this matter, and issue thence to form the track of light we call the tail . . . In this manner the comet is consumed by breathing out its own tail.'

Pen & Ink Rendition of Engraved Frontispiece from "The Tome of Hevelius," 1668.

The search for the correct form: On the title page of Cometographia by Johannes Hevelius, Aristotle (left), Johannes Kepler (right) and the author discuss the trajectories of comets [Credit: © Max Planck Society]

In 1873 Maxwell,¹ on the basis of the electromagnetic theory, showed that if light were an electromagnetic phenomenon, pressure should result from the absorption or reflection of a beam of light. After a discussion of the equations involved, he says: "Hence in a medium in which waves are propagated there is a pressure in the direction normal to the waves and numerically equal to the energy in unit volume." Maxwell computed the pressure exerted by the Sun on the illuminated surface of the Earth, and added:

It is probable that a much greater energy of radiation might be obtained by means of the concentrated rays from an electric lamp. Such rays falling on a thin metallic disk, delicately suspended in a vacuum, might perhaps produce an observable mechanical effect.



James Clerk Maxwell celebrated at final resting place in Parton

© 10 March 2020

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THE PRESSURE DUE TO RADIATION.¹

By E. F. NICHOLS and G. F. HULL.

ASTROPHYSICAL JOURNAL

THE NOBEL PRIZE IN PHYSICS 2018

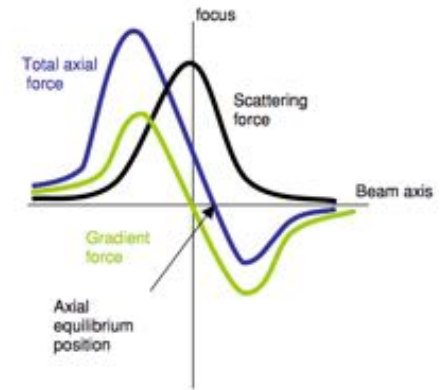
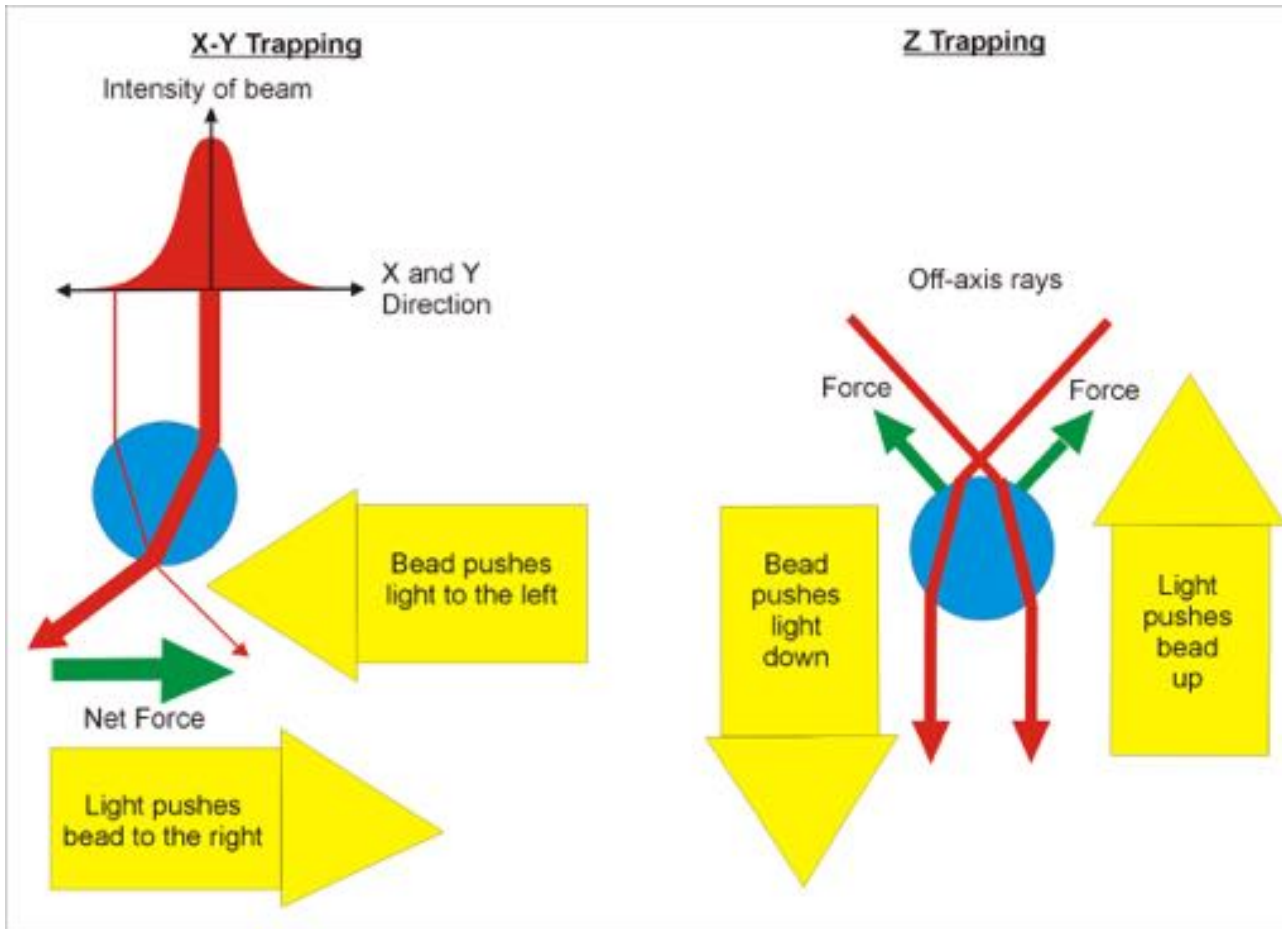


▲ Arthur Ashkin, Gérard Mourou and Donna Strickland, 2018's Nobel laureates in physics. Photograph: Nobel Assembly.

Optical Tweezers

SINGLE BEAM GRADIENT TRAP: Ashkin et al, Opt Lett 11, 288 (1986)

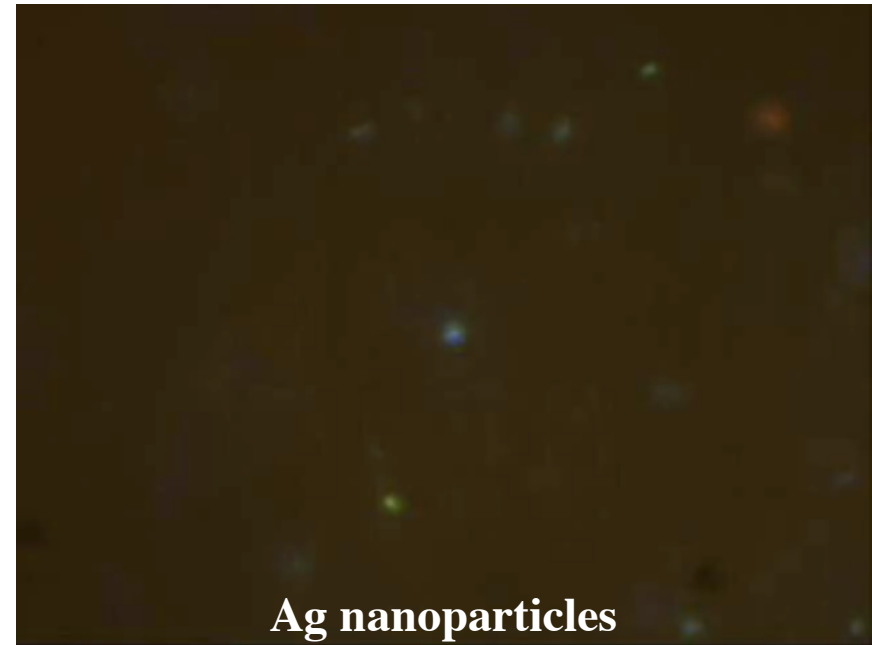
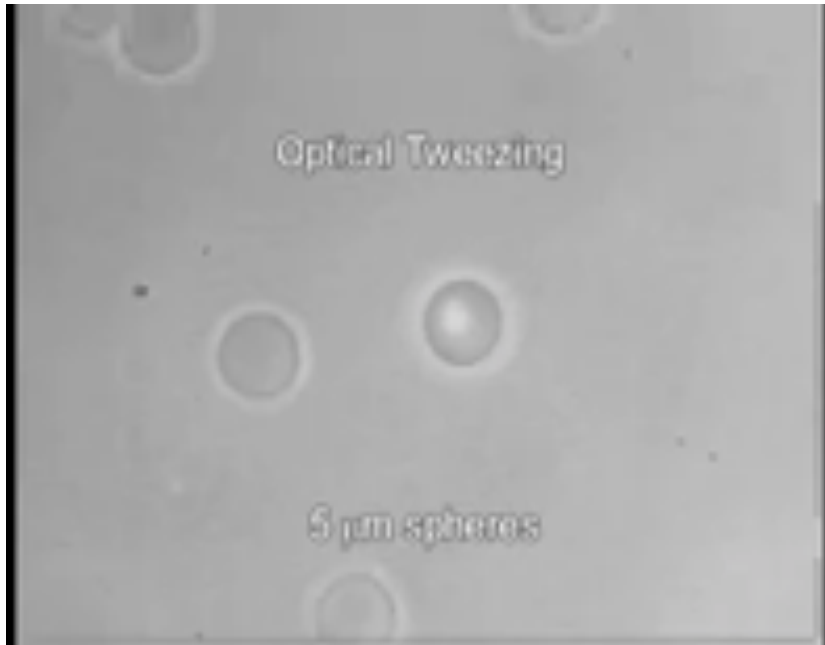
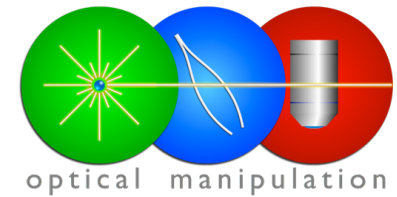
Co-recipient of the Nobel Prize 2018



$$F_{\text{gradient}} \propto \alpha \nabla I$$

$$F_{\text{scattering}} \propto I$$

Microscale to the Nanoscale

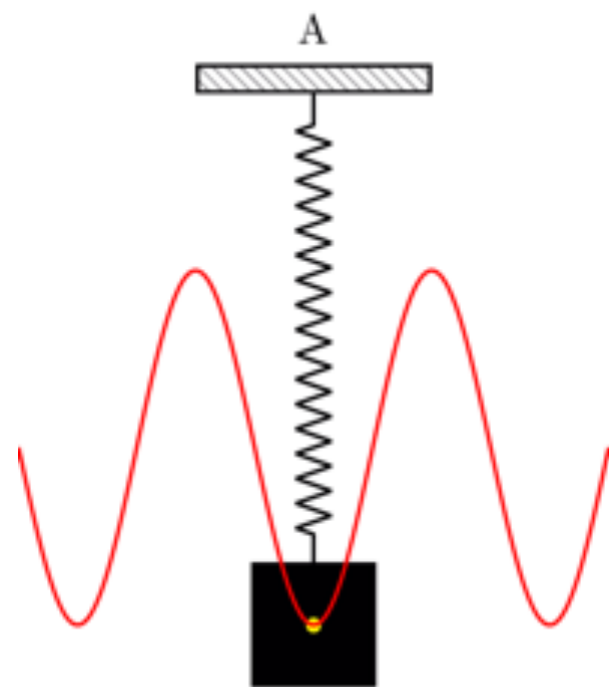


Care must be taken with heating! (trap light used from 700nm-1100nm)
Particle polarisability is key, as are field gradients

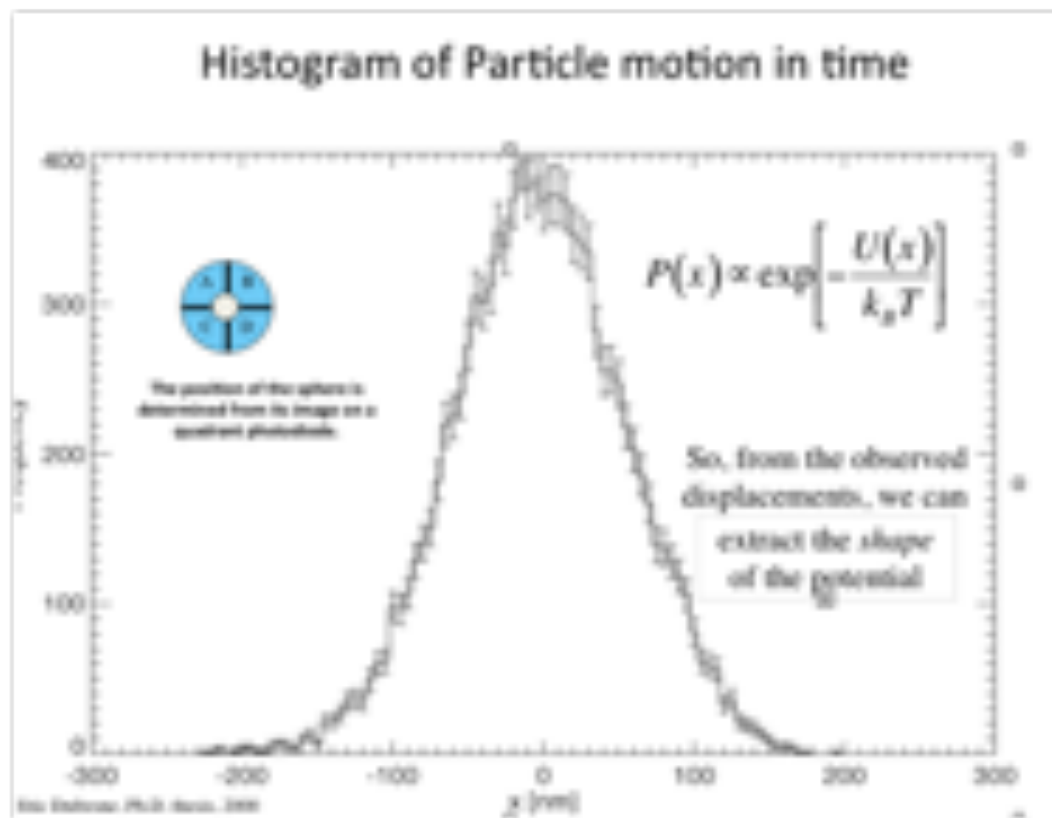
$$F_{\text{grad}} \propto \alpha \cdot \nabla I(r)$$

Reason why we go to plasmonics, nanoapertures, waveguides etc

Optical tweezers: the World's most elegant example of a Hookean spring?



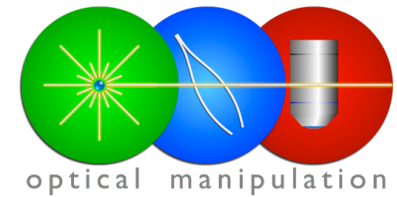
$$F_{\text{trap}} = -kx$$





Optical tweezers has revolutionised single molecule motion: e.g. kinesin steps on microtubules ...8nm..

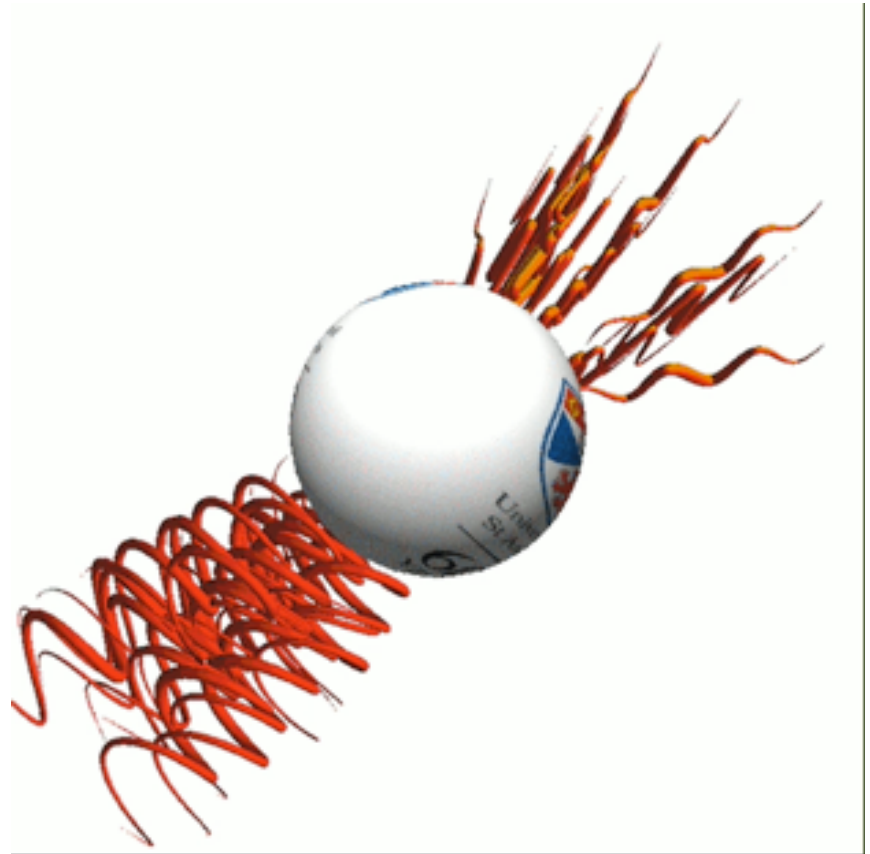
How do we cause particles to rotate in optical traps?



Two main methods:

**Asymmetric
Rotating light
patterns that cause
objects to align (e.g.
form birefringence)**

**Transfer of optical
angular momentum:
spin and orbital**



All-optical control of microfluidic components using form birefringence

STEVEN L. NEALE*, MICHAEL P. MACDONALD, KISHAN DHOLAKIA AND THOMAS F. KRAUSS

School of Physics and Astronomy, St Andrews University, St Andrews, Fife, KY16 9SS, UK

*e-mail: sln2@st-and.ac.uk

Nature Materials **4**, 530–533 (2005)

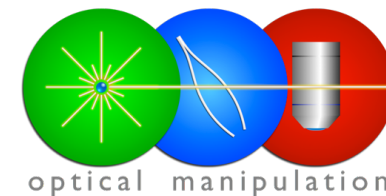
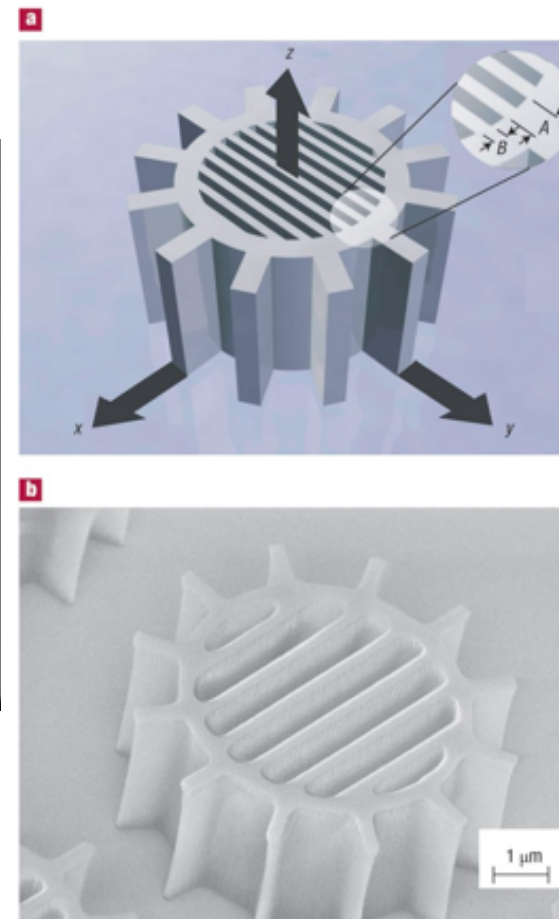


Figure 1: Form birefringent microgears.



a, Diagram of a form birefringent microgear. A is the pitch of the photonic lattice, B is the rib width and the fill factor is B/A . **b**, SEM micrograph of an actual microgear before release from the substrate.

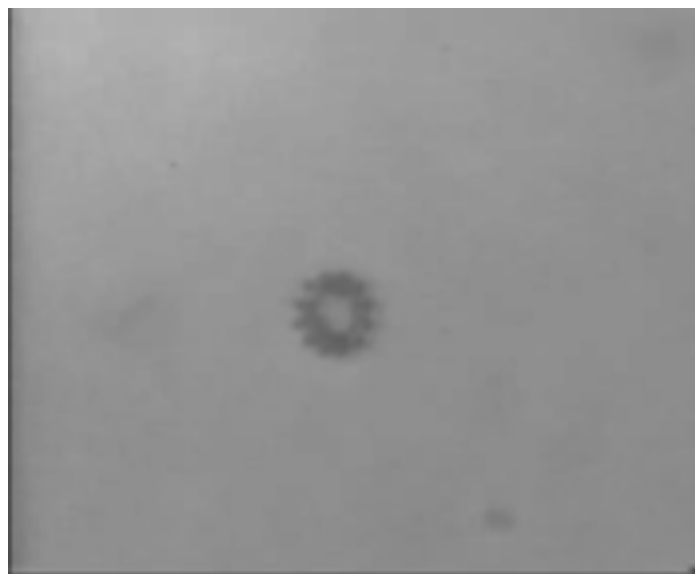
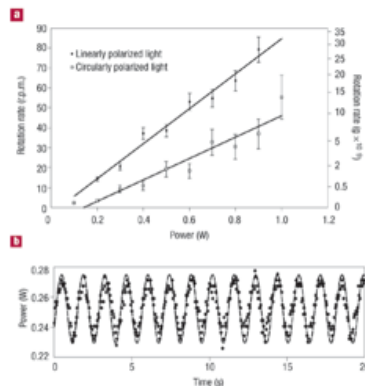


Figure 4: Rotation results.



a, Maximum rotation rates with increasing power for both linearly and circularly polarized light. **b**, Light transmitted through the microgear that has been analysed by passing it through a polarizer. The amplitude of the power variation indicates the magnitude of the birefringence, and the rotation rate can be gained from the period. The error bars represent the deviation from the mean rotation rate when averaging over 12 revolutions of the microgear.

Creating multiple traps: time sharing

Acousto-Optic Deflectors (AODs) can be scanned at *hundreds* of kHz



(video in collaboration with I Poberaj group).



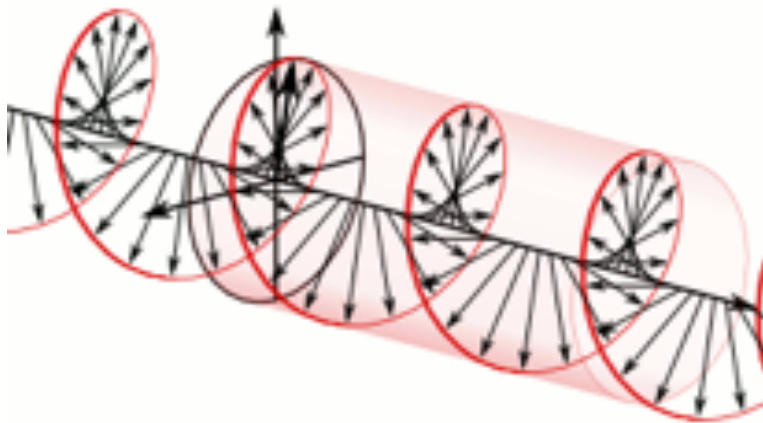
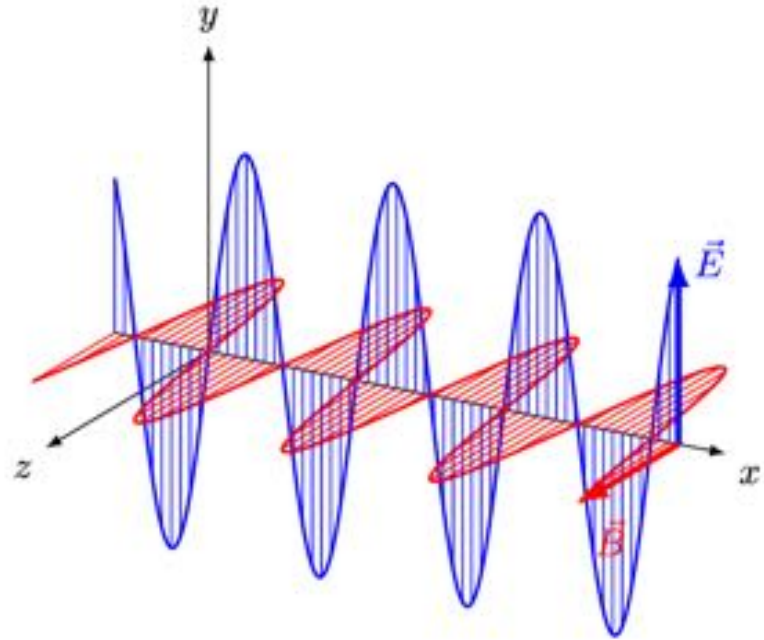
A microfluidic pump made from glass beads the size of a heart valve. (DM Marr et al. Science 2002)

Table 1. A quantitative description of parameters for time-shared devices.

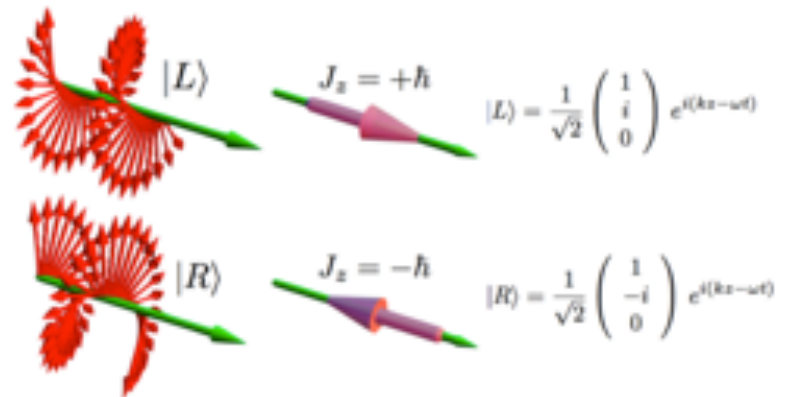
Device	Switching rate (kHz)	Power efficiency (%)	Angular range of deflection (mrad)	Resolution (μrad)	Reference
Galvo-mirror	1	— ^a	500	10	[9, 62]
Piezo-mirror	1	— ^a	50	0.1	[63]
AOD	10–50	50	30	<1 nm ^b	[64]
EOD	>100	80	1	<1 nm ^b	[69]

Light is a transverse electromagnetic wave

The E-field can oscillate up and down (linear polarization) or rotate (circular polarization – spin angular momentum)



$\pm \hbar$ per photon



Structured light can possess angular momentum: rotation

$$\mathbf{j} = \epsilon_0 \left[\mathbf{r} \times \langle \mathbf{E} \times \mathbf{B} \rangle \right]$$

Allen et al Phys Rev A (1992)

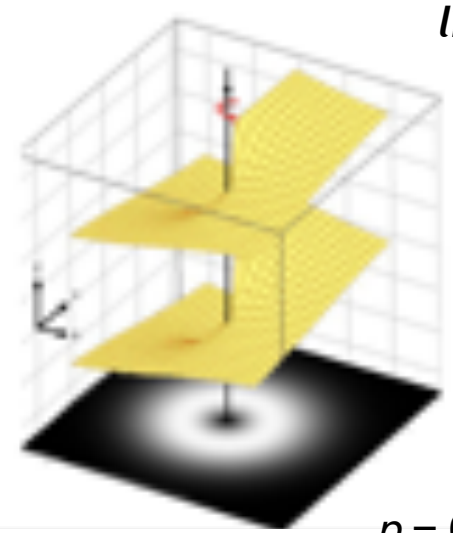
Spin: due to polarisation state (rotating E-field)

Orbital: due to inclined wavefronts

$\pm \hbar$ per photon

$|L\rangle \quad J_z = +\hbar \quad |L\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \\ 0 \end{pmatrix} e^{i(kz - \omega t)}$

$|R\rangle \quad J_z = -\hbar \quad |R\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \\ 0 \end{pmatrix} e^{i(kz - \omega t)}$



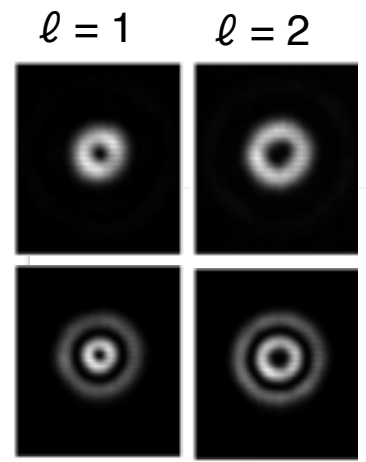
$l\hbar$ per photon

Laguerre-Gaussian (LG) beam

$$\mathbf{S} = \epsilon_0 \int (\mathbf{E} \times \mathbf{A}) d^3 \mathbf{r}$$

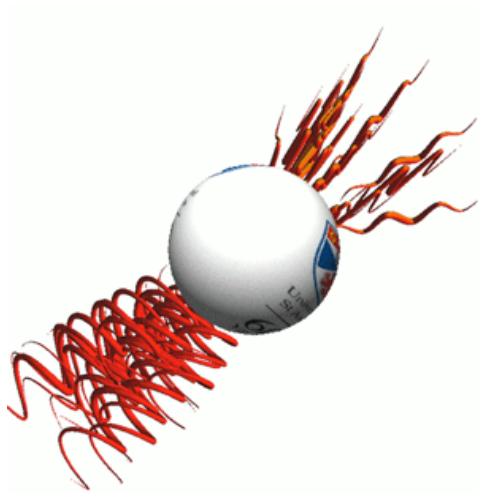
→
$$\mathbf{S} = \frac{\epsilon_0}{2i\omega} \int (\mathbf{E}^* \times \mathbf{E}) d^3 \mathbf{r}$$

$$\mathbf{S}_z = \pm \hbar$$

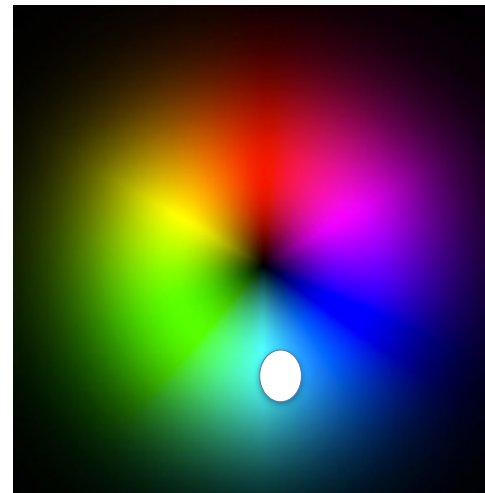


LG modes

Transfer of angular momentum to trapped particles



SPIN



ORBITAL

Structured light can possess angular momentum: rotation

$$\mathbf{j} = \epsilon_0 \left[\mathbf{r} \times \langle \mathbf{E} \times \mathbf{B} \rangle \right]$$

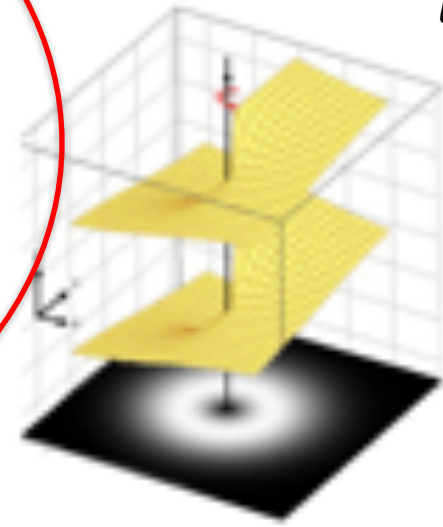
Allen et al Phys Rev A (1992)

Spin: due to polarisation state (rotating E-field)

Orbital: due to inclined wavefronts

$\pm \hbar$ per photon

$l\hbar$ per photon



Laguerre-Gaussian (LG) beam

$$\mathbf{S} = \epsilon_0 \int (\mathbf{E} \times \mathbf{A}) d^3 \mathbf{r}.$$

➔
$$\mathbf{S} = \frac{\epsilon_0}{2i\omega} \int (\mathbf{E}^* \times \mathbf{E}) d^3 \mathbf{r}.$$

$$\mathbf{S} = \sum_{\mathbf{k}} \hbar \mathbf{u}_{\mathbf{k}} \left(\hat{a}_{\mathbf{k},L}^\dagger \hat{a}_{\mathbf{k},L} - \hat{a}_{\mathbf{k},R}^\dagger \hat{a}_{\mathbf{k},R} \right)$$

$$S_z = \pm \hbar.$$

Demonstration of spin angular momentum: R. Beth, 1936, Princeton

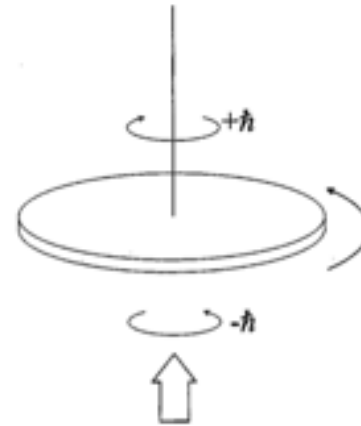
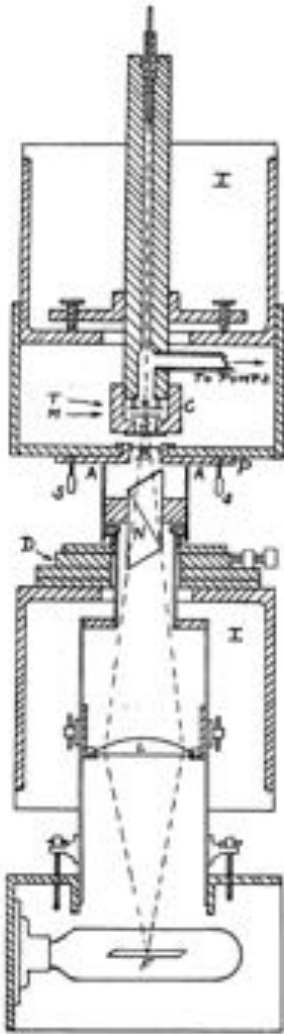


Fig. 1. Beth's experiment from 1936 (Ref. 2) was based on measuring the torque as circularly polarized light passed through a suspended quarter-wave plate.

Experimental proof of that theoretical prediction was done by R. Beth in 1936 in Princeton. As Beth announces in his paper (R. A. Beth, *Mechanical Detection and Measurement of the Angular Momentum of Light*, Physical Review, v. 50, July 15, 1936) he had several discussions about the experiment with Einstein.

In this experiment Beth showed that when linearly polarized light is converted to circularly polarized one by doubly refracting slab, the slab experiences a reaction torque.

Experimental setup of Beth's experiment

Rotation using birefringence

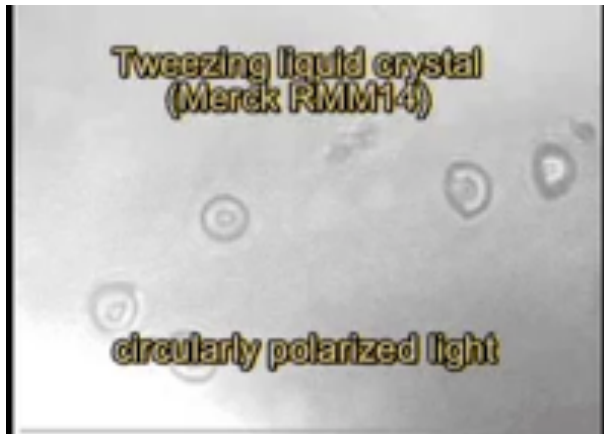
Right handed circularly polarised light



Left handed circularly polarised light



Trapped birefringent microsphere ('golf ball')

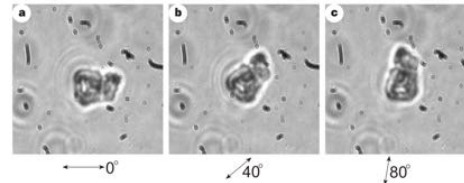


Optical alignment and spinning of laser-trapped microscopic particles

M. E. J. Friese, T. A. Nieminen, N. R. Heckenberg & H. Rubinsztein-Dunlop

Centre for Laser Science, Department of Physics, The University of Queensland, Brisbane, Queensland 4072, Australia

Friese et al
Nature 1998



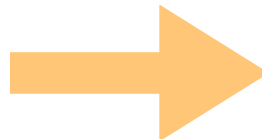
$2\hbar$ Per photon transferred

$$\phi = \frac{2\pi}{\lambda} (n_e - n_o) d$$

$$\tau_{\text{Optical}} = \frac{\Delta L}{\Delta t} = (\eta \Delta \sigma) \left(\frac{P}{h\omega} \right) \dot{t}$$

and for a spherical particle:

$$\tau_{\text{drag}} = 8\pi\eta\Omega R^3$$



$$\eta_{\text{local}} = \frac{\Delta\sigma P}{8\pi R^3 \Omega \omega}$$

Nanofabricated quartz cylinders for angular trapping: DNA supercoiling torque detection

Christopher Deufel¹, Scott Forth¹,
Chad R Simmons^{1,2}, Siavash Dejgoshia¹ &
Michelle D Wang¹

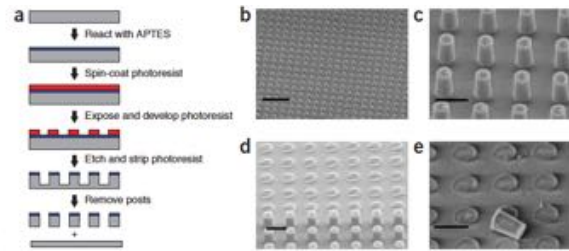
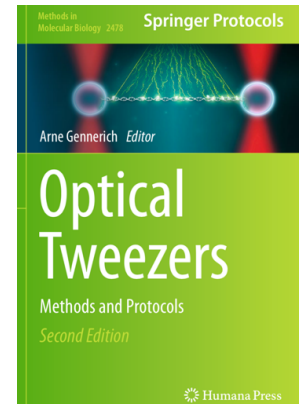
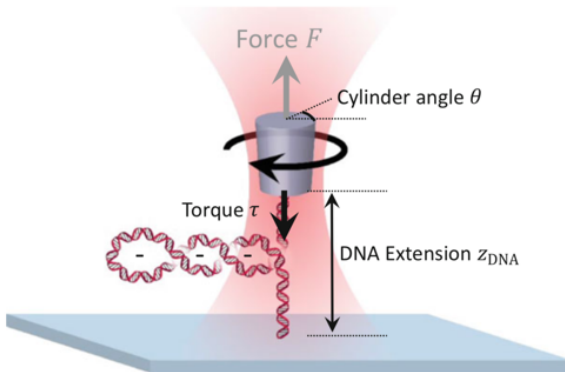


Figure 2 | Nanofabrication of quartz cylinders. (a) Schematic outline of the nanofabrication protocol. (b–e) Scanning electron micrographs of nanofabricated cylinders. Nanofabricated cylindrical posts on the wafer (b,c). The cylinders were 1.1 μm high and 0.53 μm in diameter. Quartz substrate after a portion of the posts was removed from the wafer (d). The quartz posts fractured evenly at their bases in a consistent manner. A single quartz cylinder after mechanical removal (e). Scale bars, 5 μm in b and 1 μm in c–e.



NATURE METHODS | VOL.4 NO.3 | MARCH 2007



As twist was introduced to a DNA molecule, torque increased essentially linearly until the DNA buckled to form a plectoneme, after which the torque plateaued. Vertical dashed lines: buckling transitions

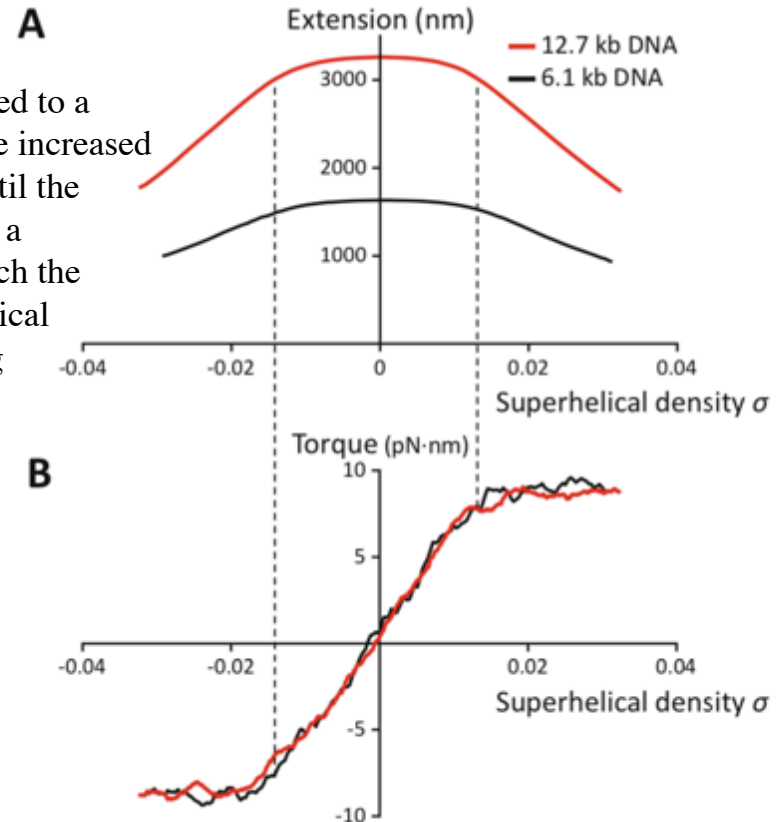
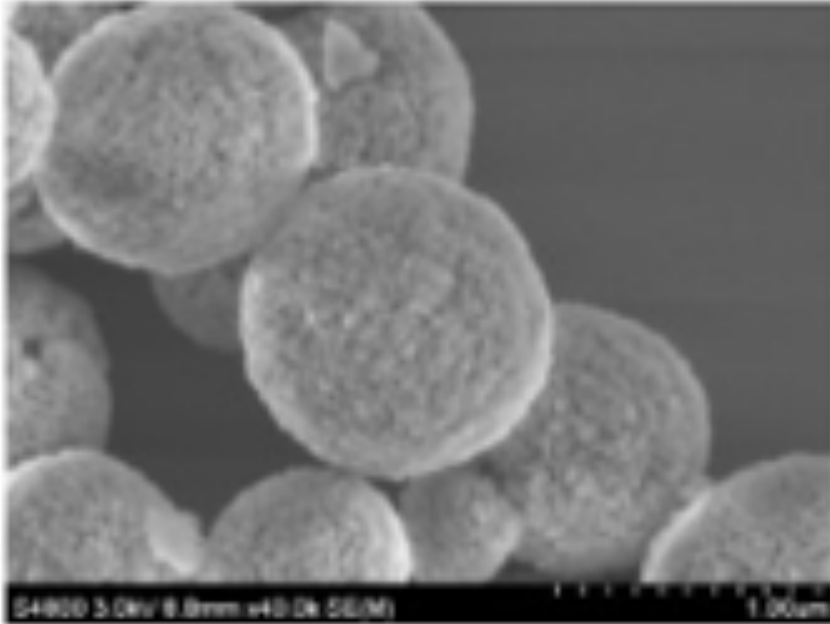


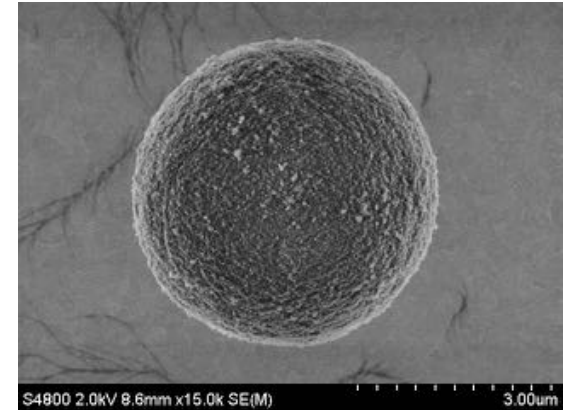
Fig. 14 Measured torque and extension of DNA molecules with different lengths. Extension (a) and torque (b) measurements for naked DNA molecules of 12.7 kbp (red) and 6.1 kbp (black) in length are shown. The

1. Le TT, et al., . Cell 179(3):619–631.e615 (2019)

Birefringent spheres: synthesis of vaterite and nanovaterite



The nucleation and growth rate of the vaterite spheres is determined by the supersaturation level of the dissolved amorphous CaCO₃.



4.4 μm vaterite (see later)

A key issue is avoidance of rapid recrystallisation into the calcite phase, which can occur due to the enhanced solubility of the particles as the particle size decreases.

To avoid this, ethylene glycol was added to the water used as the solvent for this reaction.

SEM for
nanovaterite at St
Andrews

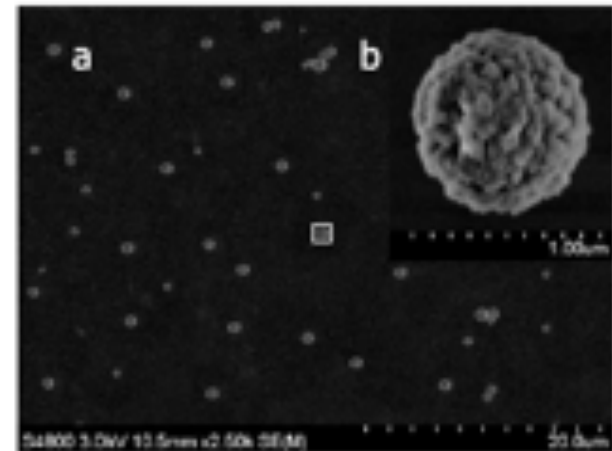


Fig. 2. SEM image of nanovaterite particles produced via the co-precipitation method in 83% ethylene glycol solvent. Mean particle size was found to be 846 nm.

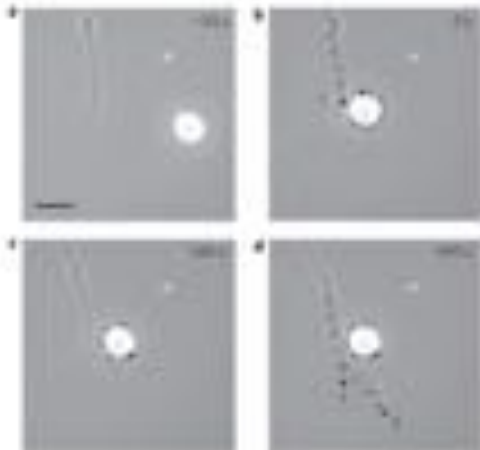
Optical forces and torques can direct the growth of neurons

A photon-driven micromotor can direct nerve fibre growth

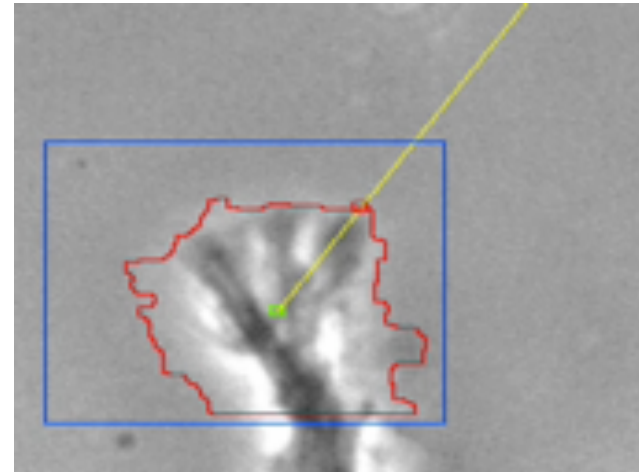
Tao Wu, Timo A. Nieminen, Samarendra Mohanty, Jill Miotke, Ronald L. Meyer, Halina Rubinsztein-Dunlop & Michael W. Biers

Nature Photonics **6**, 62–67 (2012) | doi:10.1038/nphoton.2011.287

Figure 1: Time-lapse images when a vaterite particle is rotated anticlockwise and positioned to the left of the axon defined by the growth direction of the axon (dashed arrow 1).



- a. Before the trapped vaterite particle was moved near the axon.
- b. The vaterite particle was moved to the left of the axon and rotated anticlockwise at -1 Hz.
- c. After 340 s, the axonal growth cone had already turned to a new direction...



Carnegie, D. J., Stevenson, D. J., Mazilu, M., Gunn-Moore, F. & Dholakia, K. Guided neuronal growth using optical line traps. *Opt. Express* **16**, 10507–10517 (2008).

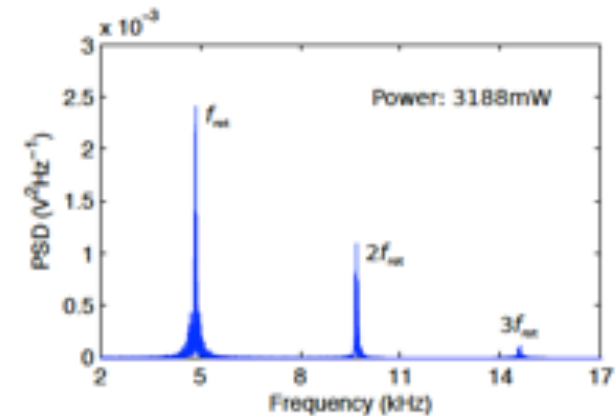
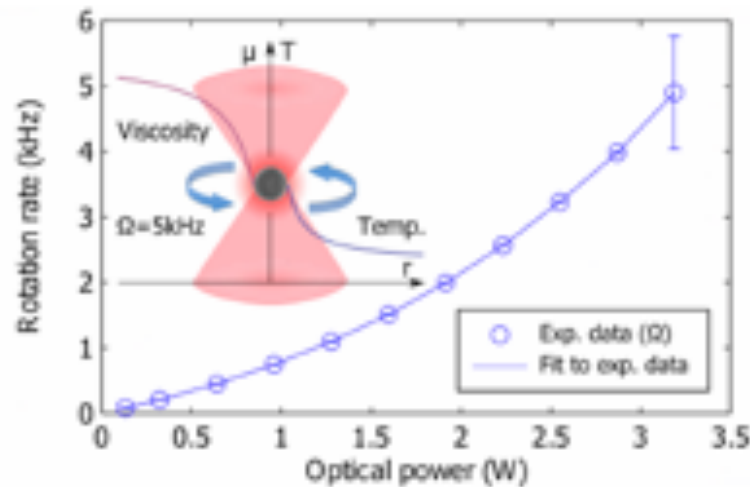
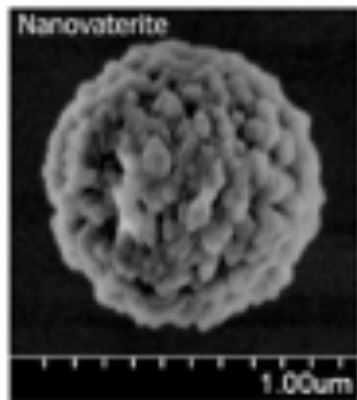
original paper:

Ehrlicher, A. *et al.* Guiding neuronal growth with light. *Proc. Natl Acad. Sci. USA* **99**, 16024–16028 (2002).

Optical trapping and rotation of nanovaterite particles

Nanovaterite trapped by a 532nm circularly polarised beam in D₂O.

→ No heating to the medium $< 0.1\text{ }^{\circ}\text{C}\text{W}^{-1}$



See also Schmidt group work, e.g. 8mK/mW (1064nm) - dominated by fluid
Biophysical Journal 84, 1308 (2002) using normal water

Previous vaterite study: Parkin, S. J. et al., Phys. Rev. E 2007, 76, 041507: For vaterite crystals,
a temperature increase of 66 CW^{-1} was inferred @ 1064nm

Rings, D.; Chakraborty, D.; Kroy, K. *New J. Phys.* **2012**, *14*, 053012.

Our study isolates thermal effects in the microsphere

Determining properties of nanovaterite

A laminar Navier-Stokes model is used which can then deliver the overall drag torque or force for the rotational and translational motion of the nanoparticle.

Finite element method (implemented in COMSOL) to calculate the overall drag torque or drag force for different residual optical absorption powers.

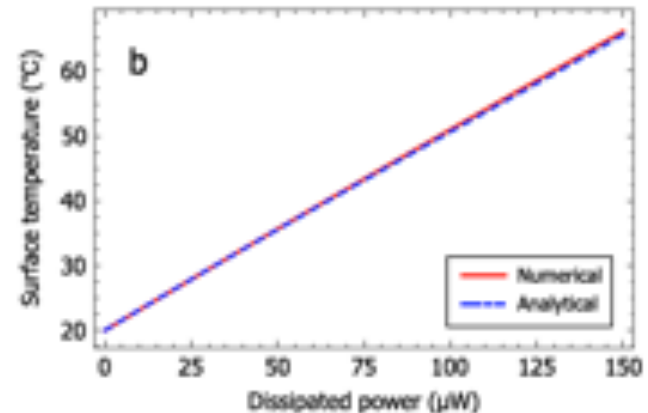
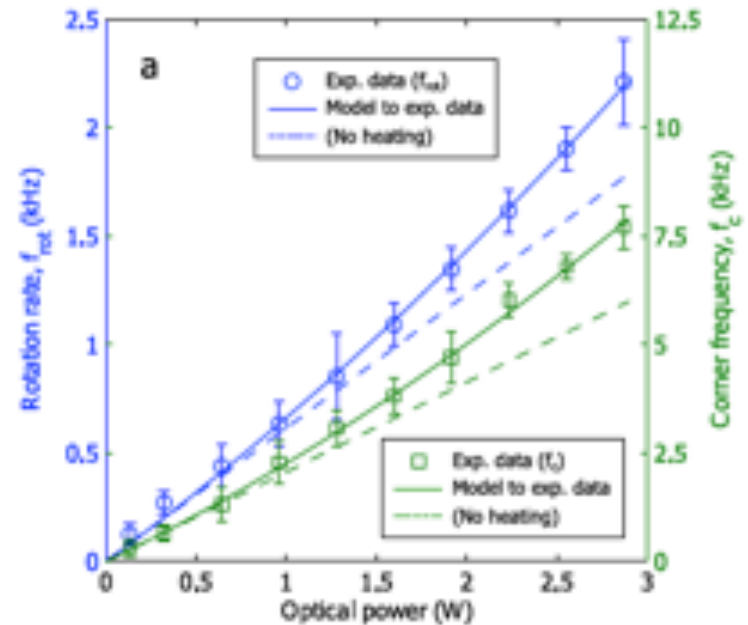
$$\tau = x/v = 6\pi\mu r/\kappa.$$

$$f_c = \tau^{-1} = \kappa/6\pi\mu r \propto P.$$

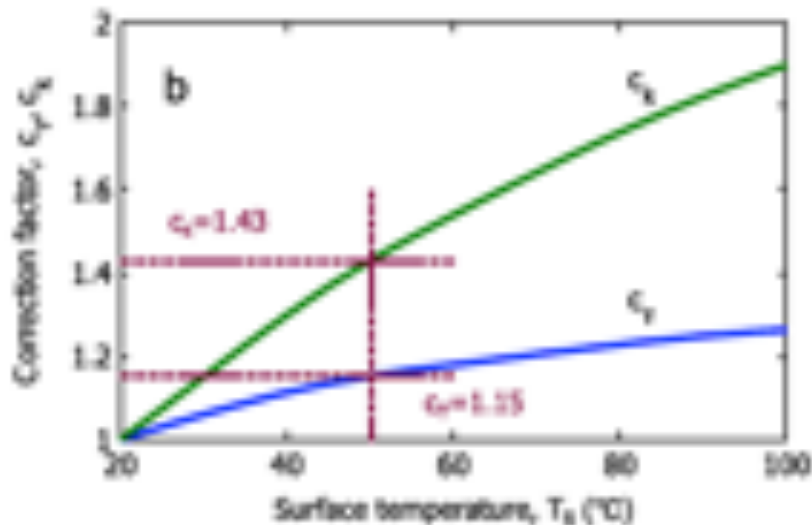
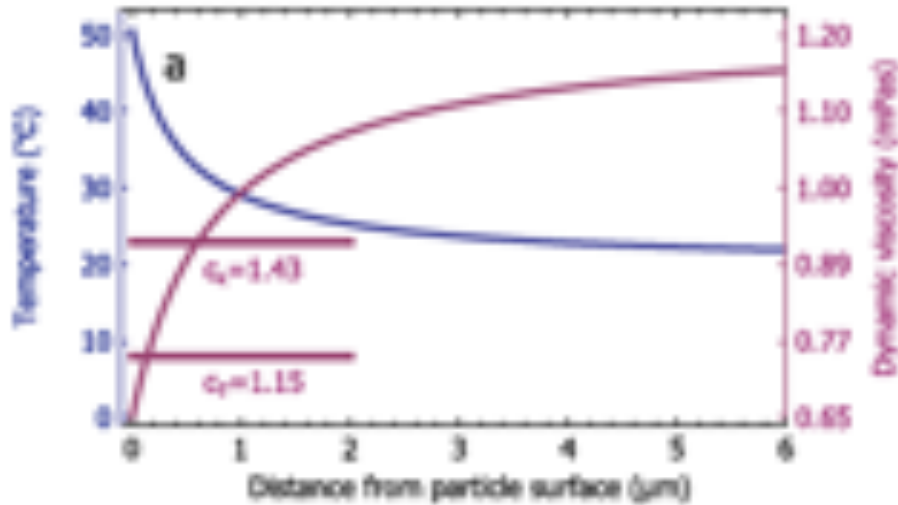
$$T_s = \frac{Q_{\text{abs}}}{4\pi r k(T_s)},$$

absorption coefficient of
 $1:59 \times 10^{-5}$ for
 nanovaterite

i.e. for an incident beam with a power of 1W, the nanoparticle dissipates 15.9 μ W, which leads to the surface temperature of 25:6 C.



Optical trapping and rotation of nanovaterite particles



Temperature dependent dynamic viscosity correction factor.

(a) Radial temperature profile and the corresponding dynamic viscosity of heavy water surrounding a nanoparticle dissipating 100 μW.

(b) Viscosity correction factors for rotational (blue solid line) and translational (green solid line) motion of the nanoparticle (radius of 423nm) as a function of its surface temperature.

$$\tau_{drag} = c_r(T_s) 8\pi\mu(T_s)r^3\Omega,$$

$$F_{drag} = c_t(T_s) 6\pi\mu(T_s)rv,$$

Arita et al., ACS Nano 10, 11505 (2016)

rotational/translational motion differs

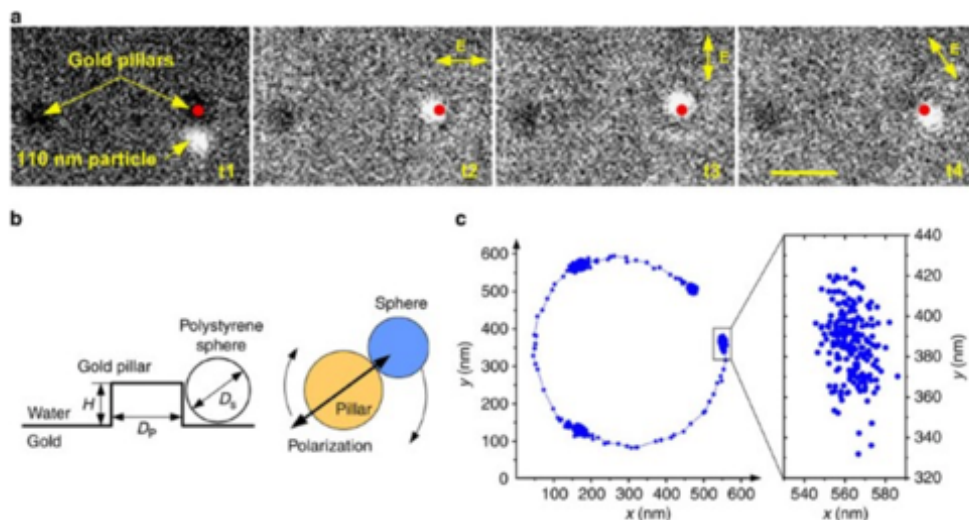
Rings, D.; Chakraborty, D.; Kroy, K. *New J. Phys.* **2012**, *14*, 053012.

Trapping and rotating nanoparticles using a plasmonic nano-tweezer with an integrated heat sink

Kai Wang, Ethan Schonbrun, Paul Steinvurzel & Kenneth B. Crozier

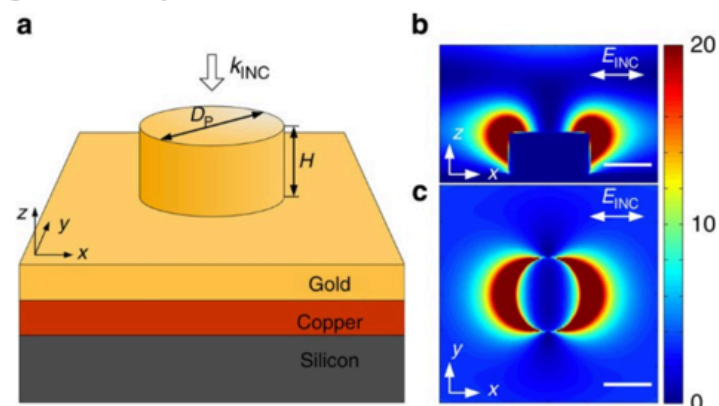
Nature Communications 2, Article number: 469 (2011) | Cite this article

Figure 3: Trapping and manual rotation of polystyrene nanospheres by gold nanopillars.



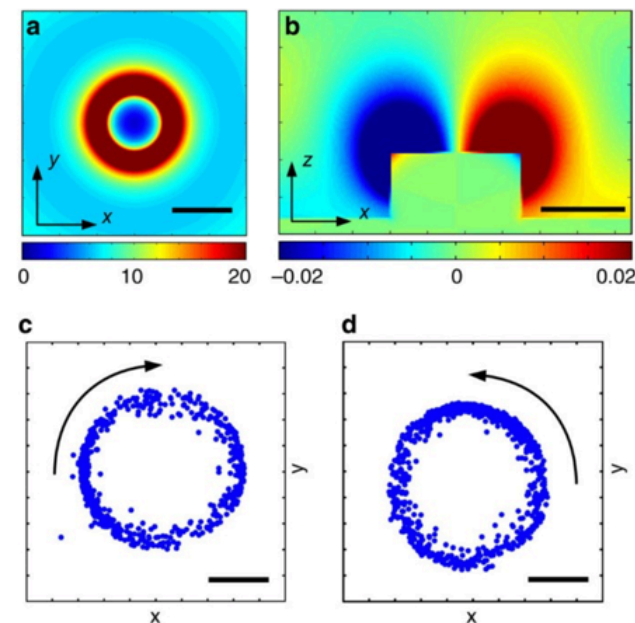
(a) Fluorescence images, obtained at successive times, of trapping and rotating 110 nm diameter polystyrene sphere by gold nanopillar. At time t_1 , sphere is close to nanopillar, but not trapped. At

Figure 1: Gold nanopillar tweezer.



(a) Plasmonic nano-tweezer comprising nanopillar formed on gold film. Underlying copper film and silicon substrate act as heat sink, conducting heat from nanopillar to substrate, thereby minimizing water heating. Nanopillar diameter D_p is 280 nm, and height H is 130 nm. (b,c) FDTD calculation of the electric field intensity distribution resulting from normal incidence plane wave illumination polarized along x -axis (E_{inc}) at $\lambda=974$ nm. Intensity enhancement, that is, intensity normalized to incident intensity $|E|^2/E_{inc}^2$, is plotted. Peak intensity enhancement is 490 times, although upper limit of colourscale is chosen to be 20 times for visualization. Scale bars, 200 nm.

Figure 4: Passive rotation of polystyrene spheres by gold nanopillar.

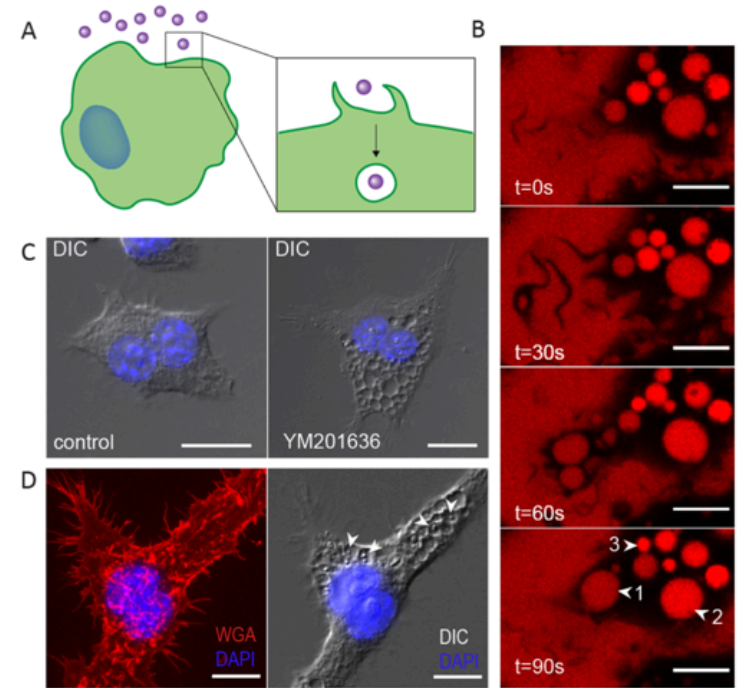
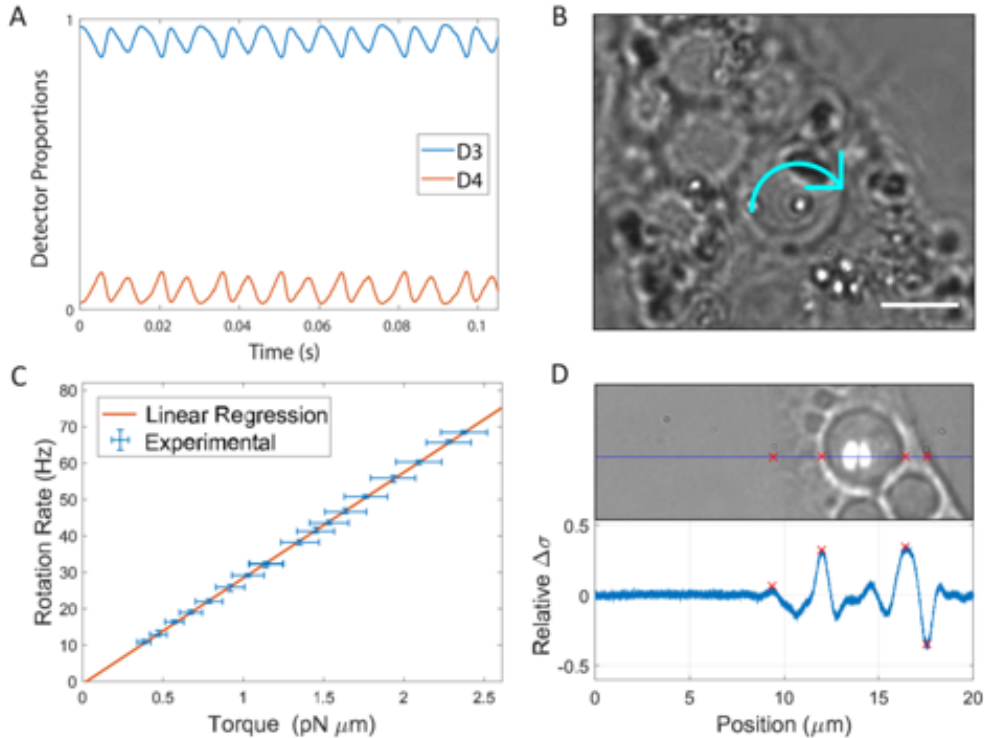
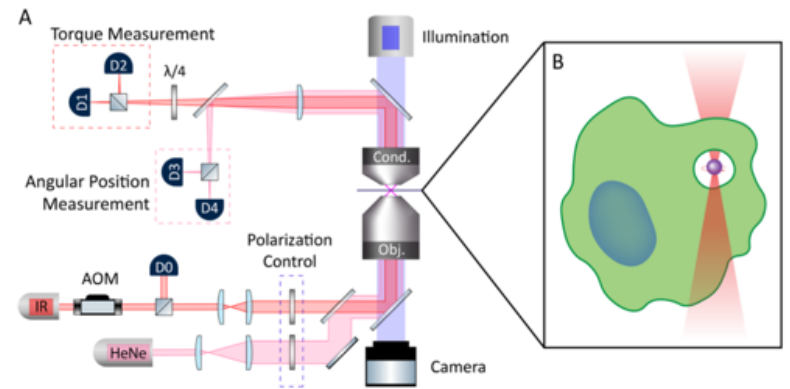


(a) FDTD calculation of electrical field intensity enhancement $|E|^2/E_{inc}^2$ distribution resulting from circularly polarized illumination at $\lambda=974$ nm. (b) FDTD calculation of y -component of Poynting

Rotational optical tweezers for active microrheometry within living cells

MARK L. WATSON,¹ DARREN L. BROWN,² ALEXANDER B. STILGOE,¹ JENNIFER L. STOW,² AND HALINA RUBINSZTEIN-DUNLOP^{1,3,*}

$$(D1 - D2)/(D1 + D2)$$



Advantage over translational motion: microrheology centred on one particle. Potential reduction of boundary effects

Structured light can possess angular momentum: rotation

$$\mathbf{j} = \epsilon_0 \left[\mathbf{r} \times \langle \mathbf{E} \times \mathbf{B} \rangle \right]$$

Allen et al Phys Rev A (1992)

Spin: due to polarisation state (rotating E-field)

$\pm \hbar$ per photon

$|L\rangle$ $J_z = +\hbar$ $|L\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \\ 0 \end{pmatrix} e^{i(kz - \omega t)}$

$|R\rangle$ $J_z = -\hbar$ $|R\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \\ 0 \end{pmatrix} e^{i(kz - \omega t)}$

Orbital: due to inclined wavefronts

$l\hbar$ per photon

Laguerre-Gaussian (LG) beam

$$\mathbf{S} = \epsilon_0 \int (\mathbf{E} \times \mathbf{A}) d^3 \mathbf{r}.$$

➔
$$\mathbf{S} = \frac{\epsilon_0}{2i\omega} \int (\mathbf{E}^* \times \mathbf{E}) d^3 \mathbf{r}.$$

$$\mathbf{S} = \sum_{\mathbf{k}} \hbar \mathbf{u}_{\mathbf{k}} \left(\hat{a}_{\mathbf{k},L}^\dagger \hat{a}_{\mathbf{k},L} - \hat{a}_{\mathbf{k},R}^\dagger \hat{a}_{\mathbf{k},R} \right)$$

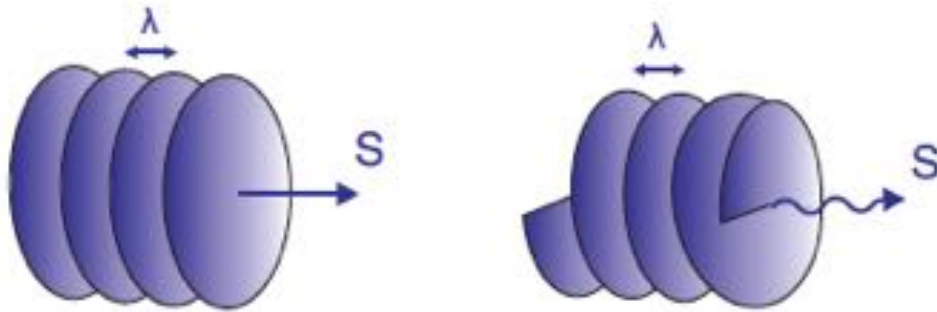
$$\mathbf{S}_z = \pm \hbar.$$



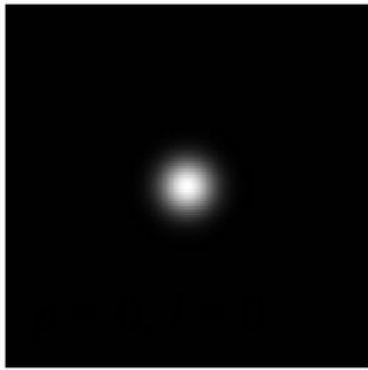
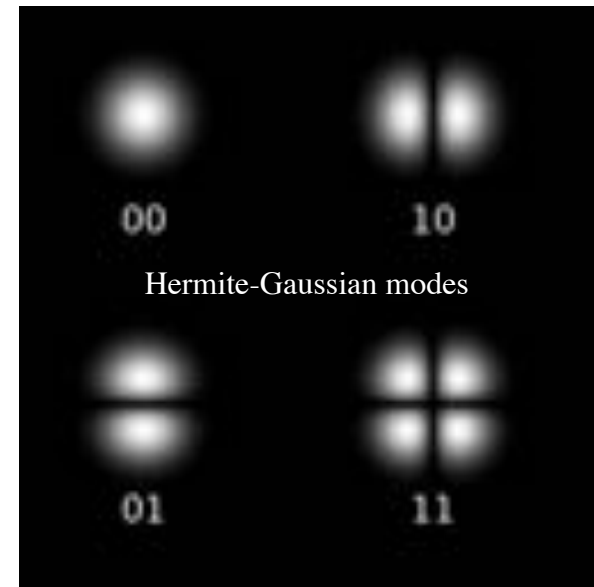
Laguerre-Gaussian modes

L. Allen et al., Physical Review A **45**, 8185 (1992)

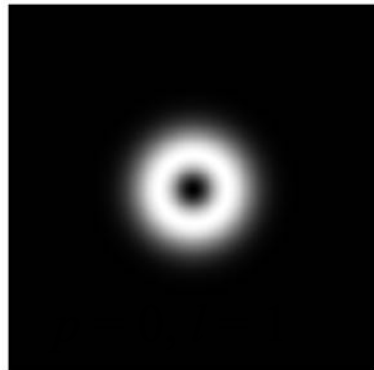
- radial mode index p (determines radial structure)
- azimuthal mode index l (determines helicity)



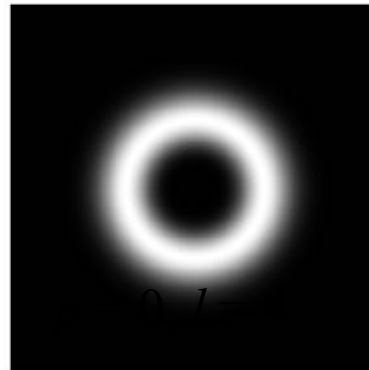
$l\hbar$ per photon



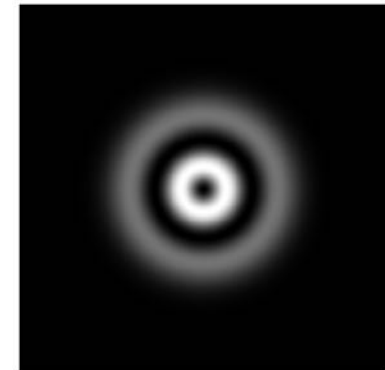
$p=0, l=0$



$p=0, l=1$

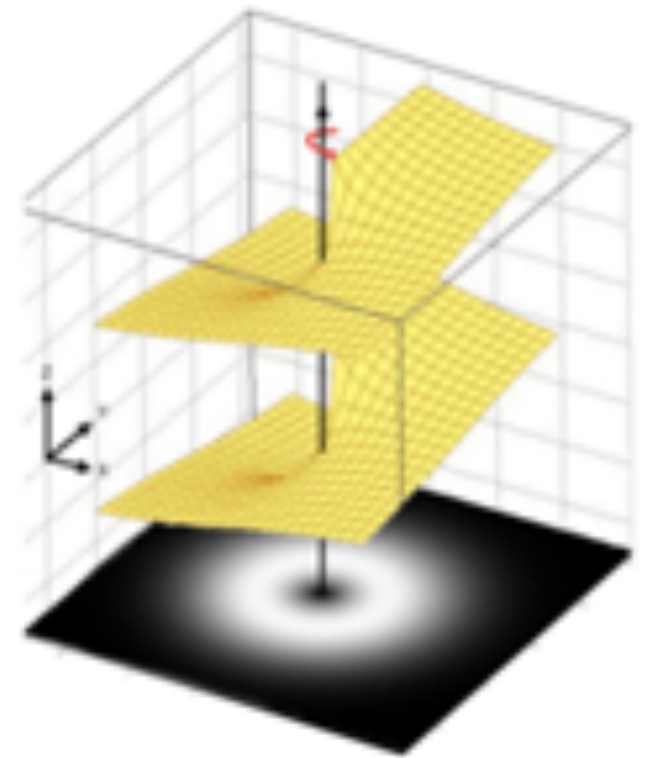


$p=0, l=3$



$p=1, l=1$

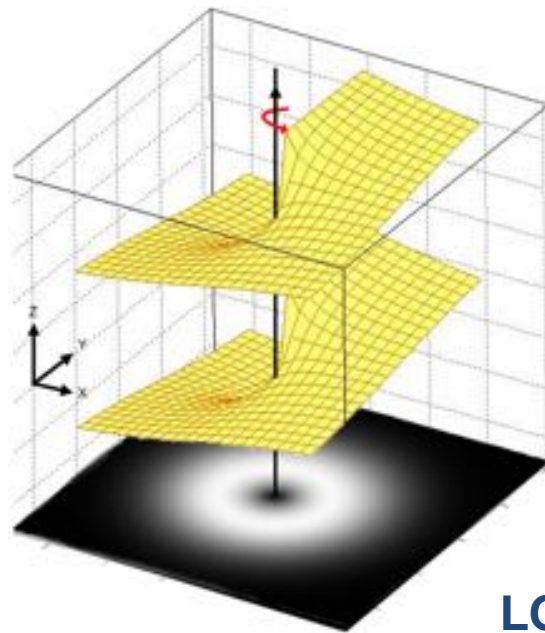
Laguerre-Gaussian modes have **orbital angular momentum**
– **due to inclined wavefronts**



Laguerre-Gaussian (LG) beams
Allen et al., Phys Rev A 1992

Light for rotation: orbital angular momentum transfer

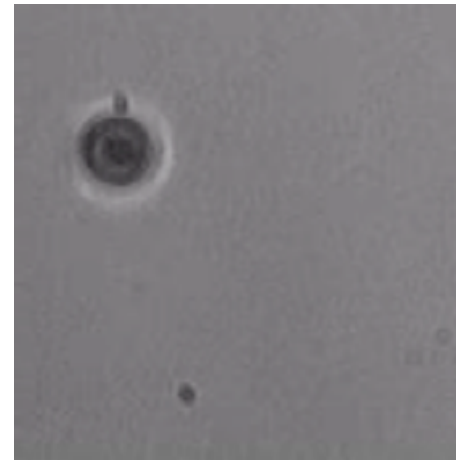
Orbital angular momentum (**OAM**) transfer by inclined wavefronts via light scattering



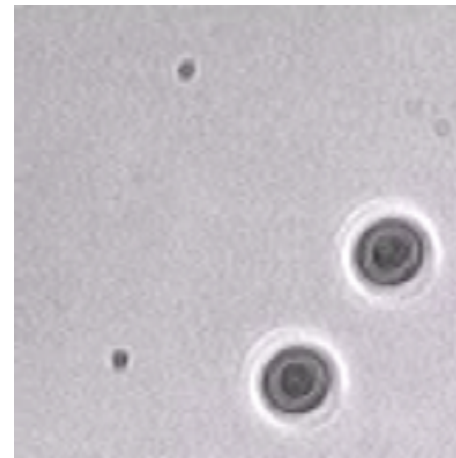
LG beam
($\ell=1, p=0$)

$\ell\hbar$ per photon

Silica
($3\mu\text{m}$)

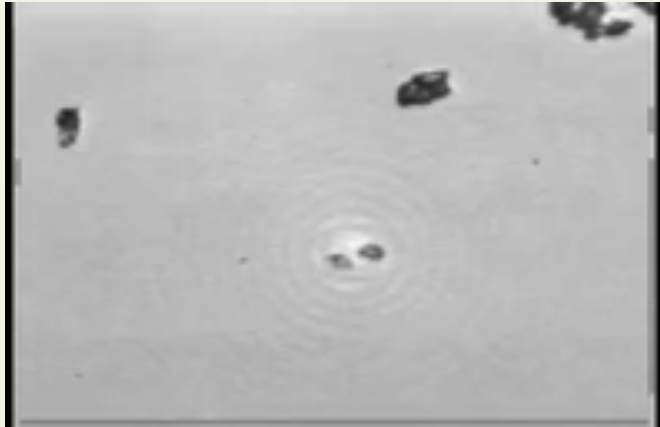
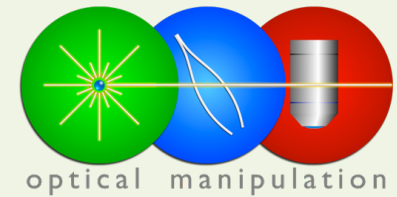


$\ell=10$

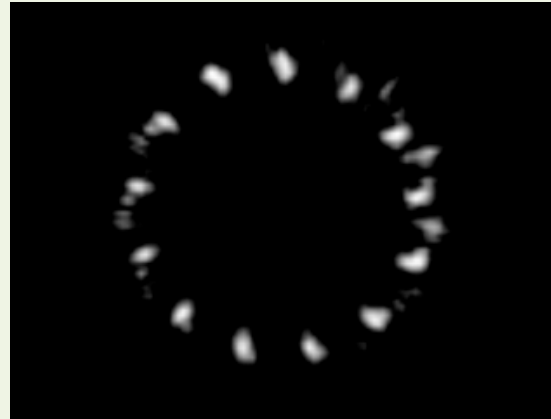


$\ell=-10$

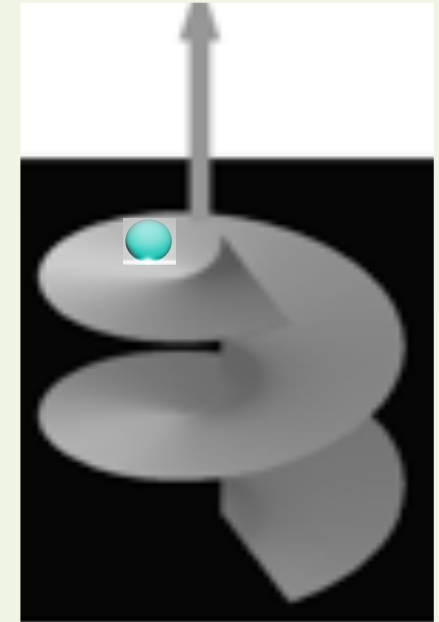
Orbital angular momentum transfer to trapped particles



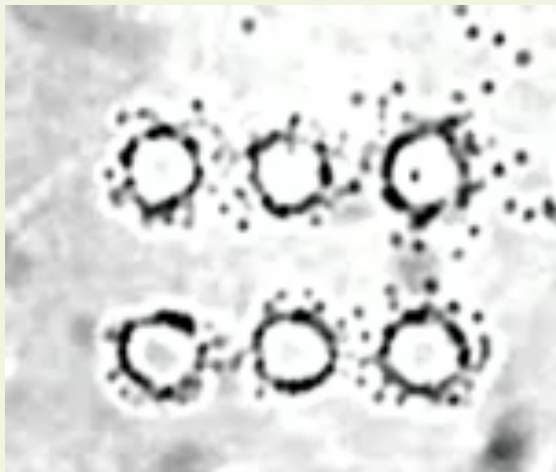
Garces-Chavez et al., Physical Review Letters (2003)
O'Neil et al., Physical Review Letters (2002)



Chen et al. Optics Letters (2013)



New concepts to look at
spin-orbit coupling
negative torque

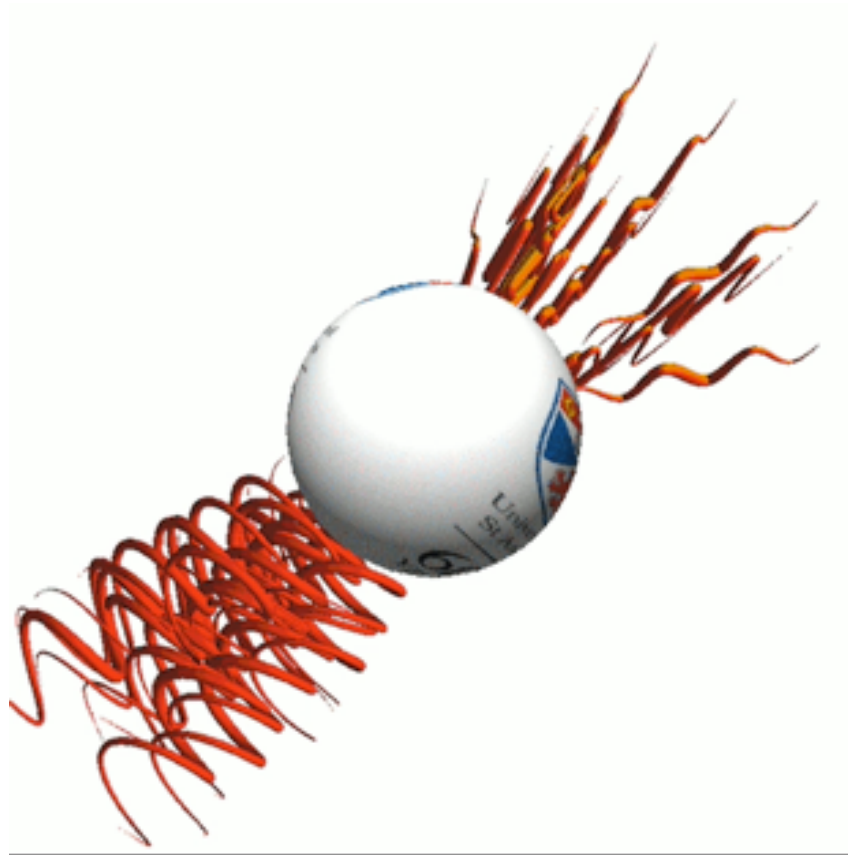


This video from Kosta Ladavac and David Grier, "Microoptomechanical pumps assembled and driven by holographic optical vortex arrays," Opt. Express **12**, 1144-1149 (2004)

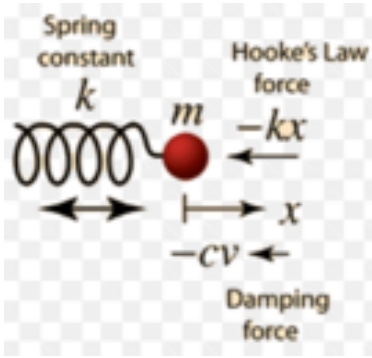
see also related work by other groups, including:

Ristch-Marte group, Innsbruck, Austria
Padgett Group, University of Glasgow, UK
Rubenzstein-Dunlop group, Brisbane (did first work on rotation - based on absorption of CuO particles (He et al PRL, 1995))

Rotation in optical traps in vacuum



Trap in vacuum? Why?



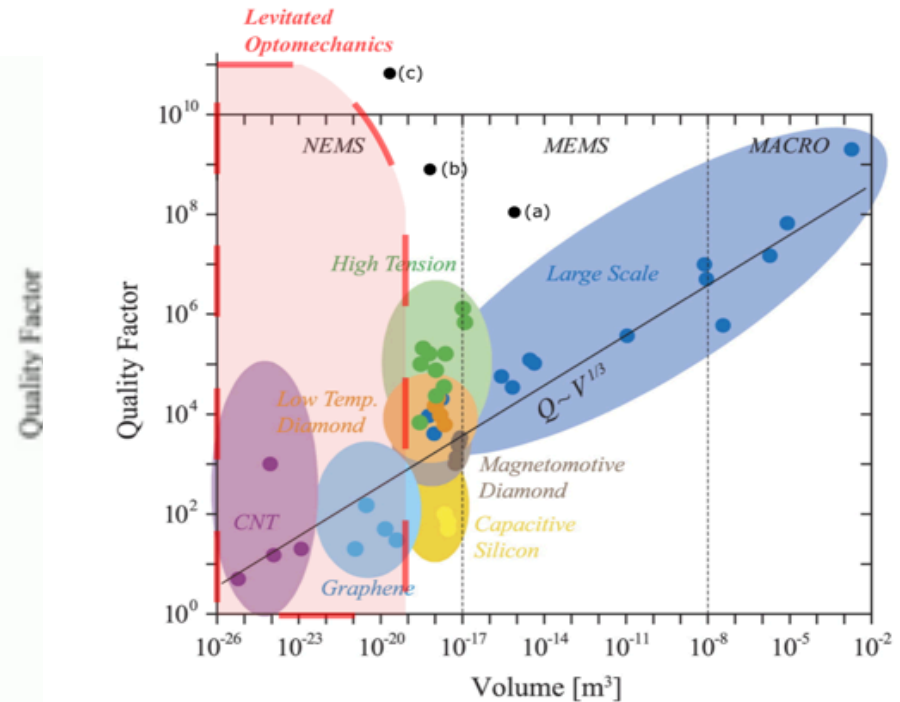
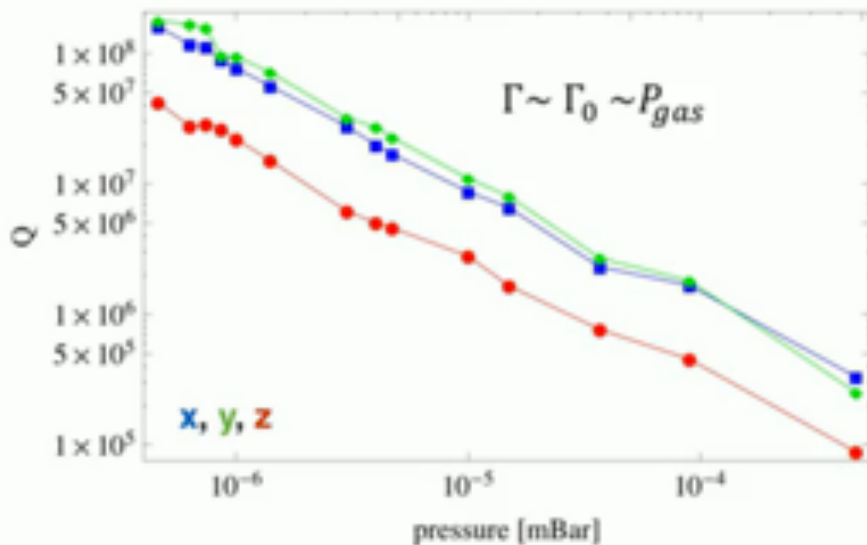
$$F \propto \frac{1}{\sqrt{Q}}$$

$$\langle n \rangle = k_B T_{\text{cm}} / (\hbar \Omega_0)$$

Optical levitation of a trapped sphere offers:

high vibrational frequencies, record rotational speeds, and mechanical Q factors exceeding 10^{12} suggested.

Geraci group: zeptonewton force sensing. St Andrews has shown $Q \sim 10^8$ (Science Advances, (2020))



Optomechanics with levitated particles

Rep. Prog. Phys. **83** (2020) 026401 (36pp)

J. Gieseler, B. Deutsch, R. Quidant, L. Novotny, Phys. Rev. Lett. **109**, 103603 (2012)

REPORT

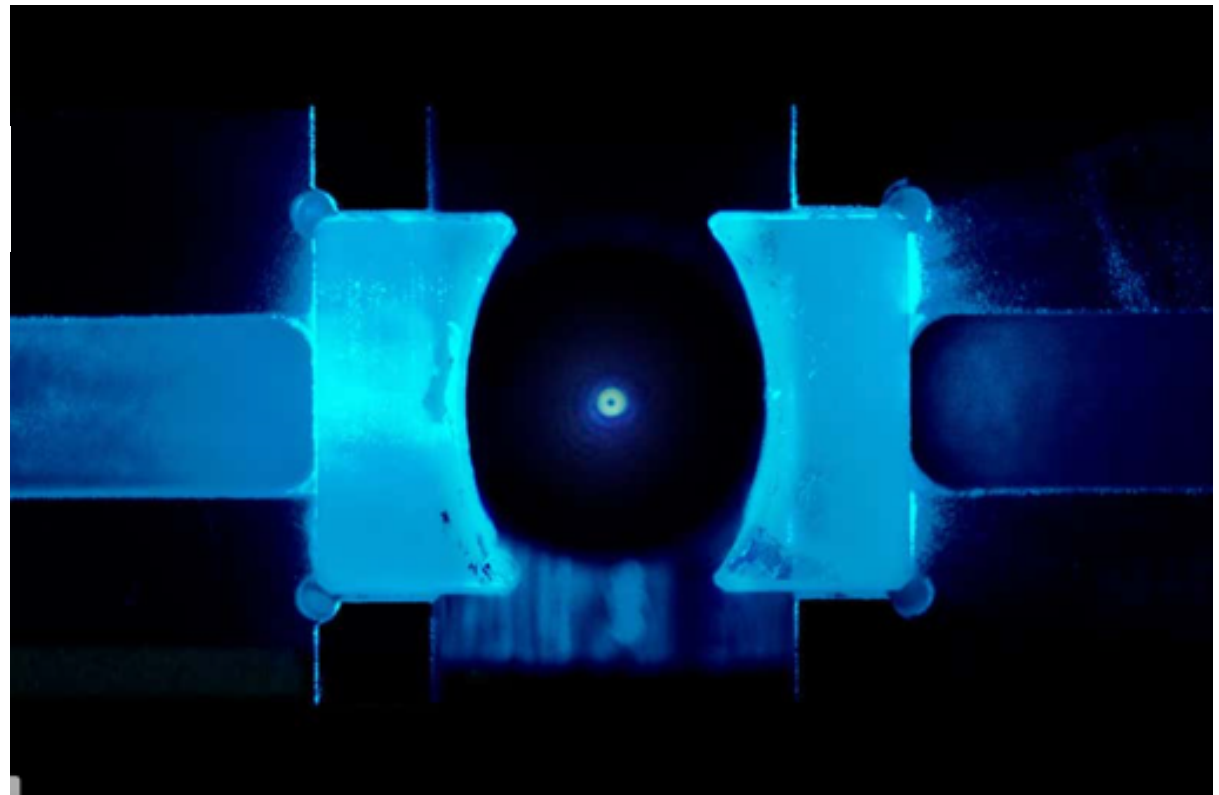
Cooling of a levitated nanoparticle to the motional quantum ground state

Uroš Delić^{1,2,*}, Manuel Reisenbauer¹, Kahan Dare^{1,2}, David Grass^{1,†}, Vladan Vuletić³, Nikolai Kiesel¹, Markus Aspelmeyer^{1,2,*}

♦ See all authors and affiliations

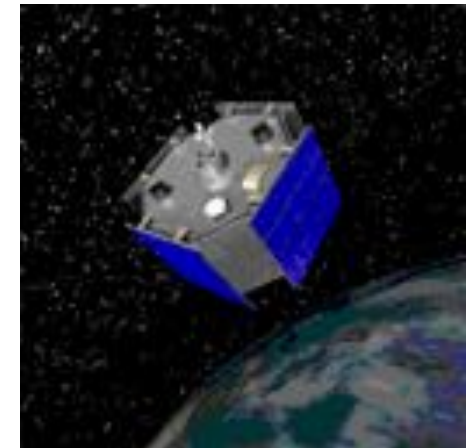
Science 21 Feb 2020:
Vol. 367, Issue 6480, pp. 892-895
DOI: 10.1126/science.aba3993

NewScientist



This tiny glass bead has been quantum chilled to near absolute zero

Rotational levitated optomechanics: all in a spin



PRL 105, 113601 (2010)

PHYSICAL REVIEW LETTERS

week ending
10 SEPTEMBER 2010

Vacuum Friction in Rotating Particles

A. Manjavacas and F.J. García de Abajo*
Instituto de Óptica—CSIC, Serrano 121, 28006 Madrid, Spain
(Received 8 March 2010; published 8 September 2010)

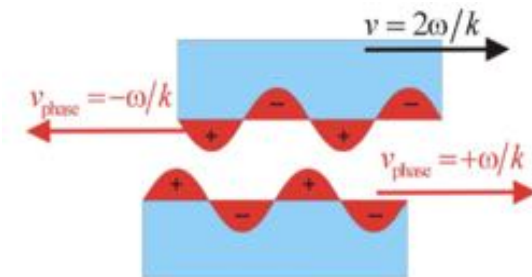
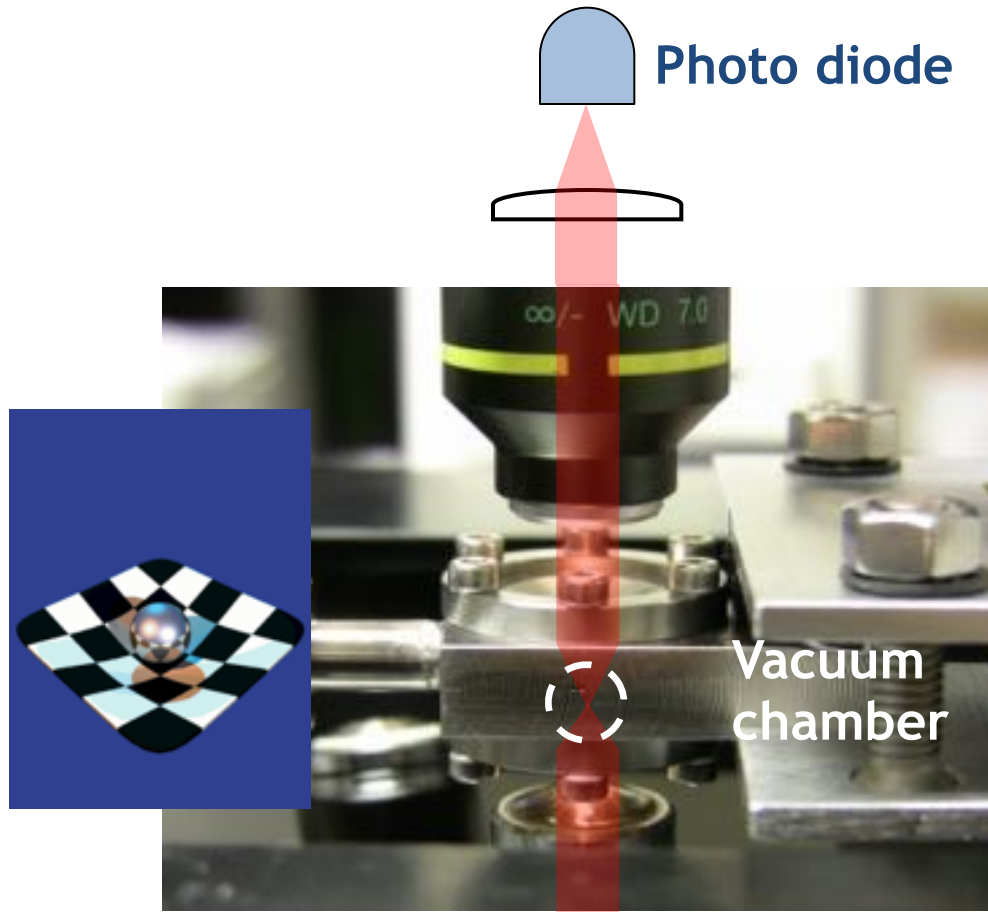
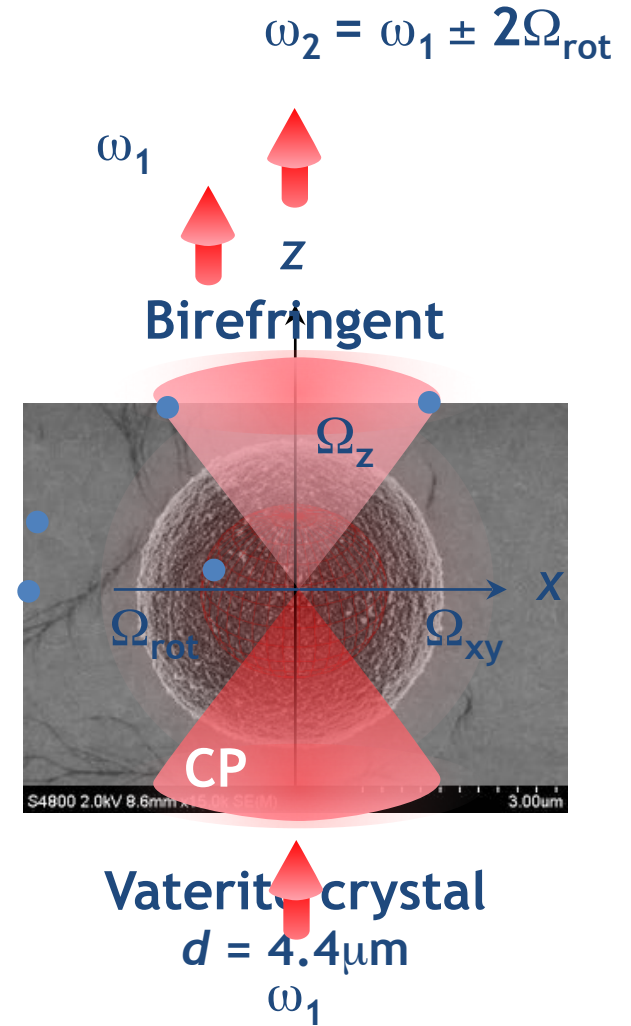


Figure adapted from: Pendry J B, *Quantum friction- fact or fiction?*, *New J. Phys.*, (2010)

Experiment: trap and rotate in air or vacuum



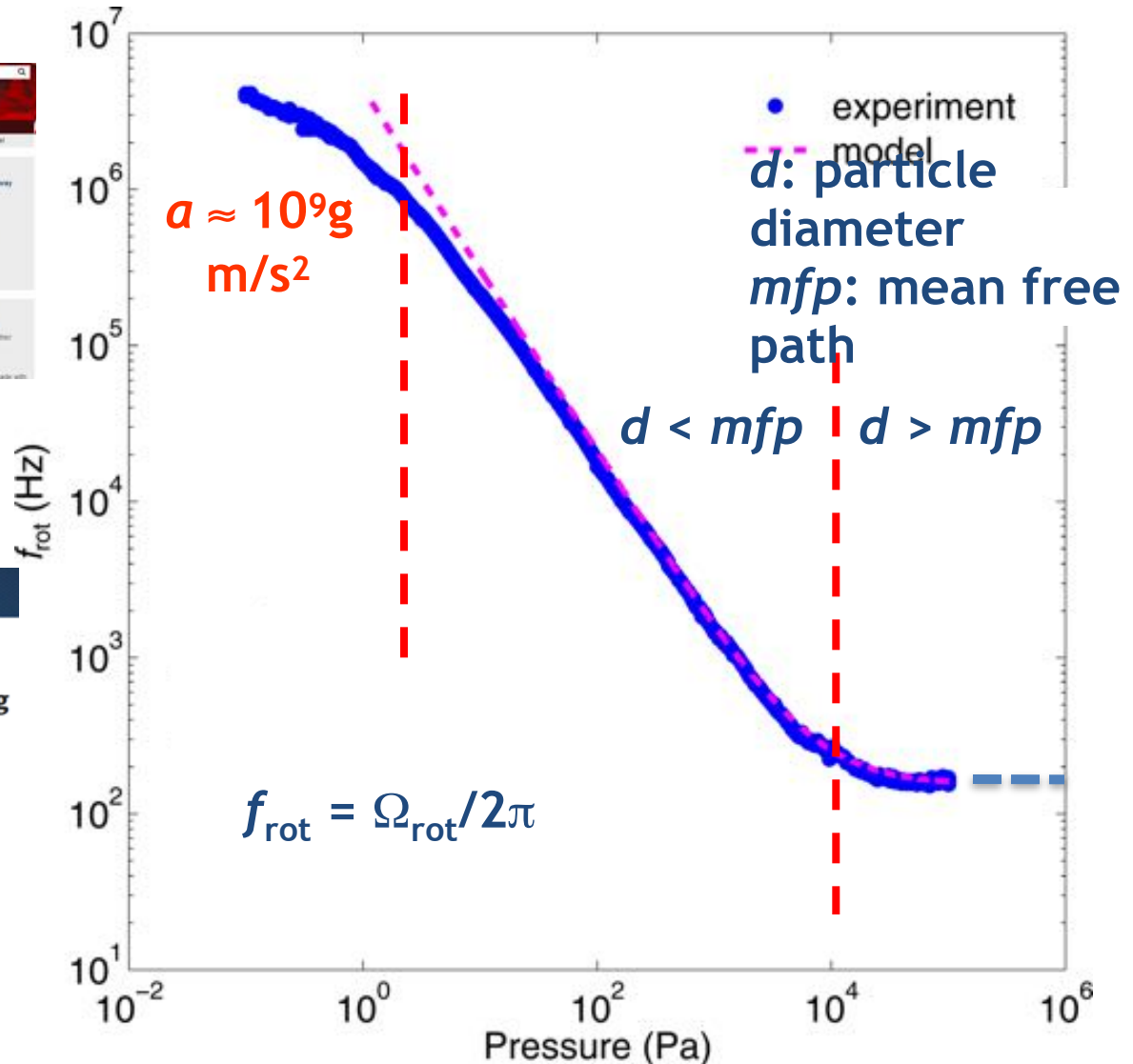
Circularly polarised (CP) trapping beam (1070nm)



Rotation versus pressure for the microgyroscope

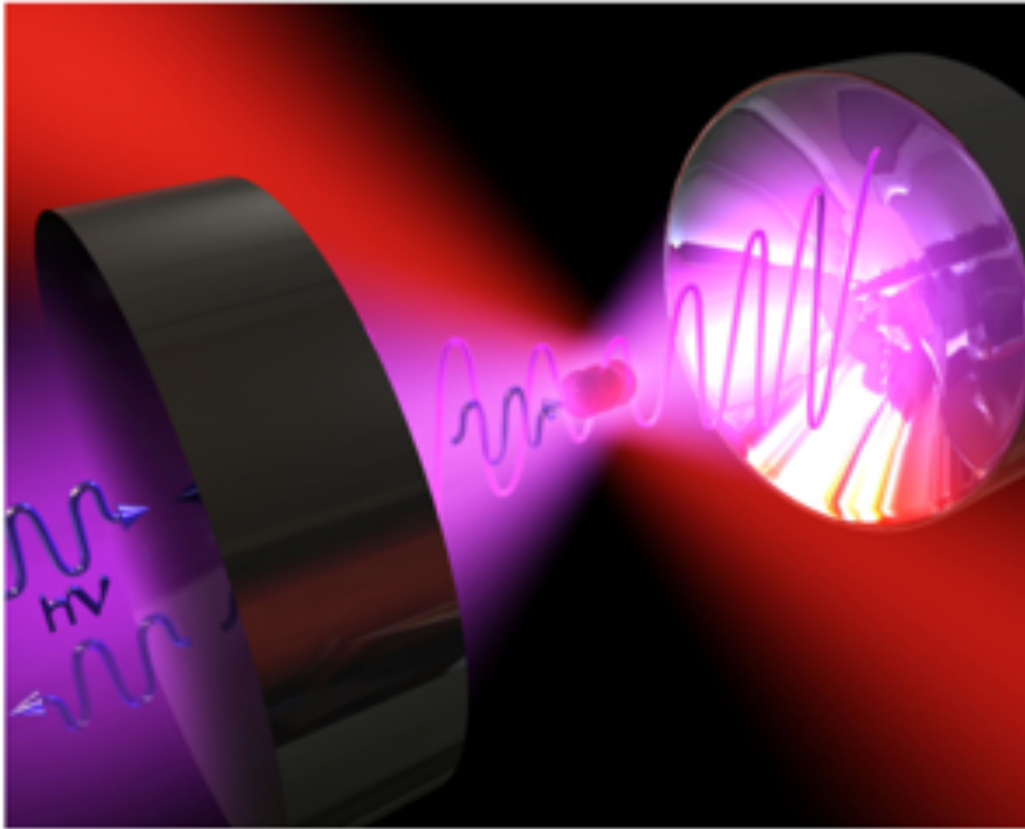


Arita et al. Nature Comm
4, 2374 (2013)



Guinness book of world records 2015 (fastest man-made rotation)





Common influenza viruses, with a size of ~ 100 nm, can be stored for several weeks in vacuum down to 10^{-4} torr.

Due to their structure (e.g. lipid bilayer, nucleocapsid protein and DNA), viruses present a transparency window at the optical wavelength which yields relatively low bulk temperatures

Figure 3. Quantum superposition of living organisms. Illustration of the protocol to create quantum superposition states applied to living organisms, such as viruses, trapped in a high-finesse optical cavity by optical tweezers.

Toward quantum superposition of living organisms

**“When I described catching living things with light
people said: ‘Don’t exaggerate Ashkin’.”**



THE NOBEL PRIZE
IN PHYSICS 2018



OE Magazine, SPIE 2013

Acknowledgements

Yoshi Arita
Mingzhou Chen
Steven Neale
Veneranda Garces-Chavez



EPSRC

Pioneering research
and skills



Australian Government

Australian Research Council

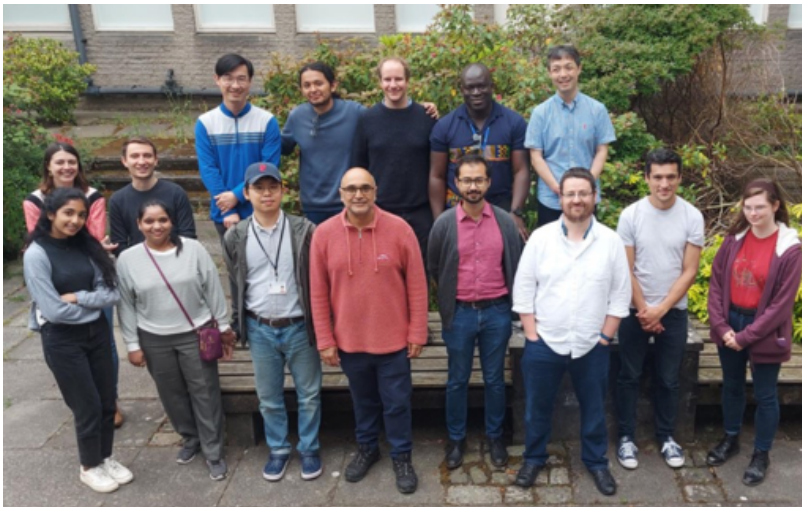
Theory collaborators

Ewan Wright
Pavel Zemanek
Stephen Simpson



University of
St Andrews | **FOUNDED**
1413 |

Early work on rotation with Michael Mazilu



ADVANCES IN PHYSICS: X
2020, VOL. 6, NO. 1, 10.1080/23746149.2020.1838322
<https://doi.org/10.1080/23746149.2020.1838322>



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REVIEW

OPEN ACCESS Check for updates

**Initiating revolutions for optical manipulation: the origins
and applications of rotational dynamics of trapped
particles**

Graham D. Bruce ^{a,*}, Paloma Rodríguez-Sevilla ^{a,*} and
Kishan Dholakia ^{a,b}