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FOUNDATION FOR RESEARCH AND TECHNOLOGY – HELLAS



**University of Crete**

# Combining chirality with Parity-Time (PT) symmetry in metamaterials

**M. Kafesaki**, I. Katsantonis, S. Droulias, E.N. Economou, and C. M.  
Soukoulis

Foundation for Research & Technology - Hellas (**FORTH**), Crete, Greece, and  
**University of Crete, Greece**

Ames Lab & Iowa State University (ISU), USA

*[kafesaki@iesl.forth.gr](mailto:kafesaki@iesl.forth.gr)*

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Photonic Metamaterials



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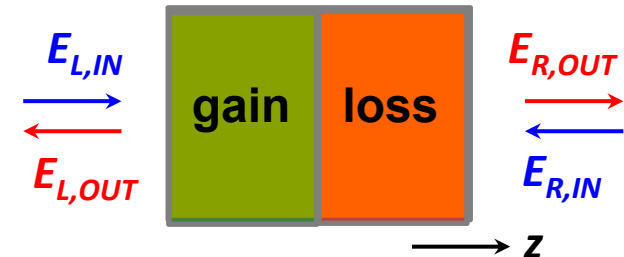
*[kafesaki@iesl.forth.gr](mailto:kafesaki@iesl.forth.gr)*

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# Outline

- Introduction
  - Parity-Time (PT) symmetry
  - **chirality** and chiral metamaterials
- Combining **chirality with PT symmetry**
  - **conditions** for PT symmetry with chirality
  - **scattering** by a PT chiral double-slab
    - normal incidence
    - oblique incidence
- Practical implementations
- PT symmetry for **molecular chirality sensing**



# The concept of Parity-Time (PT) symmetry

**Schrödinger Equation:**  $\hat{H}\Psi(x) = E\Psi(x)$

Hamiltonian

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$

**Parity-Time-symmetric Hamiltonians can have real eigenvalues even if not-Hermitian**

*Bender, Phys. Rev. Lett. 80, 5243, 1998*

**PT-symmetric Hamiltonian:**  $[\hat{P}\hat{T}, \hat{H}] = 0 \Rightarrow \hat{P}\hat{T}\hat{H} = \hat{H}\hat{P}\hat{T}$

**Parity: spatial inversion**

$$\hat{P}: \begin{cases} \mathbf{r} \rightarrow -\mathbf{r} \\ \mathbf{p} \rightarrow -\mathbf{p} \\ \mathbf{E} \rightarrow -\mathbf{E} \text{ polar vector} \\ \mathbf{H} \rightarrow \mathbf{H} \text{ axial vector} \end{cases} \quad \hat{T}: \begin{cases} t \rightarrow -t \\ i \rightarrow -i \\ \mathbf{r} \rightarrow \mathbf{r} \\ \mathbf{E} \rightarrow \mathbf{E} \\ \mathbf{H} \rightarrow -\mathbf{H} \end{cases}$$

# The concept of Parity-Time (PT) symmetry

**PT-symmetric Hamiltonians can have real eigenvalues even if not-Hermitian**

$$\hat{H}\Psi(x) = E\Psi(x)$$

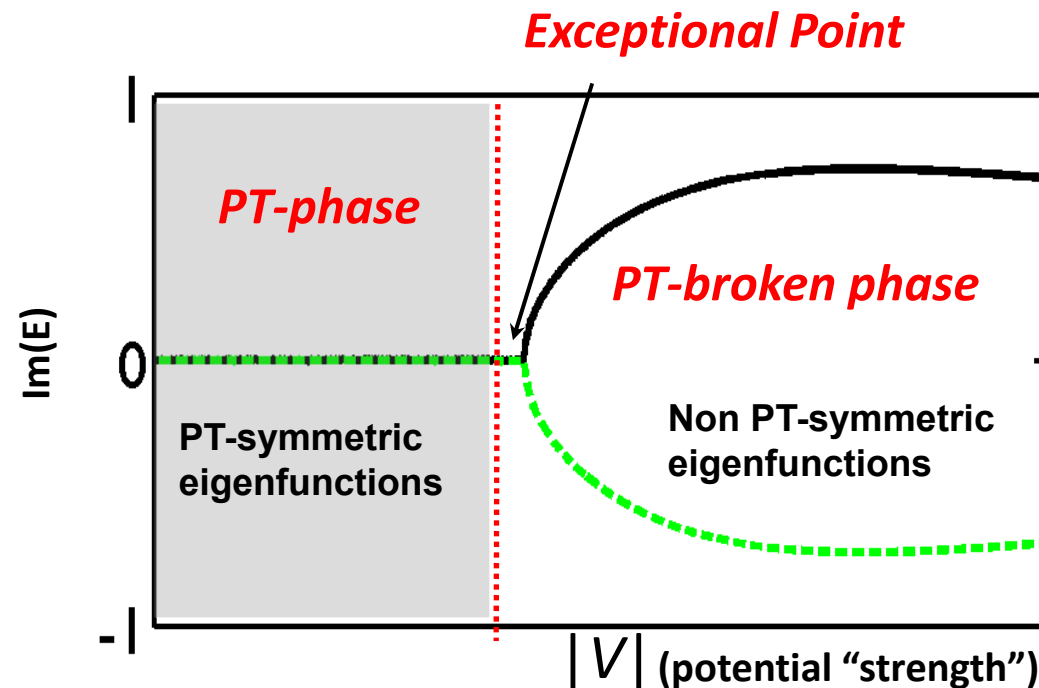
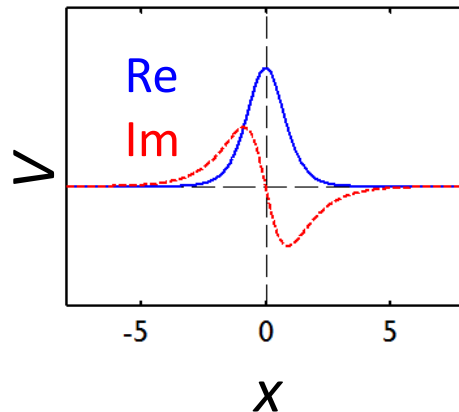
$$[\hat{P}\hat{T}, \hat{H}] = 0 \Rightarrow \hat{P}\hat{T}\hat{H} = \hat{H}\hat{P}\hat{T}$$

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$

**Condition:  $V(x) = V^*(-x)$**

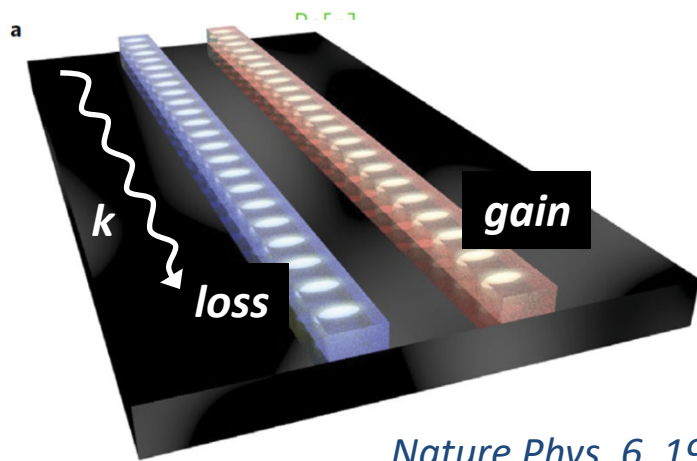
$$\begin{cases} \text{Re}[V(x)] = \text{Re}[V(-x)] \text{ (even in } x\text{)} \\ \text{Im}[V(x)] = -\text{Im}[V(-x)] \text{ (odd in } x\text{)} \end{cases}$$

Example of PT-symmetric potential



# Parity-Time (PT) symmetry in optics

*In optics, wave propagation in paraxial approximation is governed by a Schrödinger-like equation*



*Nature Phys. 6, 192 (2010)*

$n(x)$ : réfractive index  
 $k_0$ : free space wave number

$$\overbrace{\left[ -\frac{1}{2k_0 n_0} \frac{\partial^2}{\partial x^2} - \frac{\omega}{c} n(x) \right]}^{\text{Hamiltonian}} E(x) = k_z E(x)$$

*potential*

***PT-symmetric potential:  $n(x) = n^*(-x)$***

***PT-symmetry requirements  
 beyond paraxial approximation***

$$\epsilon(\mathbf{r}) = \epsilon^*(-\mathbf{r})$$

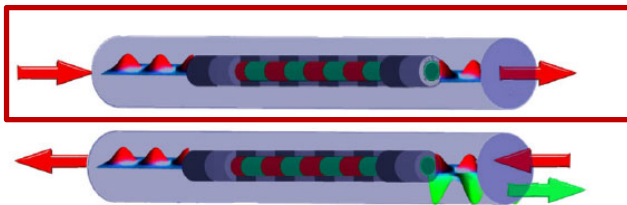
$$\mu(\mathbf{r}) = \mu^*(-\mathbf{r})$$

*Phys Rev Lett. 110,  
 173901 (2013)*

***Realization by proper  
 combination of loss and  
 gain media***

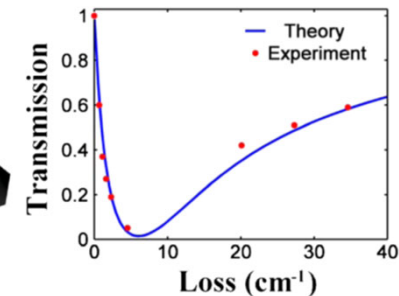
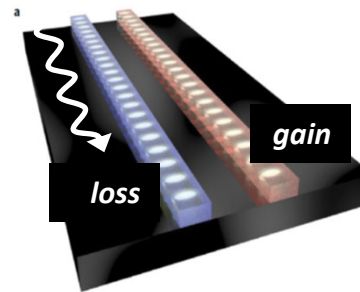
# Novel effects in PT symmetric optical systems

## Unidirectional reflectionless propagation and invisibility



PRL 106, 213901 (2011)

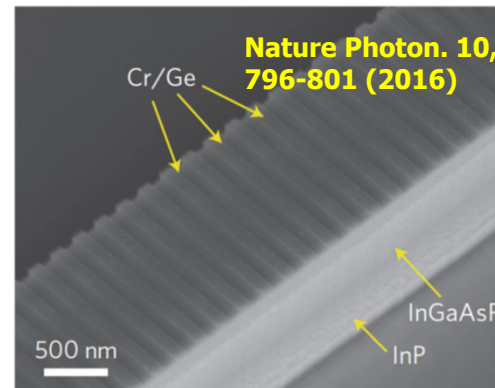
## Loss-induced transparency



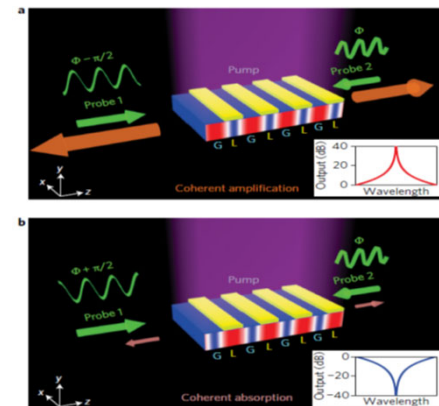
PRL 103, 093902(2009), Nature Phys. 14, 11 (2018) -review

## Coherent perfect absorption (CPA) and lasing at same frequency

PRA 82, 031801(R) (2010), PRL 106, 093902 (2011)  
Science 346 (6212), 975 (2014)



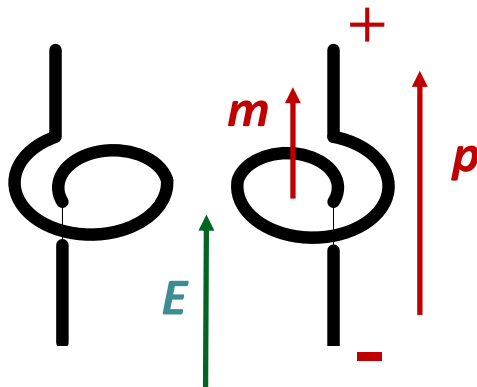
Nature Photon. 10, 796-801 (2016)



**How can we extend these effects in chiral systems?**

# Chirality and chiral metamaterials

**Chiral structure: not identical to its mirror image**



chirality is expressed via the parameter  $\kappa$

$$\mathbf{D} = \epsilon\epsilon_0\mathbf{E} + i(\kappa/c)\mathbf{H}$$

$$\mathbf{B} = \mu\mu_0\mathbf{H} - i(\kappa/c)\mathbf{E}$$

magneto-electric  
coupling

## Consequences of chirality:

1. **eigenmodes: circularly polarized waves**

- Right Circularly Polarized (RCP or +)

- Left Circularly Polarized (LCP or -)

$$E_{\pm}(\hat{x} \pm i\hat{y})$$

2. **different index for RCP (+) and LCP (-) waves!**

$$n_{\pm} = \sqrt{\epsilon\mu \pm \kappa}$$



# Chirality-originated effects

$$n_{\pm} = \sqrt{\epsilon\mu \pm \kappa}$$

**Optical activity**

Rotation of linear polarization

measured by:

**Optical Rotation**

$$\theta = \frac{1}{2} [\arg(t_{++}) - \arg(t_{--})]$$

↑  
transmittance of  
RCP (+) waves

↑  
transmittance of  
LCP (-) waves

**Circular dichroism (CD)**

Different absorption for RCP/LCP waves

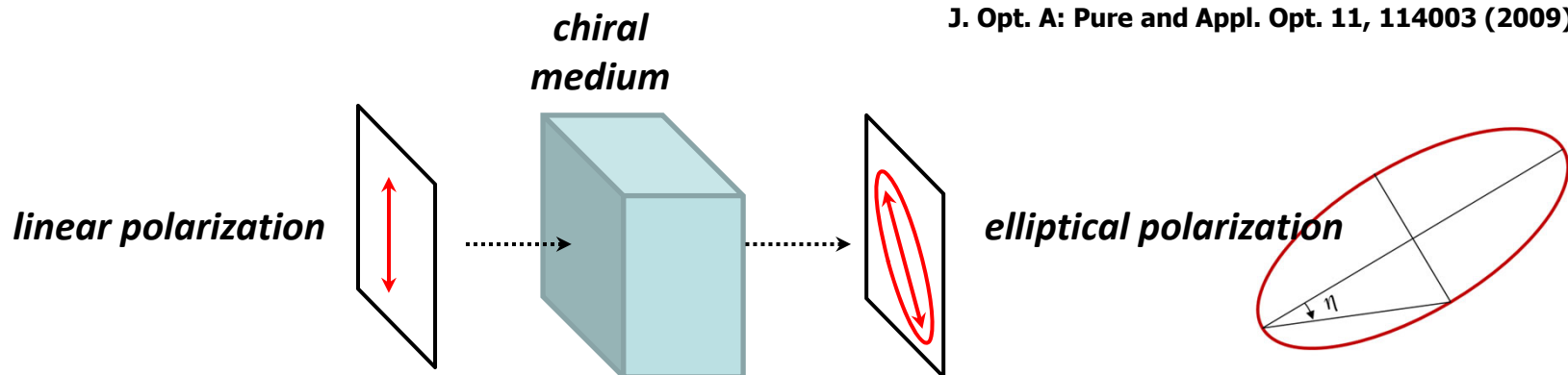
$$CD = A_{++} - A_{--}$$

measured by:

$$\text{Ellipticity } \eta = \frac{1}{2} \tan^{-1} \left( \frac{|t_{++}|^2 - |t_{--}|^2}{|t_{++}|^2 + |t_{--}|^2} \right)$$

$\eta = 0$ : linear,  $\eta = 45^\circ$ : circular

J. Opt. A: Pure and Appl. Opt. 11, 114003 (2009)



# Questions

***Can we achieve **Parity-Time symmetry** in chiral media?***

***if yes***

***Under what conditions?***

***See Droulias et. al., Phys. Rev. Lett 122, 213201 (2019)***  
***Katsantonis et. al., Photonics 7, 43 (2020)***

# To answer the questions...

**Maxwell's Eqs.**

$$\nabla \times \mathbf{E} = i\omega \mathbf{B}$$

$$\nabla \times \mathbf{H} = -i\omega \mathbf{D}$$

**Constitutive relations**

$$\mathbf{D} = \varepsilon\varepsilon_0 \mathbf{E} + i(\kappa/c)\mathbf{H}$$

$$\mathbf{B} = \mu\mu_0 \mathbf{H} - i(\kappa/c)\mathbf{E}$$

$$\hat{H} \begin{bmatrix} \mathbf{B} \\ \mathbf{D} \end{bmatrix} = \omega \begin{bmatrix} \mathbf{B} \\ \mathbf{D} \end{bmatrix}$$

$$\hat{\Omega}(\mathbf{r}) = \begin{bmatrix} 0 & -\frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & -\frac{\partial}{\partial x} \\ -\frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \end{bmatrix}$$

$$\hat{H} = \left[ \begin{array}{c|c} -i\hat{\Phi}_{11}(\mathbf{r}) - i\phi_{11}(\mathbf{r})\hat{\Omega}(\mathbf{r}) & -i\hat{\Phi}_{12}(\mathbf{r}) - i\phi_{12}(\mathbf{r})\hat{\Omega}(\mathbf{r}) \\ \hline +i\hat{\Phi}_{21}(\mathbf{r}) + i\phi_{21}(\mathbf{r})\hat{\Omega}(\mathbf{r}) & +i\hat{\Phi}_{22}(\mathbf{r}) + i\phi_{22}(\mathbf{r})\hat{\Omega}(\mathbf{r}) \end{array} \right]$$

$$\hat{\Phi}_{ij}(\mathbf{r}) = \begin{bmatrix} 0 & -\frac{\partial\phi_{ij}(\mathbf{r})}{\partial z} & \frac{\partial\phi_{ij}(\mathbf{r})}{\partial y} \\ \frac{\partial\phi_{ij}(\mathbf{r})}{\partial z} & 0 & -\frac{\partial\phi_{ij}(\mathbf{r})}{\partial x} \\ -\frac{\partial\phi_{ij}(\mathbf{r})}{\partial y} & \frac{\partial\phi_{ij}(\mathbf{r})}{\partial x} & 0 \end{bmatrix}$$

$$\hat{P}\hat{T}\hat{H} = \hat{H}\hat{P}\hat{T}$$

$$\phi_{11}(\mathbf{r}) = -\phi_{22}(\mathbf{r}) = -\frac{i\kappa/c}{(\varepsilon\mu - \kappa^2)/c^2}$$

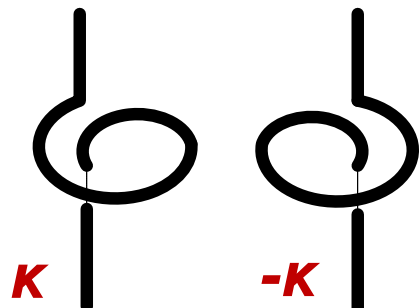
See Droulias et. al., Phys. Rev. Lett 122, 213201 (2019)

# Questions

Can we achieve **Parity-Time symmetry in chiral media?**

**if yes**

**Under what conditions?**



$$\begin{aligned}\epsilon(\mathbf{r}) &= \epsilon^*(-\mathbf{r}) \\ \mu(\mathbf{r}) &= \mu^*(-\mathbf{r}) \\ \kappa(\mathbf{r}) &= -\kappa^*(-\mathbf{r})\end{aligned}$$

as in non-chiral systems

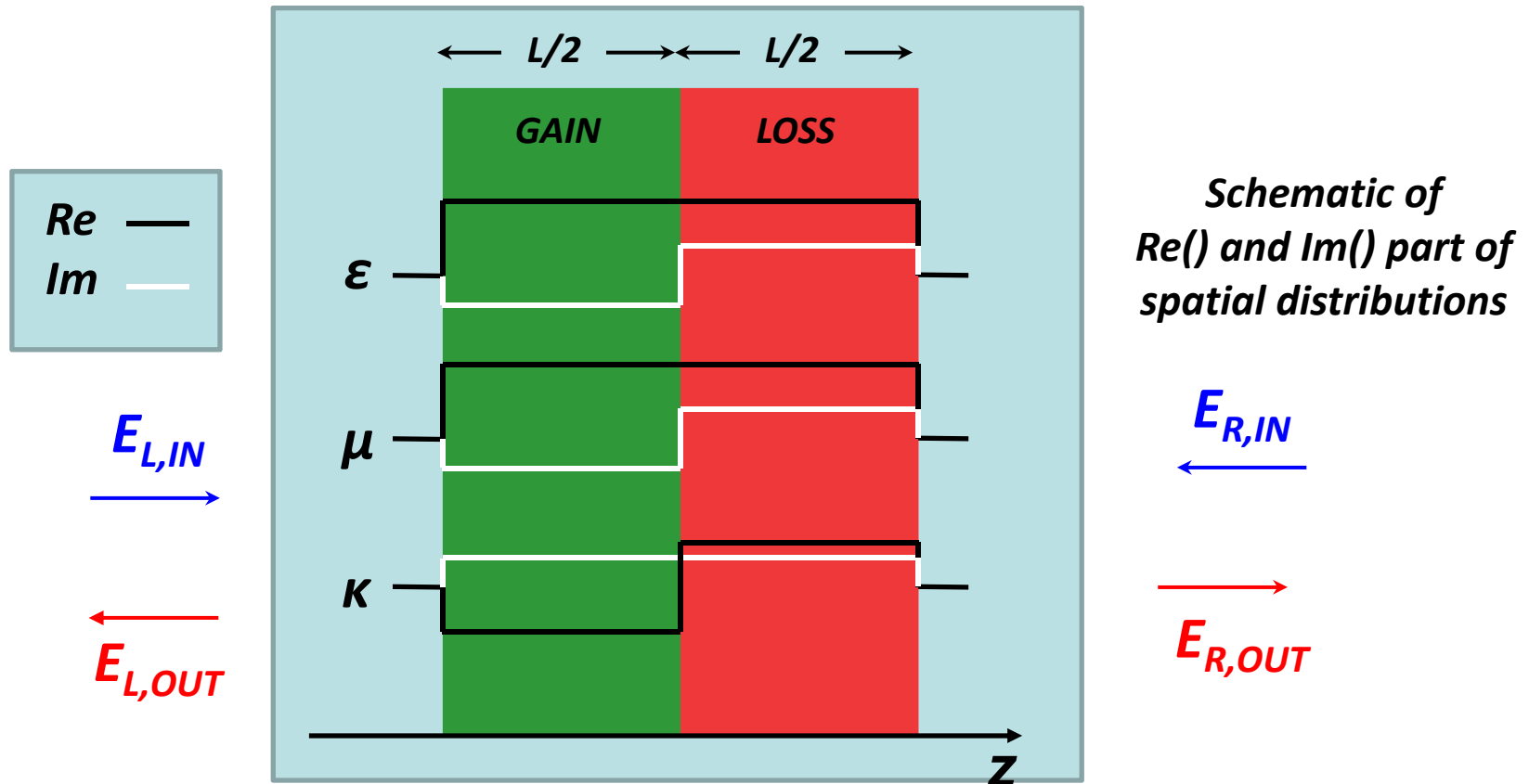
← additional condition

See Droulias et. al., *Phys. Rev. Lett* 122, 213201 (2019)  
Katsantonis et. al., *Photonics* 7, 43 (2020)

# Our PT-chiral system/bilayer

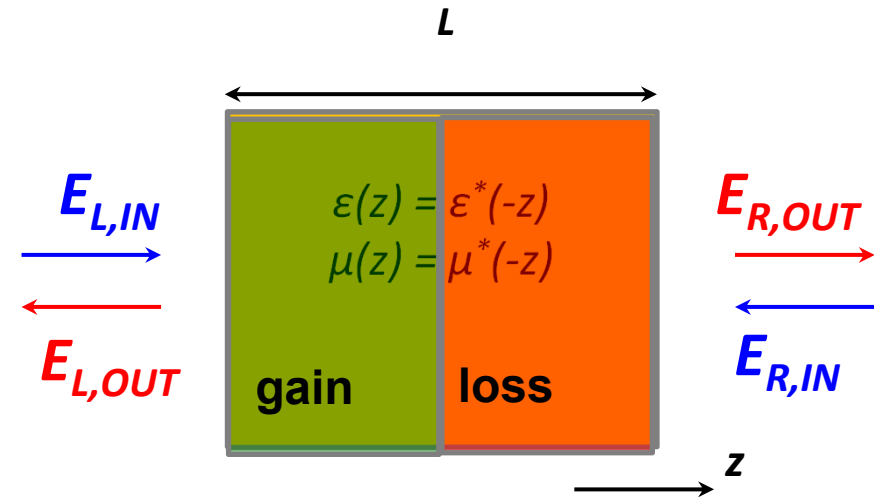
*Homogenous, non-dispersive, chiral media*

*Investigation of scattering (transmission/reflection) properties*



# Scattering by Parity-Time (PT) symmetric systems

The *PT*-symmetry phases can be identified from the eigenvalues (*s*) of the scattering matrix

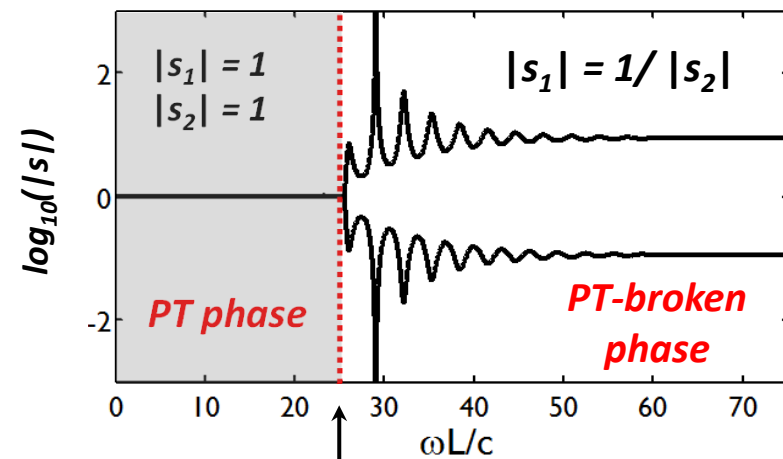


$$\begin{bmatrix} E_{L,OUT} \\ E_{R,OUT} \end{bmatrix} = \begin{bmatrix} r^{(L)} & t^{(R)} \\ t^{(R)} & r^{(R)} \end{bmatrix} \begin{bmatrix} E_{L,IN} \\ E_{R,IN} \end{bmatrix}$$

Scattering matrix

$t^{(L)}$  transmission for left-incidence

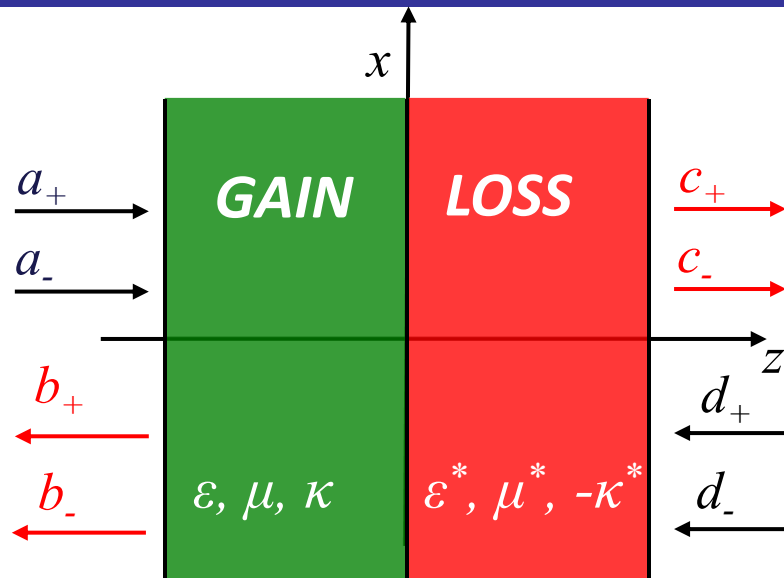
$r^{(R)}$  reflection for right-incidence



**PT-phase**  $\longleftrightarrow$  **Unimodular eigenvalues**

(\*) PRL 106, 093902 (2011)

# The "chiral" scattering matrix, $S$



Subscripts: + : RCP wave, - : LCP wave  
 a, b, c, d: amplitudes of incoming and outgoing waves

$$S \equiv 4 \times 4$$

$$\begin{pmatrix} \text{Left -} \\ \text{Right +} \\ \text{Left +} \\ \text{Right -} \end{pmatrix} = S \begin{pmatrix} \text{Left +} \\ \text{Right -} \\ \text{Left -} \\ \text{Right +} \end{pmatrix}$$

$$S = \begin{pmatrix} r^{(L)} & t_{--} & 0 & 0 \\ t_{++} & r^{(R)} & 0 & 0 \\ 0 & 0 & r^{(L)} & t_{++} \\ 0 & 0 & t_{--} & r^{(R)} \end{pmatrix}$$

Superscripts: (L)/(R): Left/Right incidence

**Generalized unitarity relation for  $S$**

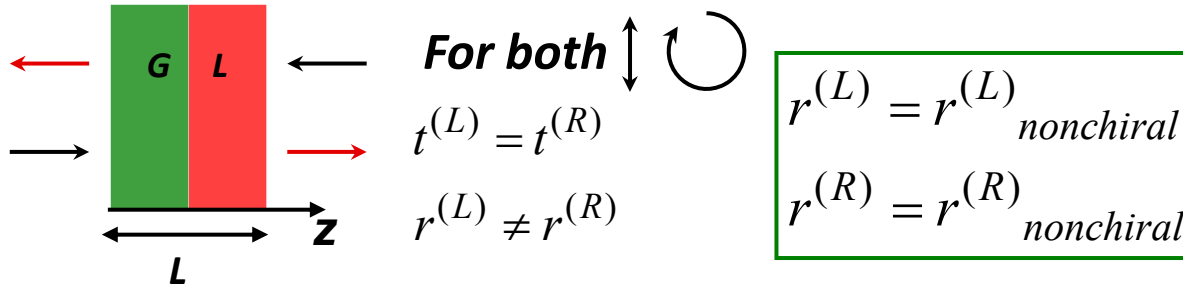
$$PTS(\omega^*)PT = S^{-1}(\omega)$$

L. Ge et al, Phys. Rev. A 85, 023802 (2012)

See Droulias et. al., Phys. Rev. Lett 122, 213201 (2019)  
 Katsantonis et. al., Photonics 7, 43 (2020)

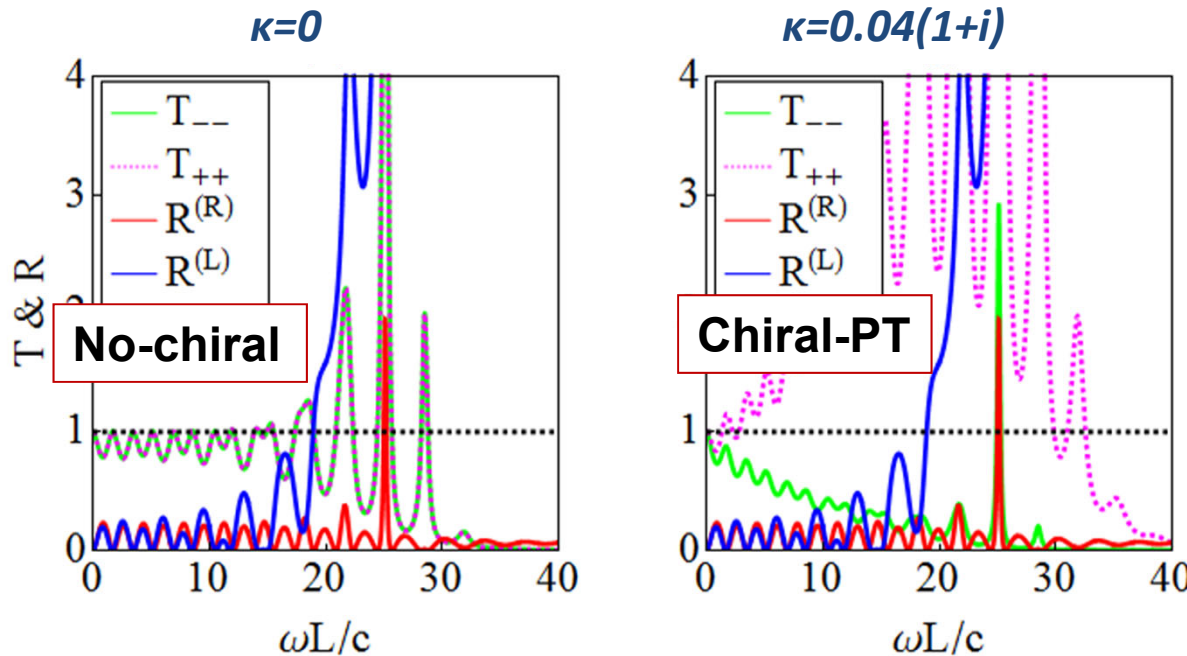
# Transmission – Reflection for normal incidence

Analytical calculations for a general chiral bi-layer



Reflection,  $r$ , independent of chirality,  $\kappa$

Superscripts (R) and (L) refer to right- and left-incidence



Transmission ( $T=|t|^2$ ) resonance-position independent of chirality

$$t_{++}t_{--} = t^2_{\text{nonchiral}}$$

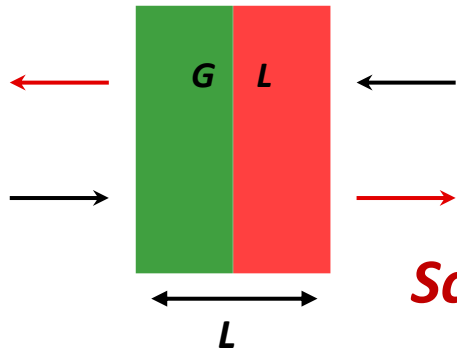
$T=|t|^2 \hat{=}$  transmission,  $R=|r|^2$  = reflection

$$\varepsilon = 3 \pm 0.4i, \mu = 1.1 \pm 0.1i$$



# S-matrix (4x4) features for normal incidence

Scattering matrix eigenvalues,  $\sigma$  (two degenerate pairs)



$$\sigma_{1,2} = \frac{1}{2} (r^{(L)} + r^{(R)} \pm \sqrt{(r^{(L)} - r^{(R)})^2 + 4t_{++}t_{--}})$$

Scattering matrix eigenvalues independent of chirality

Exceptional point, PT-phases, CPA-laser points **independent of chirality!!**

Conservation relation

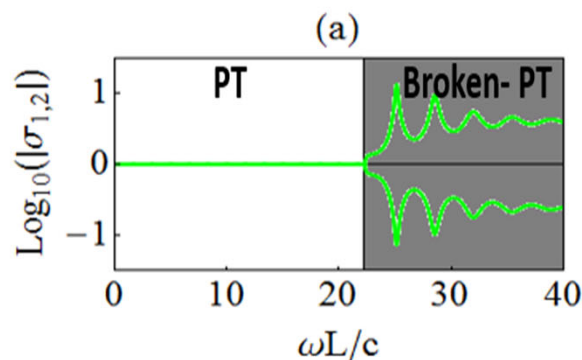
$$|\sqrt{T_{++}T_{--}} - 1| = \sqrt{R^{(L)}R^{(R)}}$$

$$T = |t|^2, R = |r|^2$$

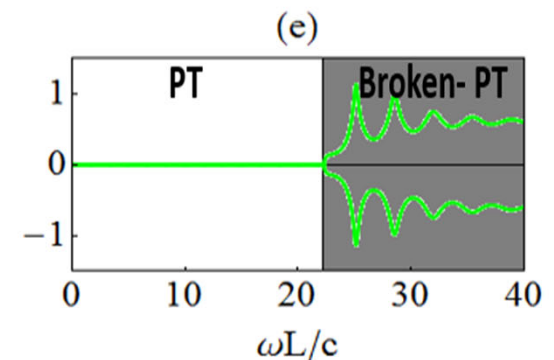
Exceptional point location

$$\frac{R^{(L)} + R^{(R)}}{2} - \sqrt{T_{++}T_{--}} = 1$$

Non-Chiral PT-system

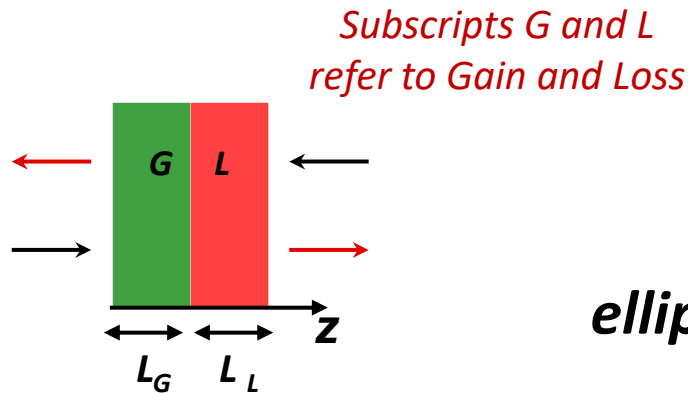


Chiral PT-system



$$\varepsilon = 3 \pm 0.4i, \mu = 1.1 \pm 0.1i, \kappa = \pm 0.04 - 0.04i$$

# Optical activity and ellipticity – normal incidence



**optical activity**

For a general chiral bi-layer

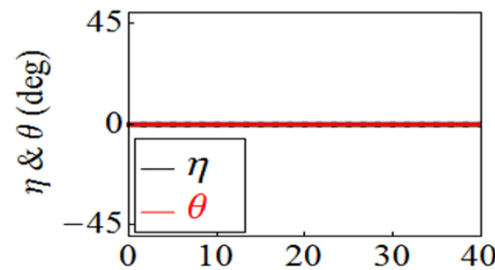
$$\theta = \frac{1}{2} k_0 (L_G \operatorname{Re}(\kappa_G) + L_L \operatorname{Re}(\kappa_L)) = 0$$

**ellipticity**

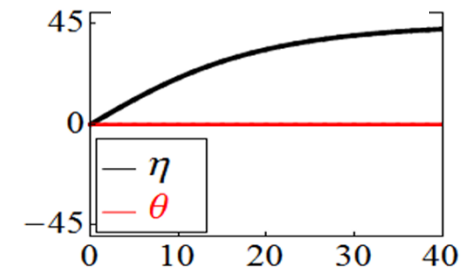
$$\eta = \frac{1}{2} \tan^{-1} \left( \frac{1 - \exp \left[ 2k_0 (L_G \operatorname{Im}(\kappa_G) + L_L \operatorname{Im}(\kappa_L)) \right]}{1 + \exp \left[ 2k_0 (L_G \operatorname{Im}(\kappa_G) + L_L \operatorname{Im}(\kappa_L)) \right]} \right)$$

**Optical activity and ellipticity depend exclusively on chirality,  $\kappa$**

**Non-chiral**



**Chiral-PT**

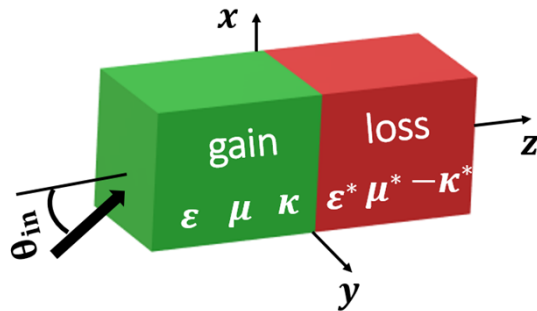


**Possibility for independent control of PT properties (exceptional point, reflection/transmission resonances, CPA-laser points) and transmitted wave polarization ( $\theta, \eta$ )!**

**E.g. CPA-laser for circularly polarized waves!**

$$\varepsilon = 3 \pm 0.4i, \mu = 1.1 \pm 0.1i, \kappa = \pm 0.04 - 0.04i$$

# Off-normal incidence: Mixed-PT phase

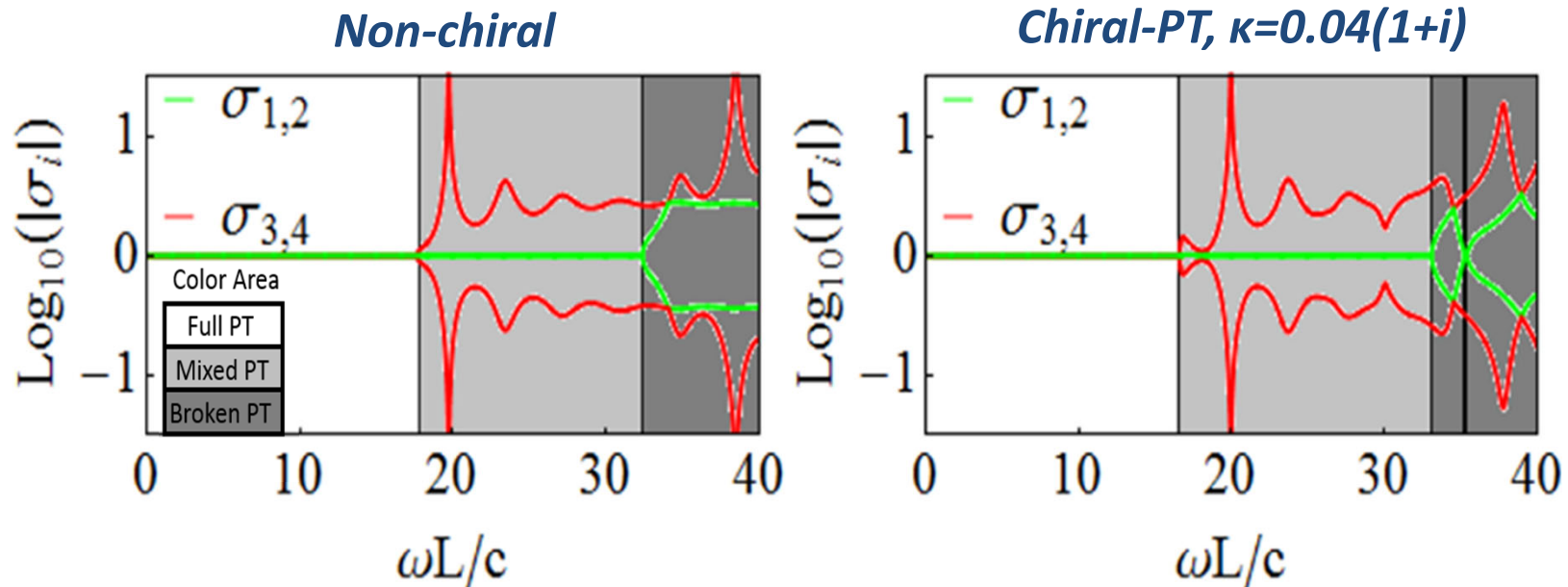


**Eigenvalues depend on chirality,  $\kappa$**

**Eigenvalues don't form two degenerate pairs**

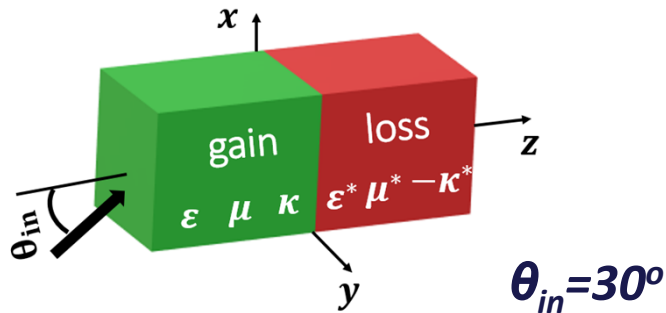
**Appearance of a mixed-PT phase**

**Scattering matrix eigenvalues,  $\sigma$**

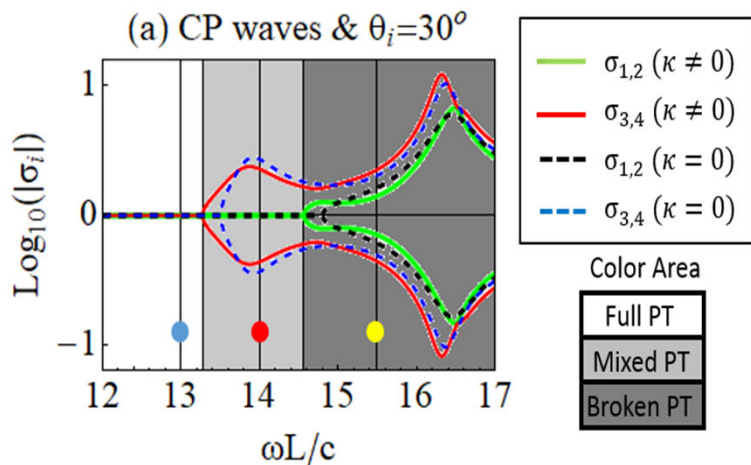


$$\epsilon = 3 \pm 0.4i, \mu = 1.1 \pm 0.1i$$

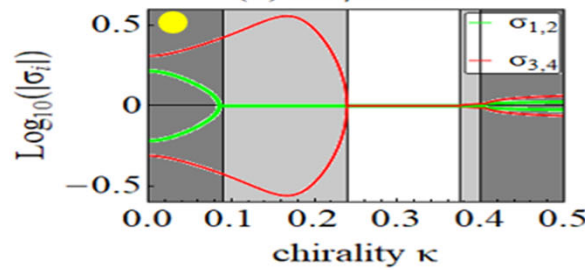
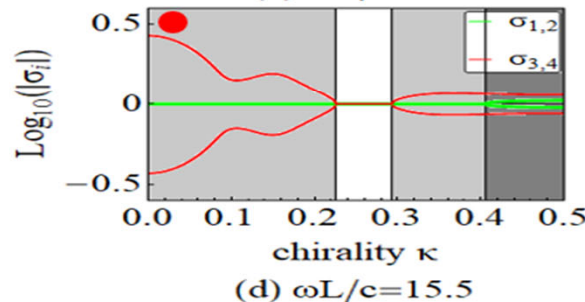
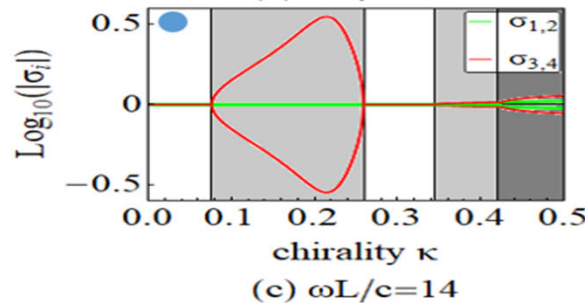
# Tuning the different phases and the exceptional points



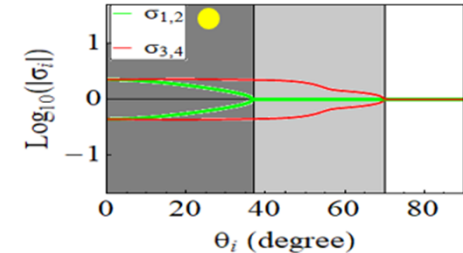
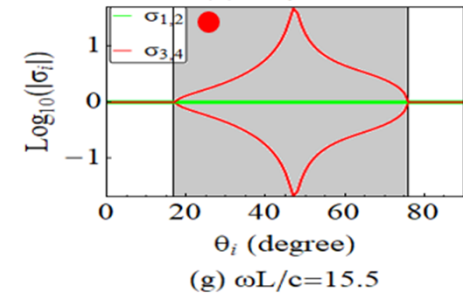
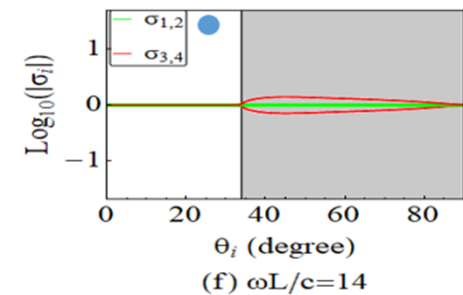
## Scattering matrix eigenvalues, $\sigma$



## Tuning through $\kappa$

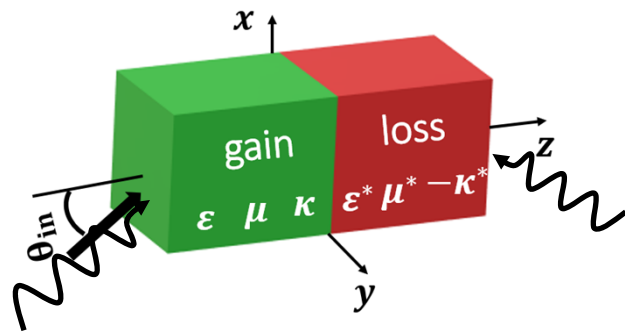


## Tuning through $\theta_{in}$

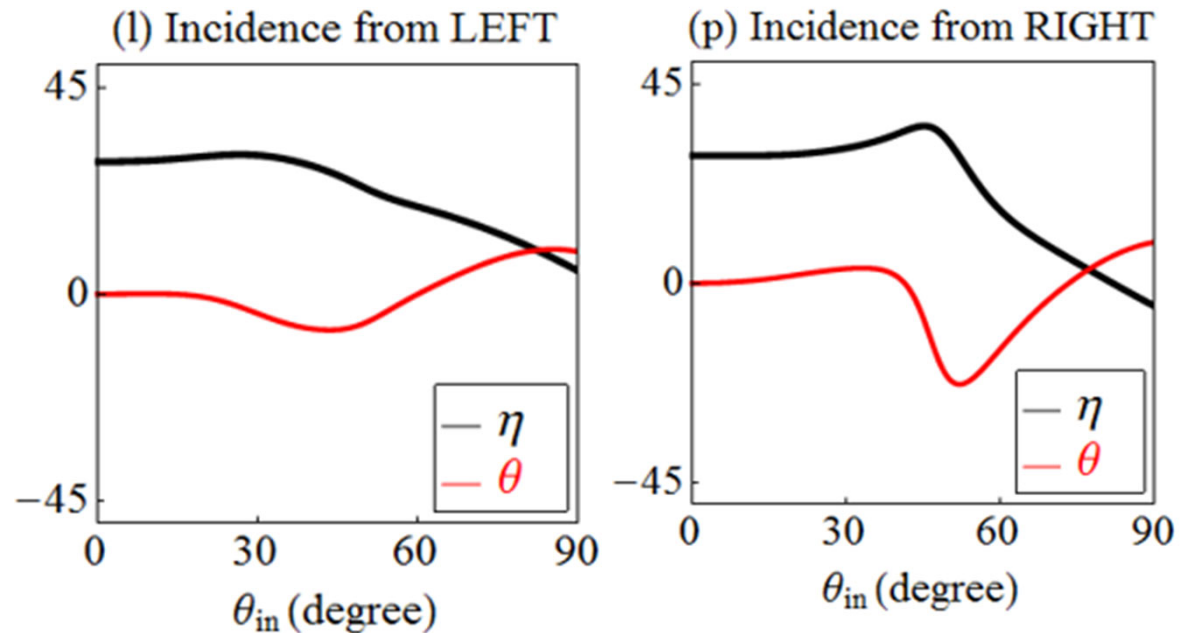


**Control of PT-features (exceptional points) through chirality and incidence angle**

# Angle-dependent and asymmetric chiral effects



## Optical activity ( $\theta$ ) - Ellipticity ( $\eta$ )



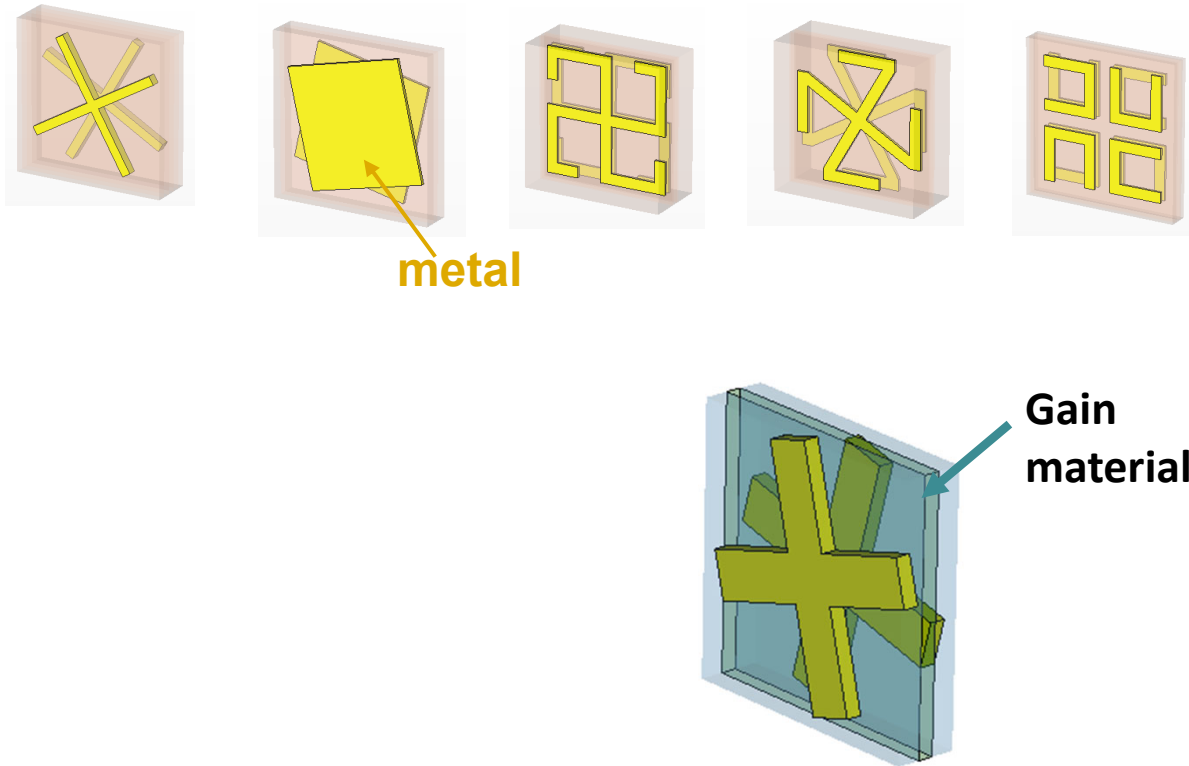
**Angle-dependent  
optical activity and  
ellipticity**

**Asymmetric optical  
activity and ellipticity**

$$\epsilon = 3 \pm 0.4i, \mu = 1.1 \pm 0.1i, \kappa = \pm 0.04 - 0.04i$$

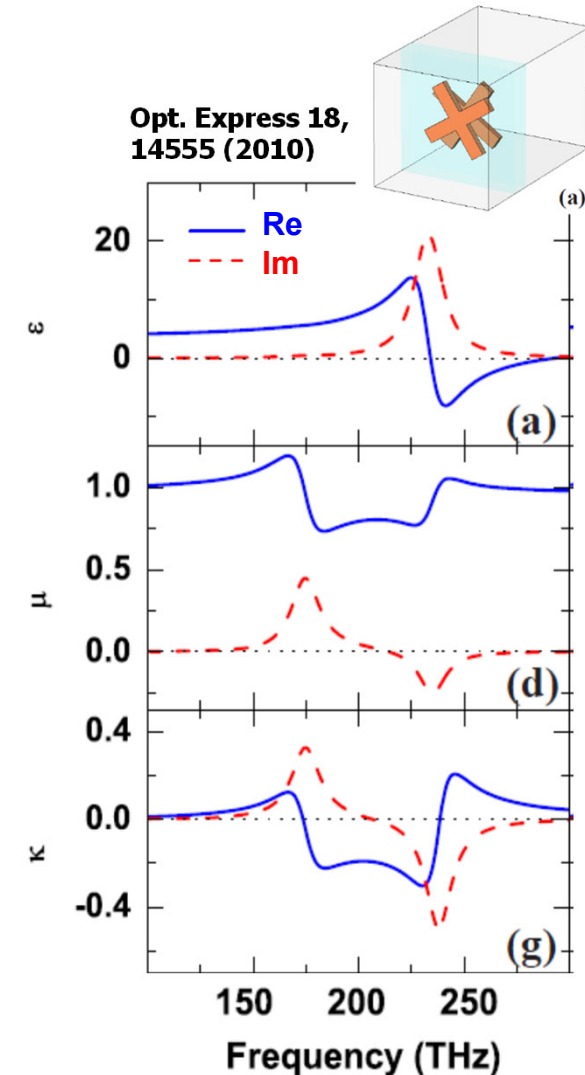
# PT symmetric chiral metamaterial practical realization?

## Bilayer-metal structures (unit cells)



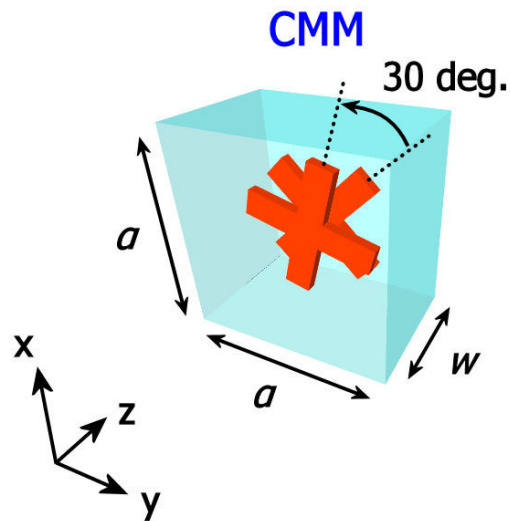
### Advantages

- Fabrication by planar technologies
- Enhancement of effective loss and gain
- “Independent” control of  $\kappa$  and  $\epsilon$ ,  $\mu$



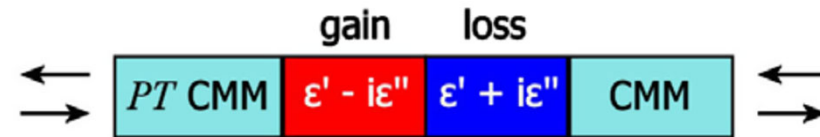
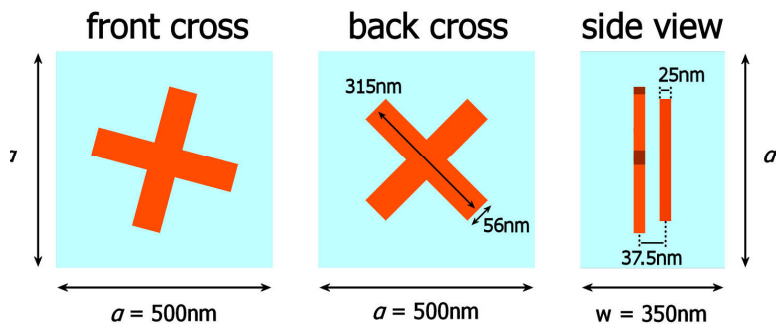
# PT-phase transition in a realistic system

**Gold crosses**  $n_{host} = 1.41 - i0.09$

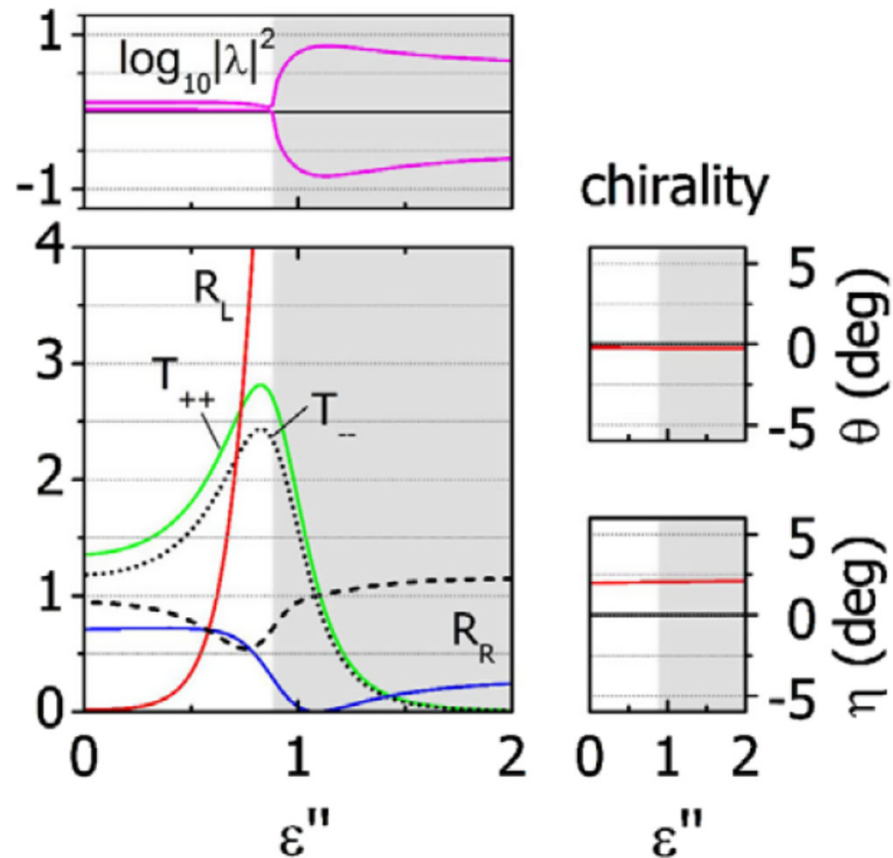


**PT @ 220 THz**

$a = 500\text{nm}$   
 $w = 350\text{nm}$



$$K_{PT-CMM} = -K_{CMM}^*$$



# Summarizing

## Combining PT-symmetry with chirality:

***Normal incidence*** → possibility of independent control of PT-effects and polarization

***Off-normal incidence*** →

- mixed PT-phases
- multiple exceptional points
- angle-dependent and chirality-dependent PT and polarization features
- asymmetric (side-dependent) polarization features

***For more info see***

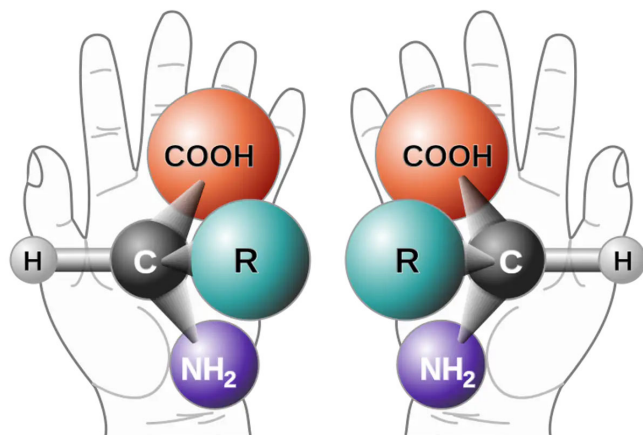
*Droulias et. al., Phys. Rev. Lett* **122**, 213201 (2019)

*Katsantonis et. al., Phy. Rev. B* **101**, 214109 (2020)

*Katsantonis et. al., Photonics* **7**, 43 (2020)



# PT symmetry for molecular chirality sensing



Left-handed  
enantiomer (S -  
sinister)

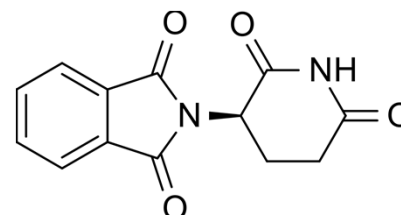
Right-handed  
enantiomer (R -  
rectus)

$$\kappa \sim 10^{-4}$$

*Why?*

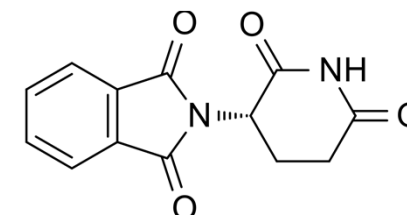
*Opposite enantiomers: Different  
biological activity*

Morning sickness  
treatment



(R)-thalidomide

Teratogenesis



(S)-thalidomide

*Enantiomer detection and discrimination is of great importance for  
biology and pharmaceutical industry*

# Chirality effects used for sensing

$$n_{\pm} = \sqrt{\epsilon\mu \pm \kappa}$$

## **Optical activity**

Rotation of linear polarization

## **Circular dichroism (CD)**

Different absorption for RCP(+)/LCP(-) waves: **CD:  $A_+ - A_-$**

*measured by:*

$$CD = A_+ - A_- = |t_{--}|^2 - |t_{++}|^2$$

↑  
*transmittance of LCP (-) waves*

↑  
*transmittance of RCP (+) waves*

***In natural chiral media  $\kappa \sim 10^{-4} \rightarrow$  weak chiro-optical signals, e.g.  $CD \sim 10^{-5}$***

# Chirality effects used for sensing

$$n_{\pm} = \sqrt{\epsilon\mu} \pm \kappa$$

**Kuhn's dissymmetry factor,  $g$**   
Relative differential absorption  
for RCP(+)/LCP(-) waves

$$g = (A_+ - A_-) / [(A_+ + A_-) / 2]$$

(percentage of enantioselectivity in  
photoexcitations)

$$g = \frac{A_+ - A_-}{(A_+ + A_-) / 2} = \frac{CD}{(A_+ + A_-) / 2}$$

**Circular dichroism (CD)**  
Different absorption for  
RCP(+)/LCP(-) waves: **CD:  $A_+ - A_-$**

measured by:

$$CD = A_+ - A_- = |t_{--}|^2 - |t_{++}|^2$$

**Enhancing  $g$  →**  
Enhanced enantioselectivity  
in **photoionization, photolysis,**  
**fluorescence,**  
**photopolymerization, etc.**

**In natural chiral media  $\kappa \sim 10^{-4}$  → weak chiro-optical signals, e.g.  $CD \sim 10^{-5}$**

# Absorption in chiral media

*Absorbed power density* of a CP (+/-) wave by a thin chiral layer of parameters  $\epsilon, \mu, \kappa$  (Poynting's th.

$$A_{\pm} \sim \frac{\omega}{2} [\text{Im}(\epsilon) |\mathbf{E}|^2 + \text{Im}(\mu) |\mathbf{H}|^2] \pm \frac{2\omega}{c} \text{Im}(\kappa) \text{Im}(\mathbf{E} \cdot \mathbf{H}^*)$$

RCP 

LCP 

Chiral  
  
 $d_c$

$$CD = A_+ - A_- = 4c \text{Im}(\kappa) \left[ \frac{\omega}{2c^2} \text{Im}(\mathbf{E} \cdot \mathbf{H}^*) \right]$$

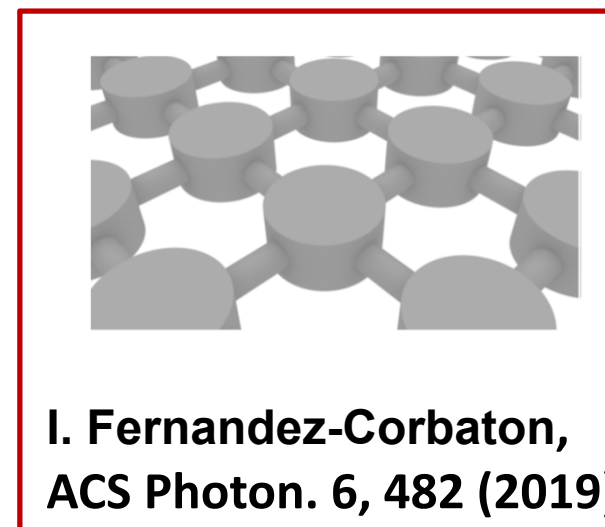
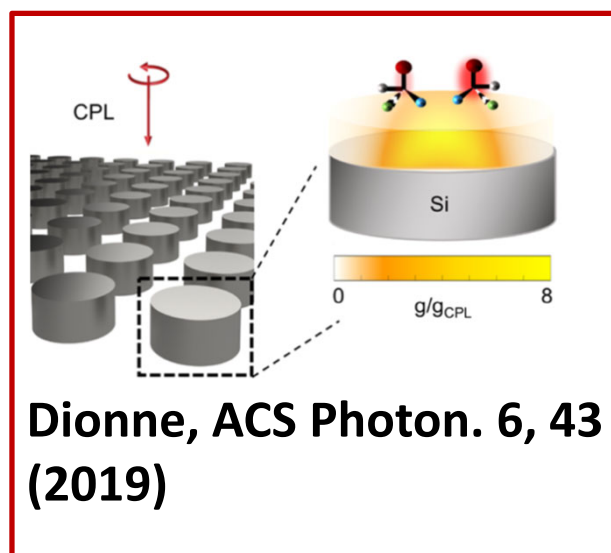
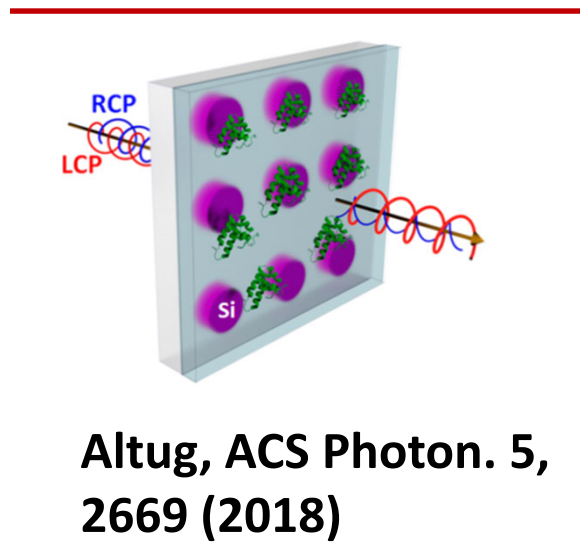
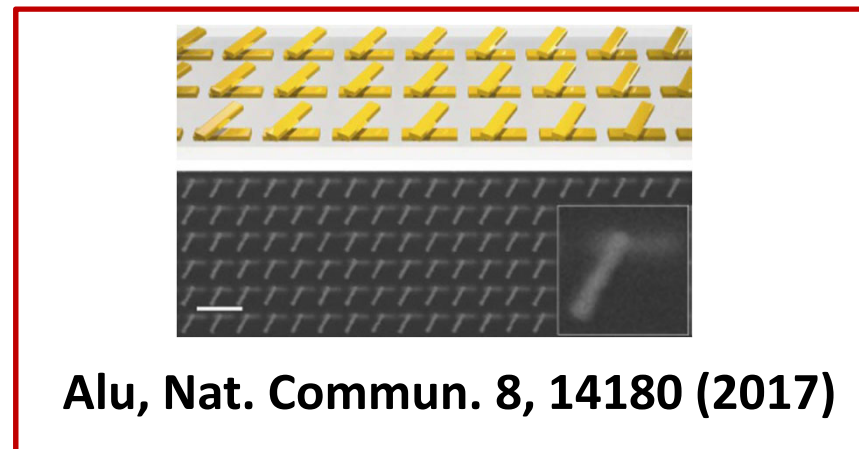
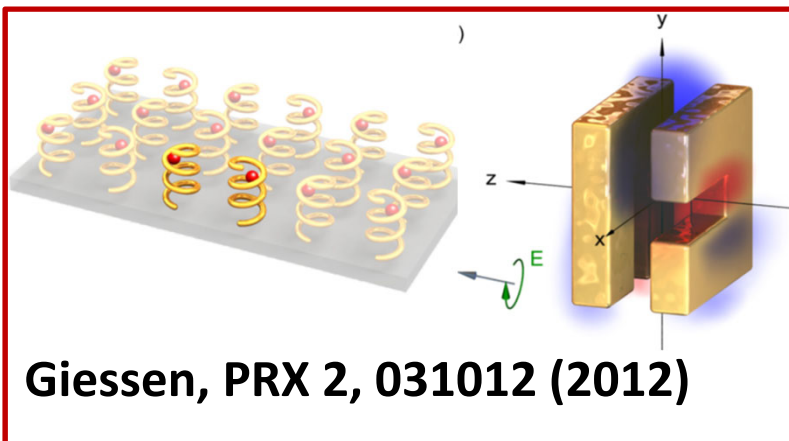
*Field chirality*

$$C = -\frac{\omega}{2c^2} \text{Im}(\mathbf{E} \cdot \mathbf{H}^*)$$

*Enhancement through proper resonances – enhanced local fields*

Tang & Cohen, PRL 104, 163901 (2010)

# Nanophotonics for enhancing circular dichroism



***Superiority of achiral structures – Overlapping resonances – up to 150x CD enhancement***

# Absorption in chiral media

Absorbed power density of a CP wave by a thin chiral layer of parameters  $\epsilon, \mu, \kappa$

$$A_{\pm} \sim \frac{\omega}{2} [\text{Im}(\epsilon) |\mathbf{E}|^2 + \text{Im}(\mu) |\mathbf{H}|^2] \pm \frac{2\omega}{c} \text{Im}(\kappa) \text{Im}(\mathbf{E} \cdot \mathbf{H}^*)$$

RCP   
LCP 

Chiral  
 $\updownarrow$   
 $d_c$

$$CD = A_+ - A_- = 4c \text{Im}(\kappa) \left[ \frac{\omega}{2c^2} \text{Im}(\mathbf{E} \cdot \mathbf{H}^*) \right]$$

Field chirality

$$C = -\frac{\omega}{2c^2} \text{Im}(\mathbf{E} \cdot \mathbf{H}^*)$$

Enhancement through proper resonances – enhanced local fields

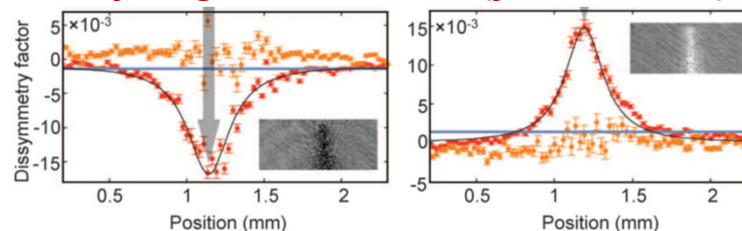
Tang & Cohen, PRL 104, 163901 (2010)

$$g_{\pm} = \frac{2(A_+ - A_-)}{A_+ + A_-} \propto \frac{8c \text{Im}(\kappa) C}{\omega [\text{Im}(\epsilon) |E|^2 + \text{Im}(\mu) |H|^2]}$$

Nodes of standing wave in specially designed cavity

Tang & Cohen, PRL104, 163901 (2010)

11-fold g-enhancement (fluorescence)



Tang & Cohen, Science 332, 333 (2011)

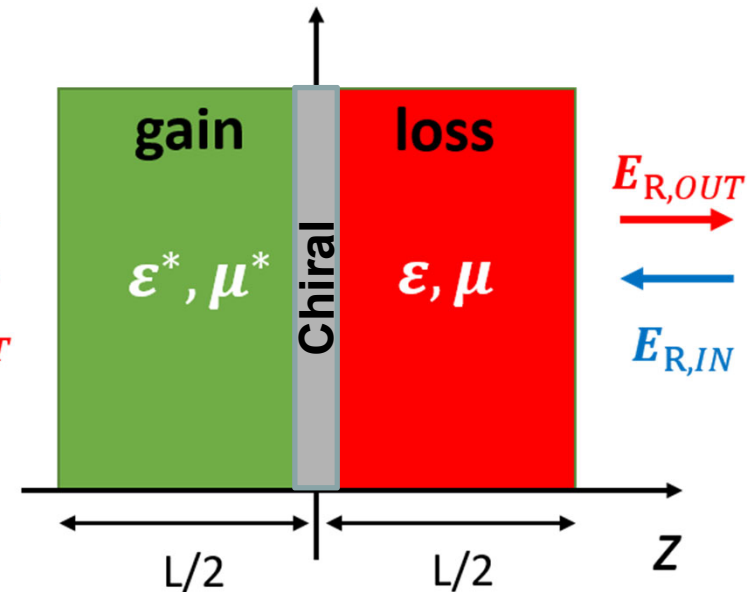
He et al., Nature Comm. (2018)

(Photopolymerization dissymmetry)

# Our approach: Combine loss and gain media

## Parity-Time (PT) symmetric bilayer

$$g \propto \frac{\frac{1}{V} \int \text{Im}(\kappa) \text{Im}(\mathbf{E} \cdot \mathbf{H}^*) dV}{\frac{1}{V} \int [\text{Im}(\epsilon) |\mathbf{E}|^2 + \text{Im}(\mu) |\mathbf{H}|^2] dV}$$



## Why Parity-Time (PT) symmetric?

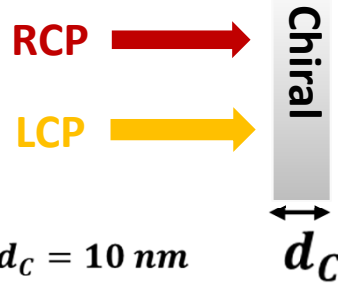
- Balanced loss and gain  $\text{Im}(\epsilon_{\text{gain}}) = -\text{Im}(\epsilon_{\text{loss}})$
- Strong fields at gain-loss interface
- **Peculiar**, potentially useful **electromagnetic effects**  
(unidirectional invisibility, simultaneous coherent perfect absorption & lasing)

$$\begin{aligned} \epsilon(z) &= \epsilon^*(-z) \\ \mu(z) &= \mu^*(-z) \end{aligned}$$

I. Katsantonis et. al., Phys. Rev. B 105, 174112 (2022)

# CD enhancement in the PT-symmetric 3-layer system

Chiral layer alone

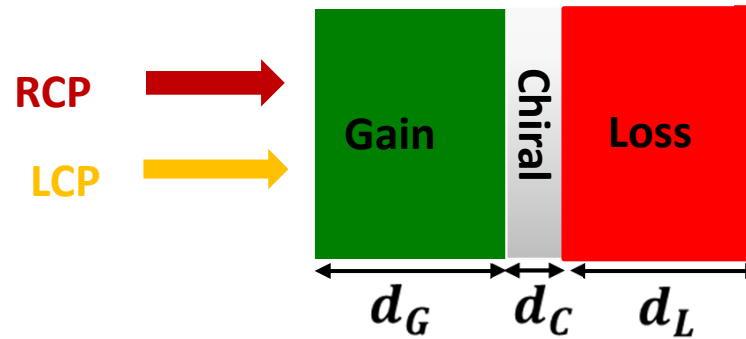


$$d_c = 10 \text{ nm}$$

$$n_c = 1.45 + 0.01i$$

$$\kappa = \pm 5 \times 10^{-5}(1 + i)$$

PT-symmetric trilayer

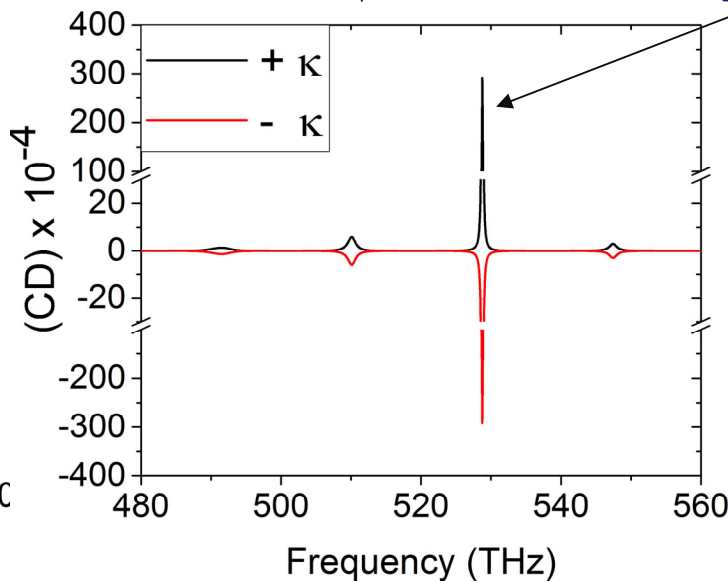
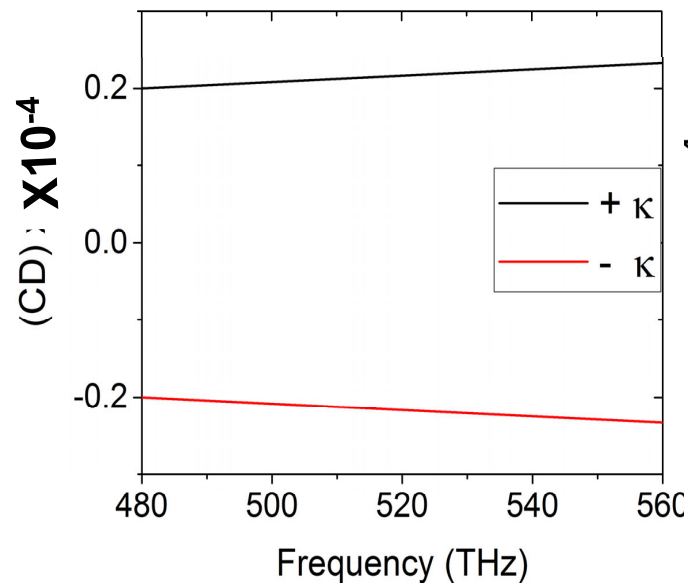


$$d_{G/L} = 2.5 \mu\text{m}$$

$$n_{G/L} = 3.2 -/+0.05i$$

$$CD = A_+ - A_- = |t_-|^2 - |t_+|^2$$

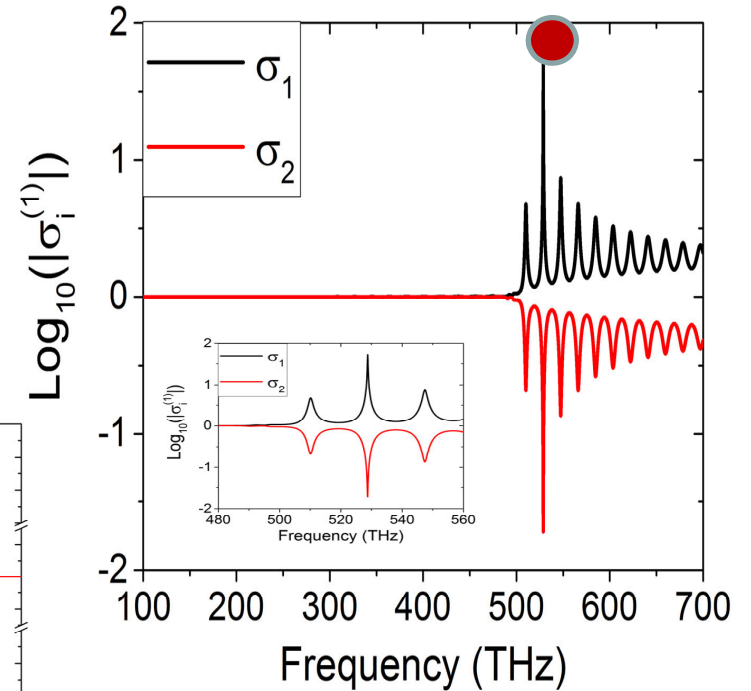
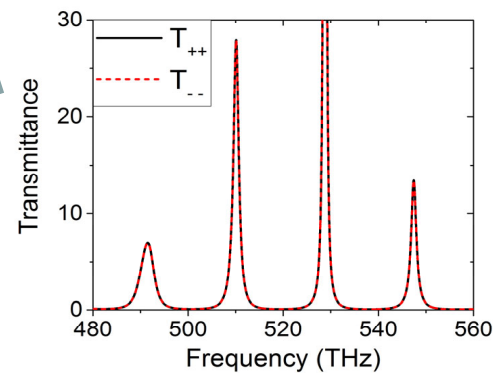
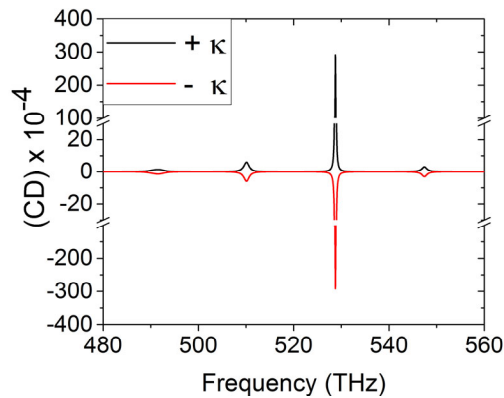
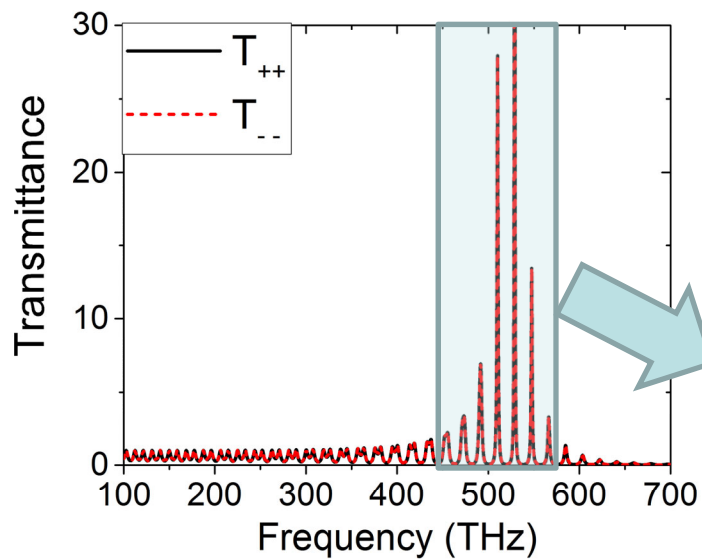
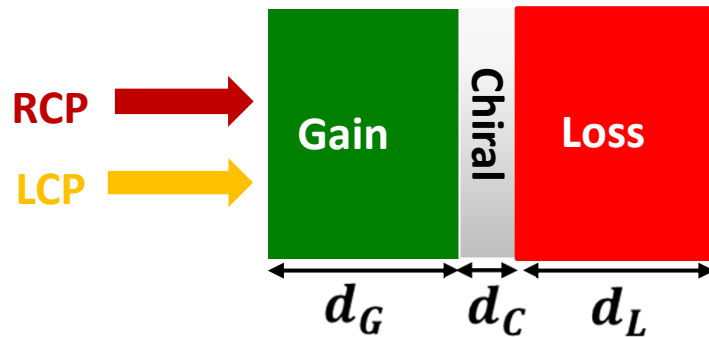
**~1500 times CD enhancement!**



**Many CD peaks of ~20x enhancement**

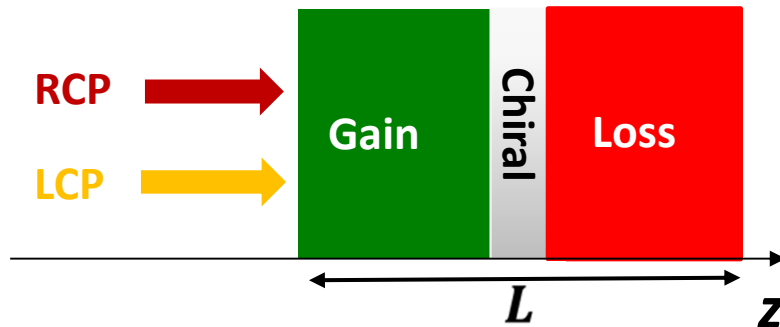


# Transmission / reflection / scattering matrix for the 3-layer PT system

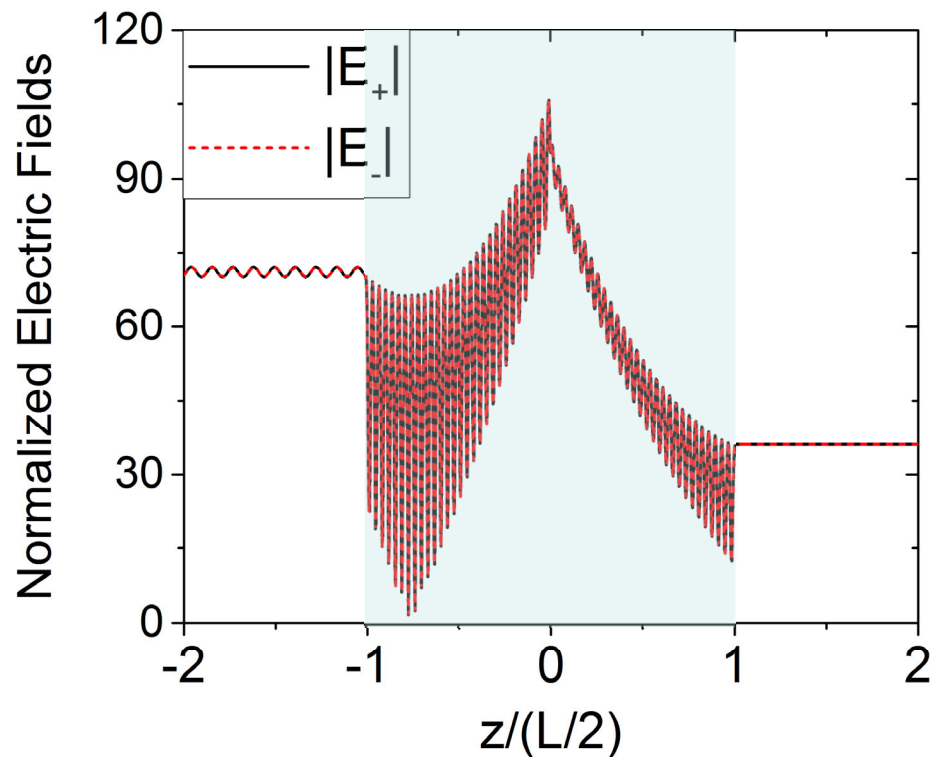


**Largest CD enhancement at lasing threshold**

# Electric field at the CD peak of 3-layer PT system



*Larger field in gain region*



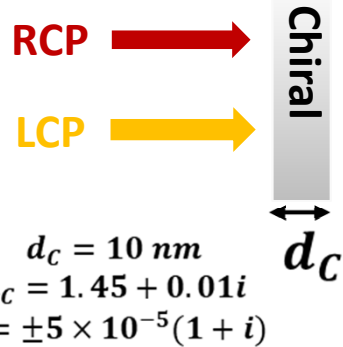
*Large interfacial field intensities*



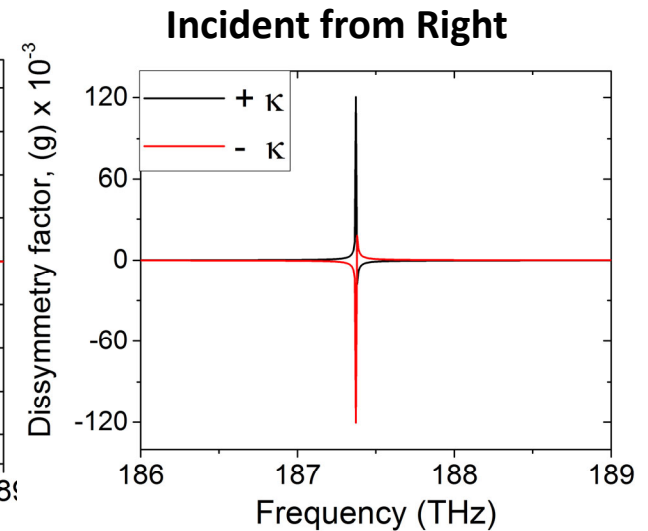
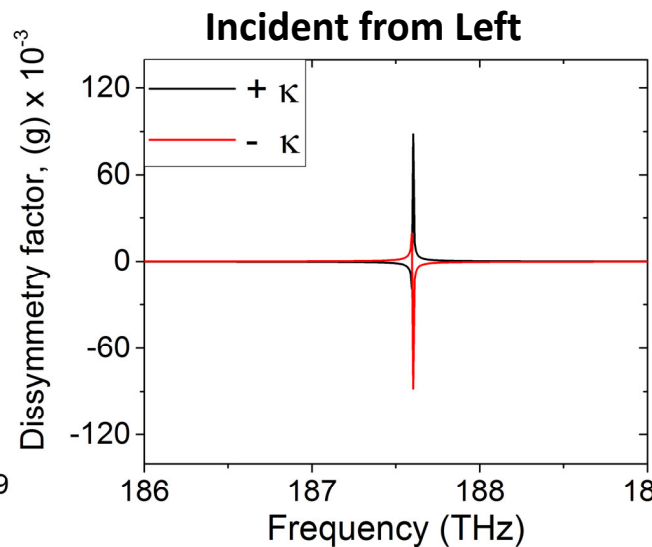
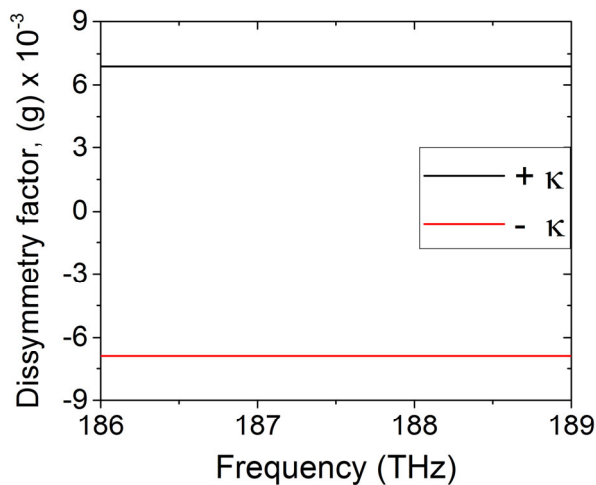
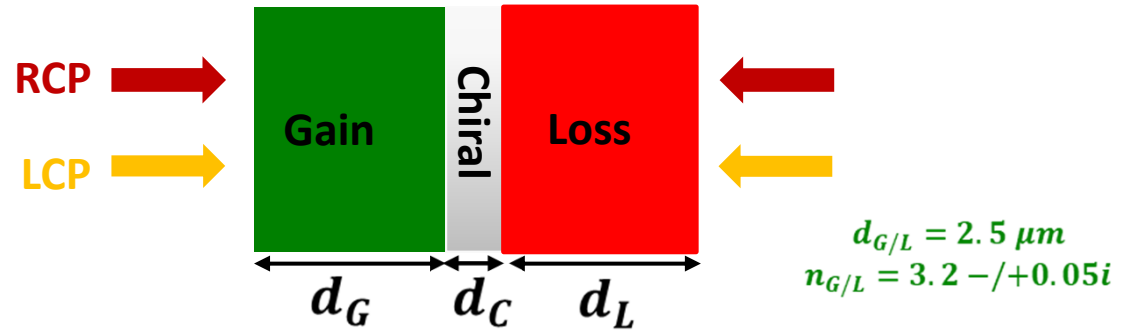
*Large optical chirality*

# $g$ -enhancement in the PT-symmetric 3-layer system

Chiral layer alone



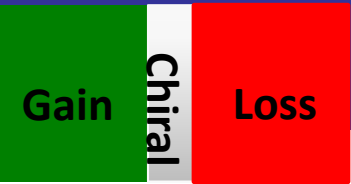
PT-symmetric trilayer



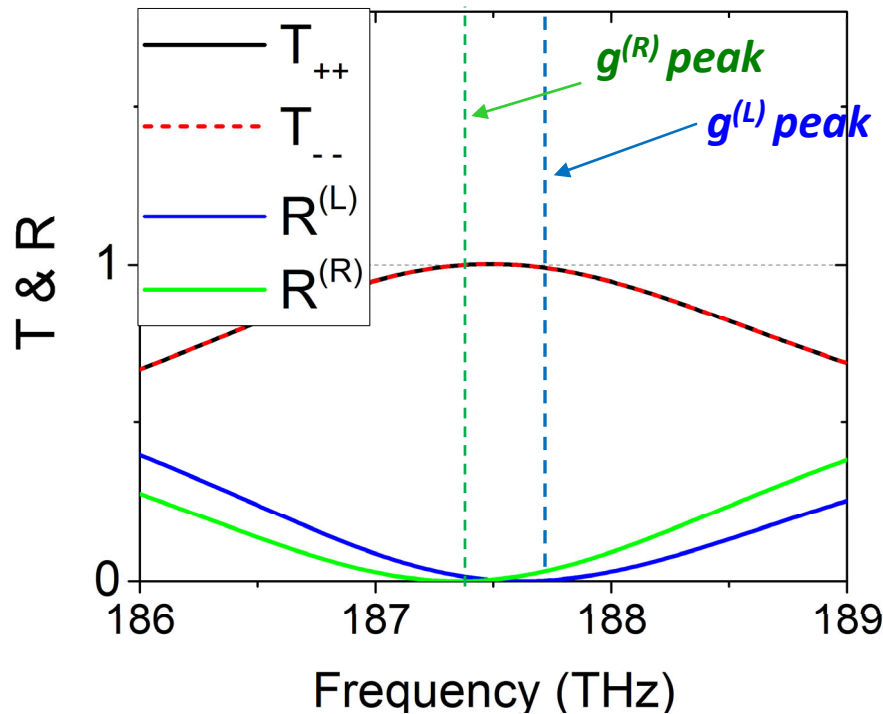
$$g^{(L/R)} = \frac{2(|t_-|^2 - |t_+|^2)}{2 - |t_-|^2 - |t_+|^2 - 2|r^{(L/R)}|^2}$$

***~15x dissymmetry enhancement compared to chiral layer alone***

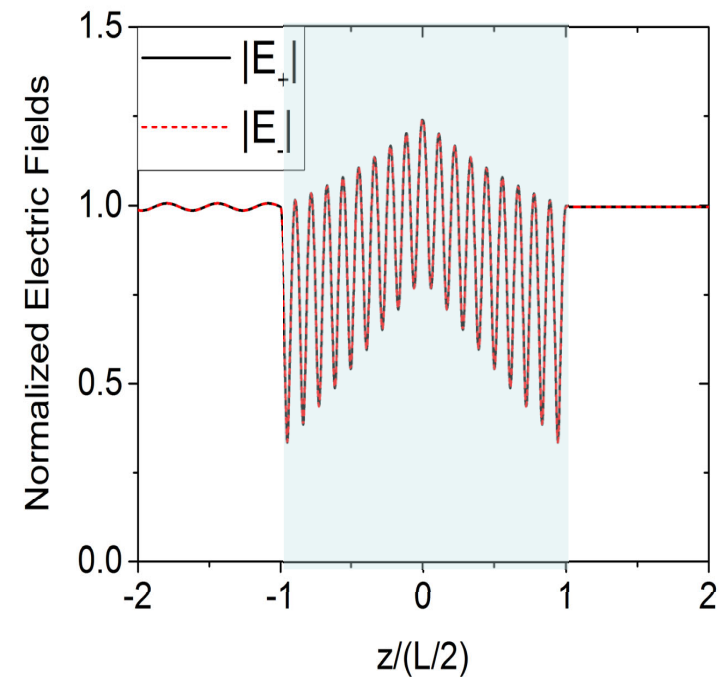
# $g$ -enhancement investigation for the PT-symmetric 3-layer system



$g$  enhancements @  $R=0, T\approx 1$



*Field intensity along system @  $g$ -peak*



*$g$  enhancements centered at flux conserving points with  $R=0, T=1, A=0$  - anisotropic transmission resonances*

*$g$ -enhancement associated with appreciable field intensity*

# Conclusions/summary

**PT-symmetric media** provide a **promising avenue** for **molecular chirality sensing** and **enantioselective field-molecule interactions**

*In such systems:*

Large **Circular Dichroism** enhancements

Large **dissymmetry factor ( $g$ )** enhancements with appreciable field intensities

**For details: I. Katsantonis et. al.,  
Phys. Rev. B 105, 174112 (2022)**

*Thank you!*

nanoPOLY

visor  
SURF

**Ultrachiral**