Ultrafast X-rays: What are they good for?

Phil Bucksbaum Siegman School 2016





Atomic, Molecular, and Optical Science (AMO)





X-ray interactions with atoms





Auger spectrum from N₂







Transient state studies: N₂⁺⁺ Potential Energy Surfaces





R. W. Wetmore and R. K. Boyd, J. Phys. Chem. 90, 5540 (1986).

Auger electron energies observed in the molecular frame from $N_2 \rightarrow N_2^+ \rightarrow N_2^{2+}$ at 1.1keV





⁽Cryan et al, J. Phys. B 45 055601 (2012))

Strong fields: 1-photon, 1-electron ionization



consider a 1-photon K-shell transition: $\sigma_{K} \approx 10^{-18} \text{ cm}^{2}$ $\Gamma_{K} = \sigma_{K}F_{IcIs} \approx 10^{15} \text{ s}^{-1}$ (saturated) $t_{K} = 1/\Gamma_{K} = 1 \text{ fs}$ 2s,p 1s photoionization

• rapid enough to ionize more than one electron.

• fast enough to compete with atomic relaxation.



Ne¹⁰⁺ at the focus:





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Strong and Ultrafast: Hollow atoms





Double core holes form when the photoionization rate exceeds Auger



Time

Hollow atom formation at LCLS





High Energy Auger Spectrum of N₂ from LCLS shows clearer evidence for double core hole formation





Ultrafast electron dynamics: Nucleobase photoprotection





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Ultrafast electron dynamics: Nucleobase photoprotection





Non-Born-Oppenheimer dynamics





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Auger energy is sensitive to valence charge near the oxygen atom.



Localized structural evolution: Time-resolved Auger Electron Spectroscopy





Geometry change: 0.15 Å Spectral change: several eV



McFarland et al. Nature Commun. 2014, 5, 4235

Probing ultrafast $\pi\pi^*$ - $n\pi^*$ transitions via oxygen K-edge resonant absorption:





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NEXAFS spectrum indicates a rapid decay to the $n\pi^*$ state





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Hard x-rays at X-FELs: Single particle imaging paradigm





- Ultrashort pulses outrun damage: diffract before destroy
- High intensities: A diffraction image on every spot.
- High repetition: Millions of images.

Serial femtosecond nanocrystallography





Chapman et al., Nature 470, 73 (2011)

PS-II structure from LCLS





Chapman, et al. Nature 470, 73 (2011)



CS-PAD (2.5Mpx) detector





Project each $S(Q_r,t)$ onto first 10 Legendre polynomials $P_L(\cos\theta)$





Signal comes from the entire charge distribution





Extracting the molecular movie from the scattering data.





 $\mathcal{AC}[\rho(\vec{x},t)] \equiv \rho(\vec{x},t) \otimes \rho(\vec{x},t)$ $= \mathcal{F}^{-1}(f(\vec{Q},t)f^*(\vec{Q},t))$

$$\mathcal{AC}[\rho(\vec{x},t)] = \mathcal{AC}[\rho_X(\vec{x})] + \mathcal{AC}[\rho_B(\vec{x},t)] + 2\mathcal{CC}[\rho_X(\vec{x}),\rho_B(\vec{x},t)]$$

$$2 \ \mathcal{CC}[\rho_X(\vec{x}), \rho_B(\vec{x}, t)] \simeq \mathcal{AC}[\rho(\vec{x}, t)] - \mathcal{AC}[\rho_X(\vec{x})]$$

The cross correlation is a hologram where the ground charge distribution creates the reference scattered wave.

M. J. J. Vrakking and T. Elsaesser, Nat Photon 6, 645-647 (2012).



$P_2(\cos\theta)$ component of S(Q,t)







В ЗПА 2 3 R(Å) 4 5 6 7 3.5 65 6 3 5.5 2.5 5 .5 2 1.5 3.5 3 1 0.5 0 5 1.5 2 0 WW

₹П

The movie



The future: X-ray-induced attosecond electron motion





Attosecond X-ray Free Electron Lasers



Laser pre-modulated electrons tame SASE



Ding, Y., PRST 12, (2009).

Mapping electron dynamics with core excitation



Core excitation creates localized electron disturbances. Correlation drives nonlocal electron transport in molecules

Example of how this could work: Send in three x-rays, k_1 , k_2 , and k_3 and read out the final Auger electron spectrum



Lunnemann et al., Chem Phys Lett 450 232 (2008); Mukamel et al., Ann Rev. P. Chem. 64, 101 (2013); Miyabe, S. & PHB, PRL 114, 143005 (2015)

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Lots of contributors

- Bucksbaum group
 - Adi Natan
 - Song Wang
 - Julien Devin
 - Matthew Ware
 - Adrei Kamalov
 - James Cryan
 - Vlad Petrovic
 - Chelsea Liekhus-Schmaltz
 - Greg McCracken
 - Andreas Kaldun
- PULSE and LCLS collaborators
 - Todd Martinez
 - Kellly Gaffney
 - David Reis
 - Shambhu Ghimire
 - Shungo Miyabe
 - Cristoph Bostedt (75% LCLS)
 - Mariano Trigo
 - Mike Glownia
 - Hermann Durr
 - Timor Osipov
 - Thomas Wolf
 - Ryan Coffee



- Other alumni in the past three years
 - Limor Spector (went to McKinsey)
 - Brian McFarland (went to LANL)
 - Fenglin Wang (went to CFEL)
 - Joe Farrell (graduated)
 - Doug Broge (graduated)
 - Ben Barbrel (LBNL)
 - Jaehee Kim (NSF, graduates 12/14)
 - James White (NSF, startup)
- Outside collaborators (partial list)
 - Nora Berrah et al, U. Conn
 - Lou Dimauro et al, Ohio State
 - Artem Rudenko, K-State
 - Tamar Seideman, N'western
 - Linda Young et al, ANL
 - Ilya Averbukh, Weizmann
 - Jon Marangos et al, Imperial
 - Hamed Merdji, CEA
 - Roseanne Sension, UM
 - Fenglin Wang, and others, CFEL
 - Markus Guehr (U. Potsdam)

