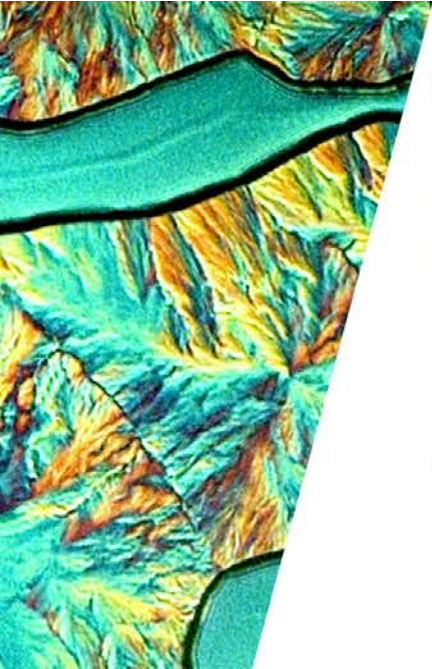


How to Enhance Nonlinear Optical Signals at the Nanoscale

Presented by:



The OSA Nonlinear Optics Technical Group Welcomes You!



WEBINAR ON HOW TO ENHANCE NONLINEAR OPTICAL SIGNALS AT THE NANOSCALE

10 January 2019 • 11:00 EST

OSA Nonlinear
Optics
Technical Group

OSA Nonlinear
Optics
Technical Group

Technical Group Leadership 2018



Chair

Ajanta Barh

DTU Fotonik, Denmark



Event officer

Ryan T. Glasser

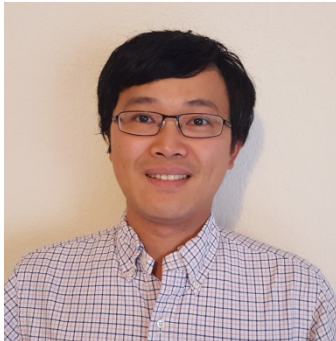
Tulane University, USA



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Vice Chair

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U. of Colorado, USA



E-magazine & Social media officer

Samuel Serna

MIT, USA



Webinar Co-officer

Eva Pogna

CNR, Pisa, Italy

Technical Group at a glance

- **Focus**

- “Physics of nonlinear optical materials, processes, devices, & applications”
- **3800** members (**largest** in OIS, 3rd largest in OSA)

- **Mission**

- To benefit YOU
- webinars, e-Presence, publications, technical events, business events, outreach
- Interested in presenting your research? Have ideas for TG events? Contact us at

- **Email:** TGNonlinearOptics@osa.org

- Find us here

- www.osa.org/NonlinearOpticsTG

- **Facebook:** www.facebook.com/osanonlinearoptics

- **LinkedIn:** www.linkedin.com/groups/8302249

Today's webinar



How to enhance non-linear optical signals at the nanoscale

Speaker's short Bio:

Graduation at Indian Institute of Technology in New Delhi, India

Ph.D. degree from ETH Zürich (CH)

Postdoc at the EPFL in Lausanne (CH)

Junior group leader at the Friedrich Schiller University, Jena (DE)

Prof. Dr. Rachel Grange

Research Group leader of Optical Nanomaterial Group

Physics Department, Institute for Quantum Electronics, ETH Zürich, Switzerland

grange@phys.ethz.ch

ETH zürich

Department of Physics | Institute for Quantum Electronics

Optical Nanomaterial Group

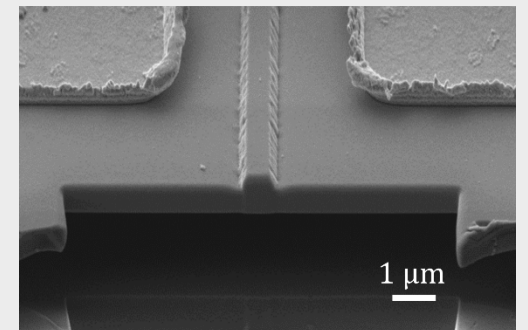
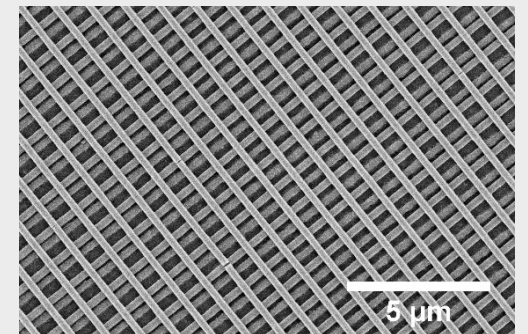
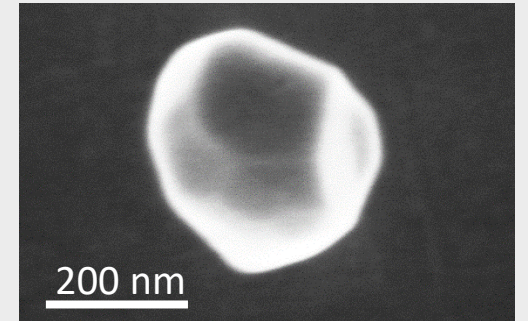
How to Enhance Nonlinear Optical Signals at the Nanoscale?

www.ong.ethz.ch

Beyond Metal and Semiconductors: Nano-Oxides for Nonlinear Photonics Webinar

Rachel Grange
Optical Nanomaterials Group
Department of Physics
ETH Zurich, Switzerland

grange@phys.ethz.ch
@rachel_grange



How to Enhance Nonlinear Optical Signals at the Nanoscale?

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Flavia Timpu



Marc Reig



Viola Vogler



Fabian Kaufmann



David Pohl



Maria Timofeeva

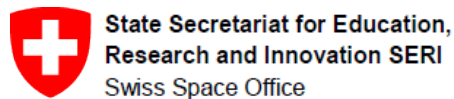


Romolo Savo

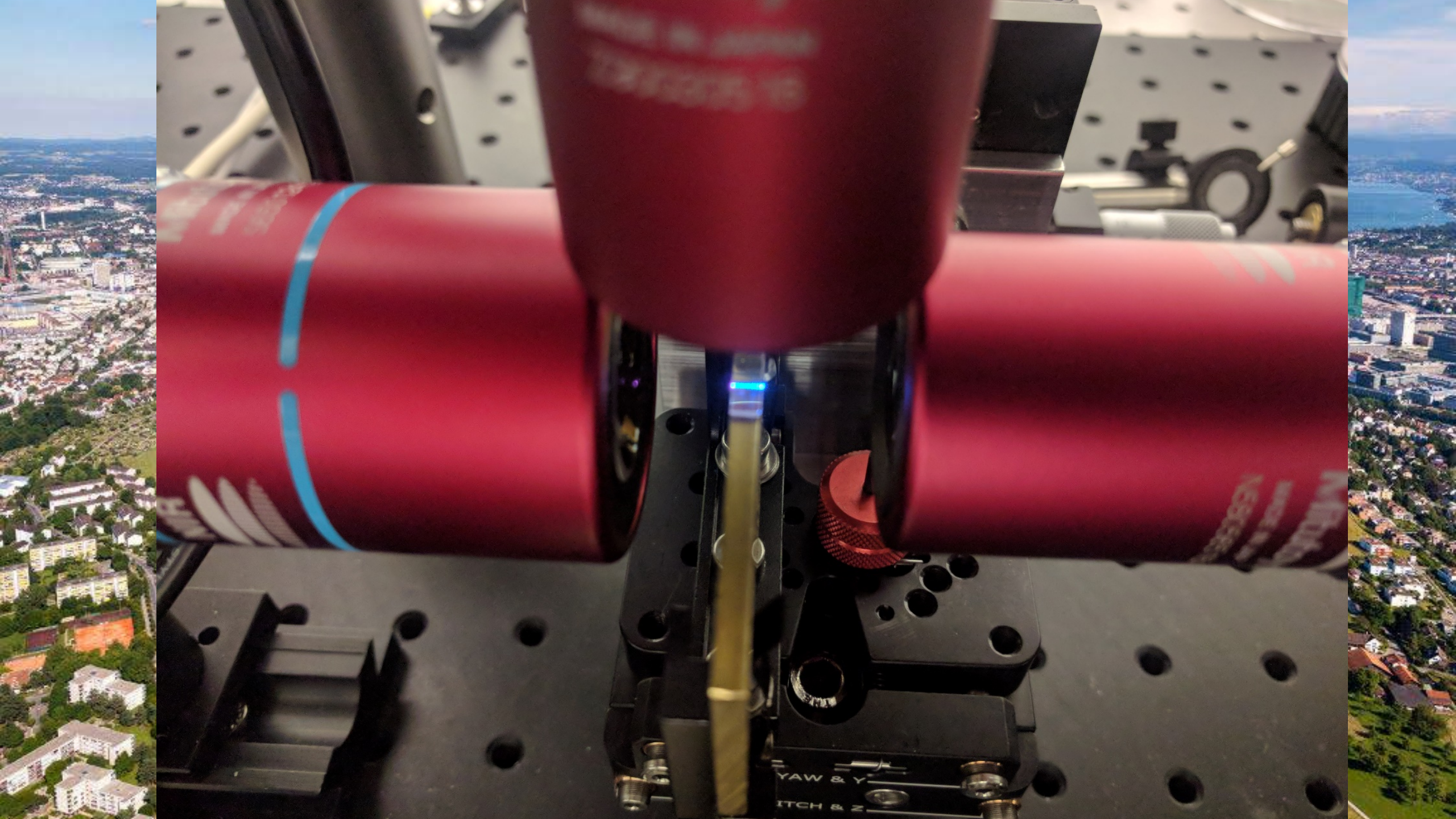
Andrea Morandi
Kevin Stalder
Damian Bucher
Iek Tang
Grégoire Saerens

Alumni: Anton Sergeyev, Nick Hendricks, Andrea Steinbrück, Eugenie Kim, Claude Renaut

Funding

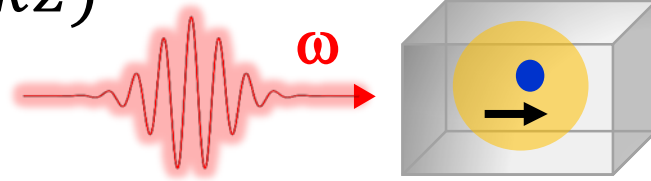






What is nonlinear optics ?

$$\vec{E}(t) = A \cos(\omega t + kz)$$

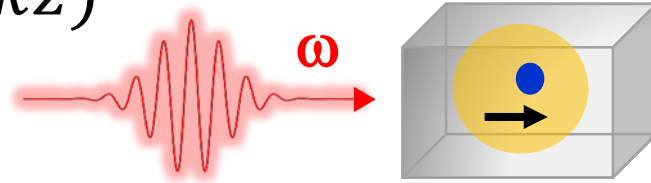


Induced dipole moment

No change under light

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No change under light

Induced dipole moment

Linear regime

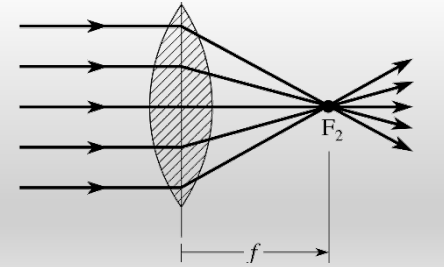
Induced Polarization

$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E}$$

dielectric constant
of the vacuum

$\chi^{(1)}$ Electric susceptibility

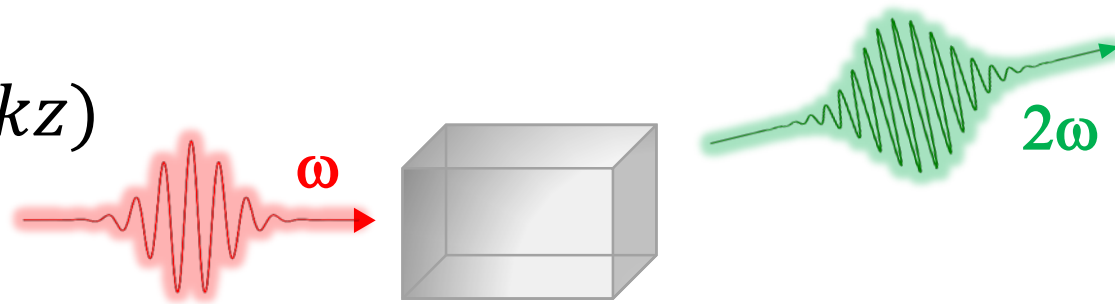
$$n^2 = 1 + \chi$$



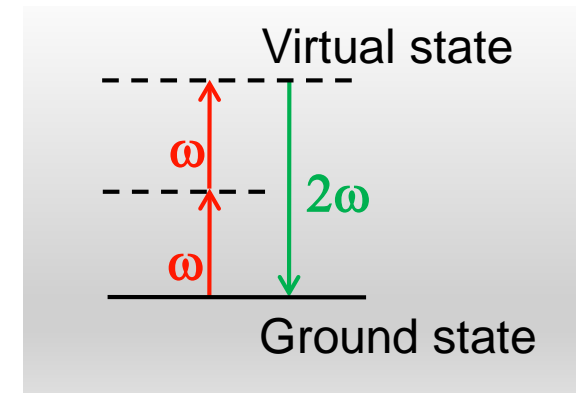
no longer valid for intense light

What is nonlinear optics ?

$$\vec{E}(t) = A \cos(\omega t + kz)$$



Second Harmonic Generation



Nonlinear regime

Induced Polarization

$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E} + \epsilon_0 \chi^{(2)} \vec{E}^2 + \epsilon_0 \chi^{(3)} \vec{E}^3 + \dots$$



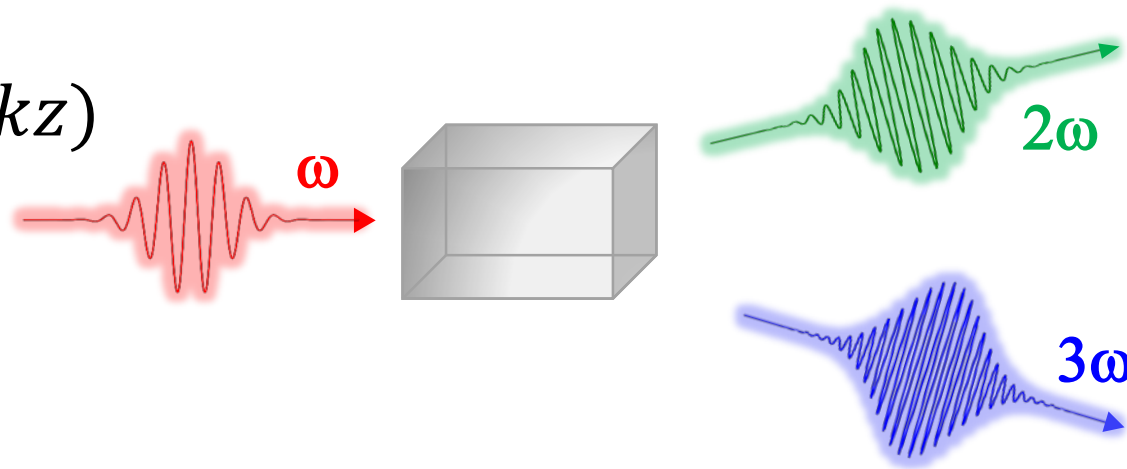
Second order susceptibility

Second Harmonic Generation (SHG)

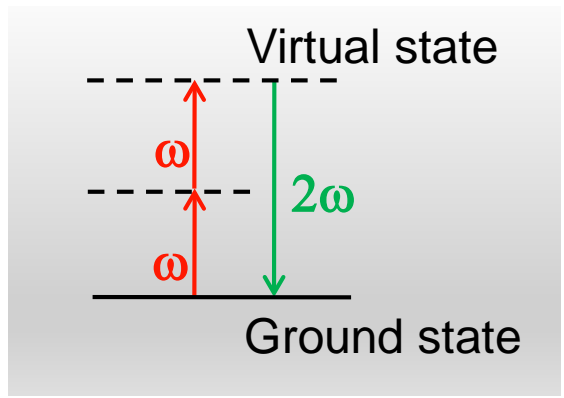
Franken et al., Phys Rev Lett, 7,4 1961

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Second Harmonic Generation



Nonlinear regime

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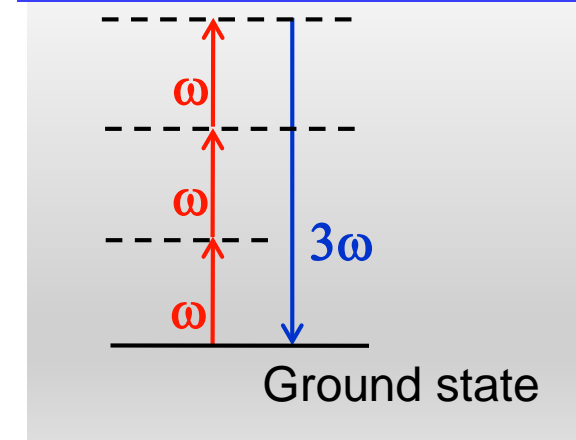
Second order susceptibility

Third order susceptibility

Second Harmonic Generation (SHG)

Third harmonic generation

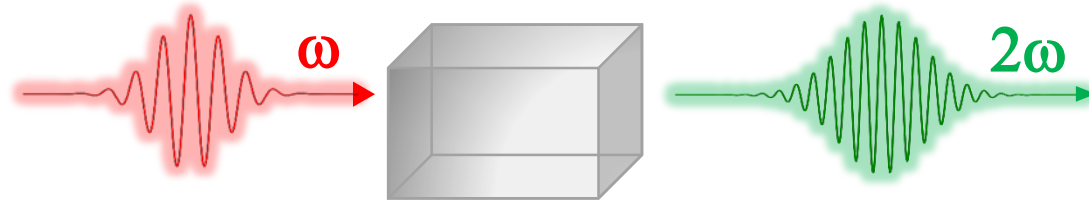
Third harmonic generation



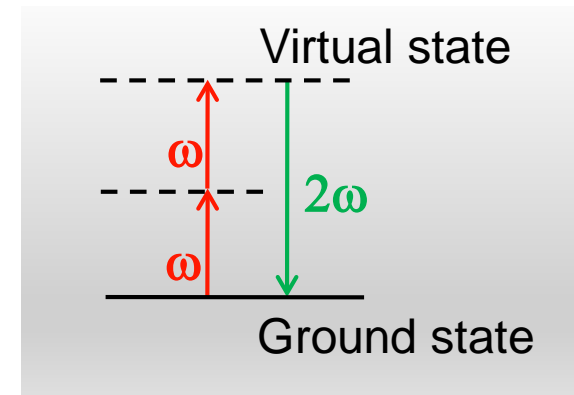
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Why is nonlinear optics useful?

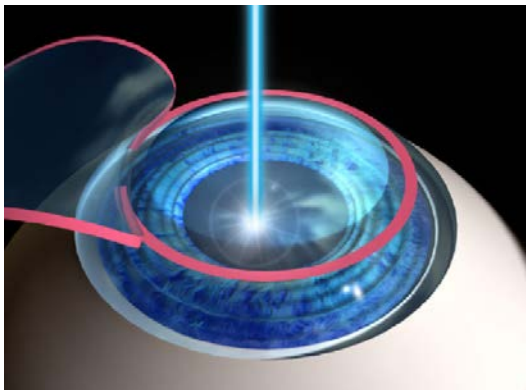
Green pointer



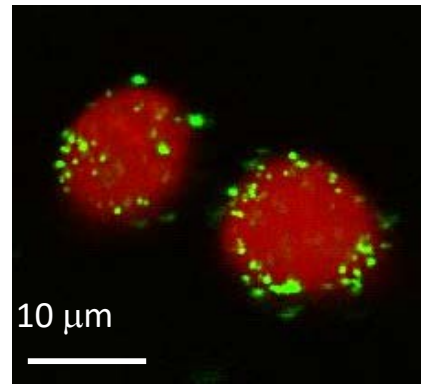
Second Harmonic Generation



Laser for eye surgery



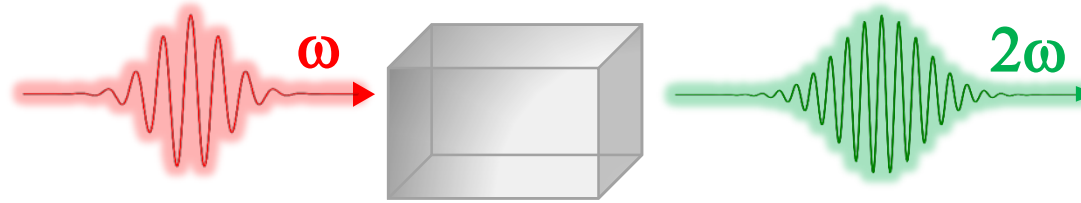
Bioimaging



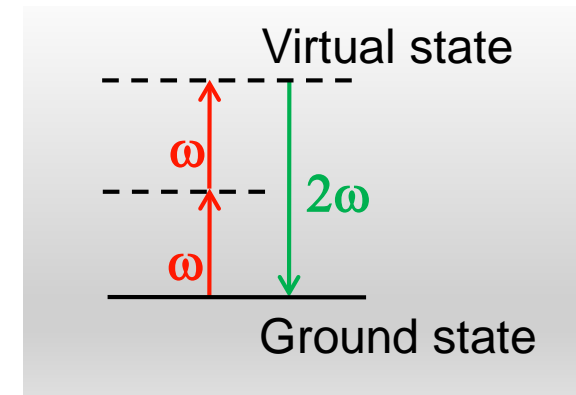
Hsieh, Grange et al, Opt. Exp. 2009

Why is nonlinear optics useful?

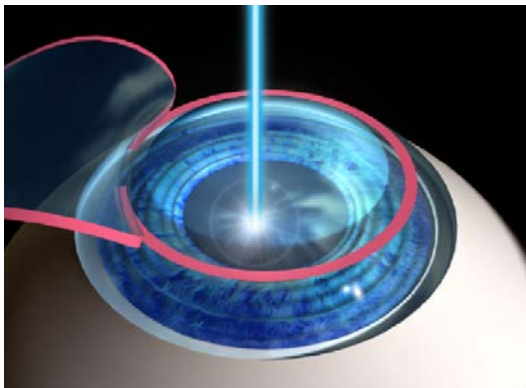
Green pointer



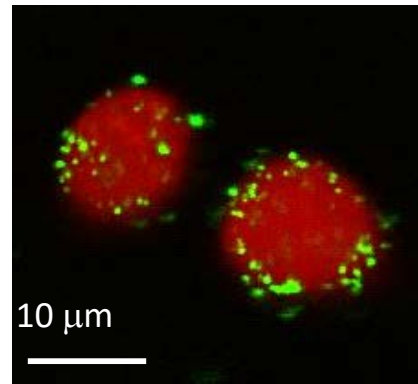
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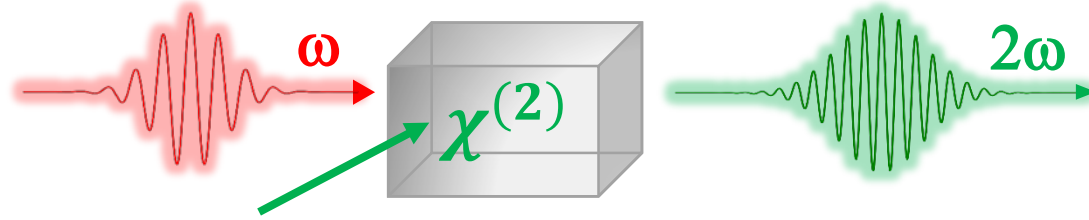


Hsieh, Grange et al, Opt. Exp. 2009

Advantages

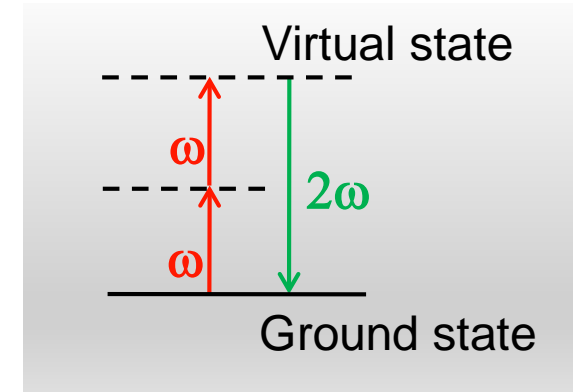
- + Photon interactions
- + Generation of new color: UV to NIR
- + Spectral filtering (no bleaching)
- + Coherent signal
- + Ultrafast response
- + Revealing crystalline structure

Why is nonlinear optics useful?

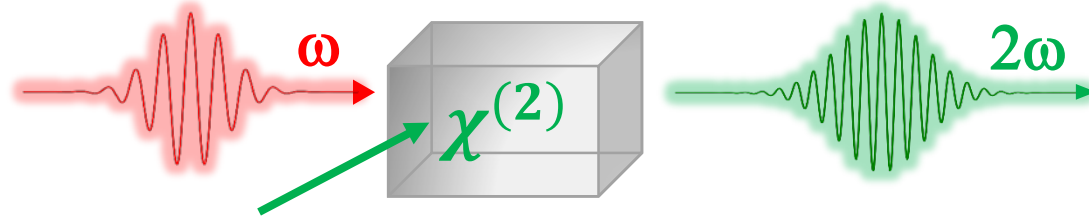


$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E} + \epsilon_0 \chi^{(2)} \vec{E}^2 + \dots$$

Second Harmonic Generation

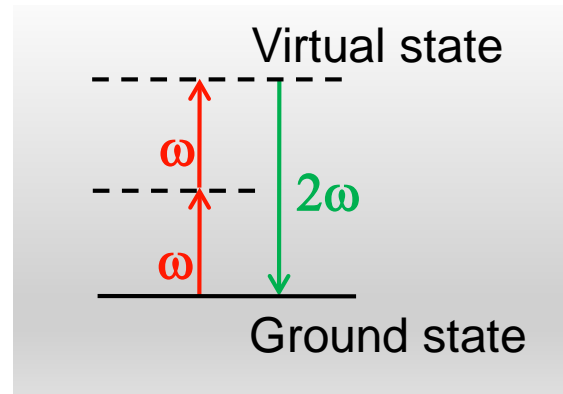


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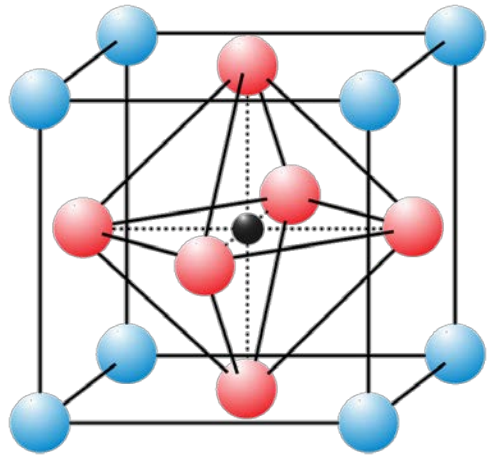
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Second Harmonic Generation



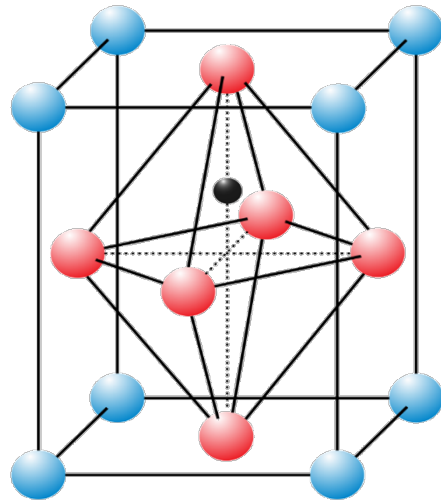
Centrosymmetric

$$\chi^{(2)} = 0$$



Non centrosymmetric

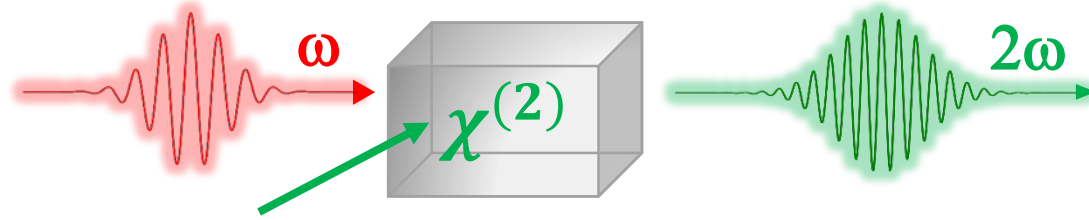
$$\chi^{(2)} \neq 0$$



Tensor coefficients

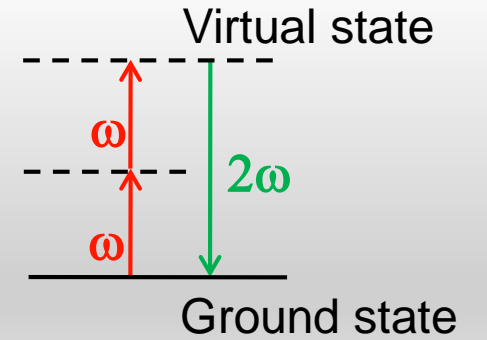
Material	$\chi^{(2)}$ coeff. (pm/V)
Silicon	0
SiO ₂ (Quartz)	0.335
BaTiO ₃	6.8 to 17
LiNbO ₃	6 to 34
GaAs	134 to 256

Why is nonlinear optics useful?



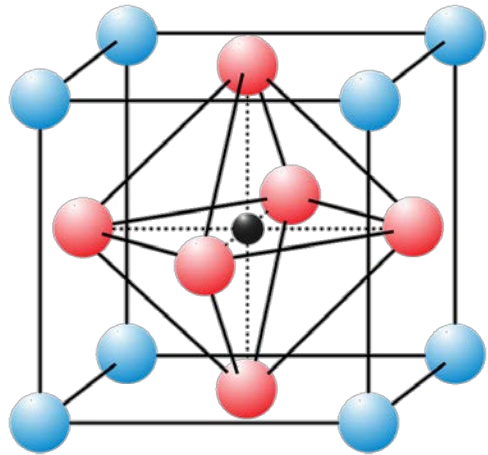
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Second Harmonic Generation



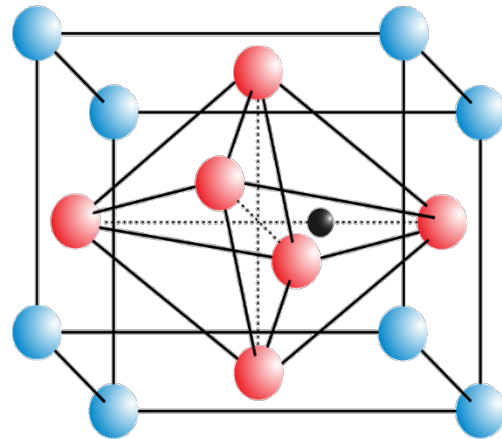
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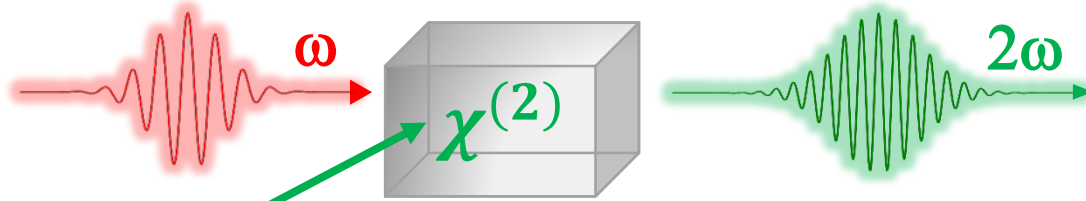
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Tensor coefficients

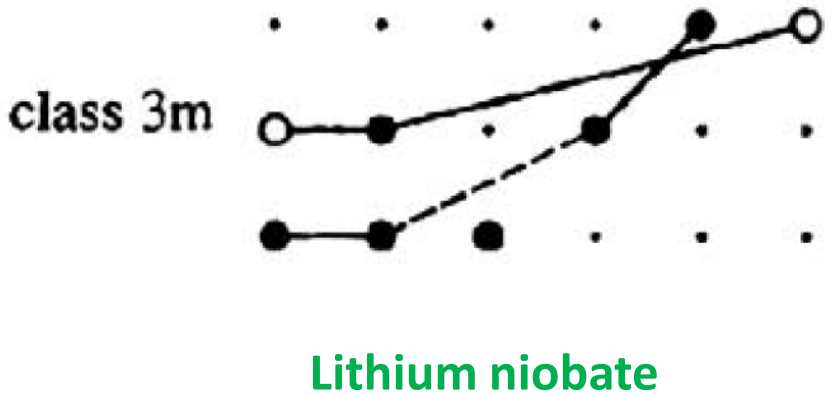
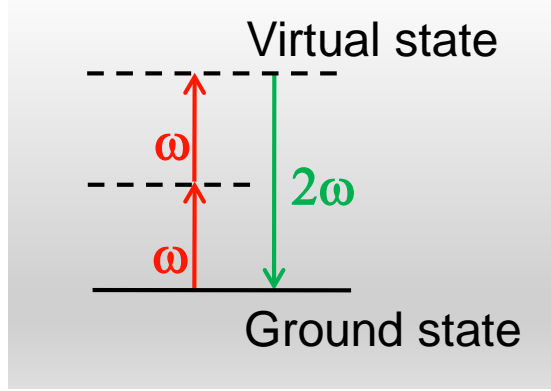
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Second Harmonic Generation

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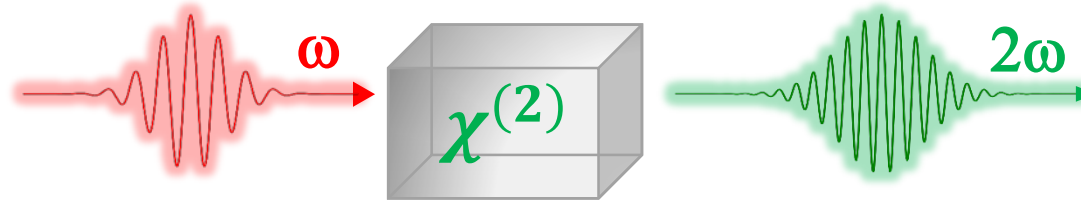


Tensor coefficients

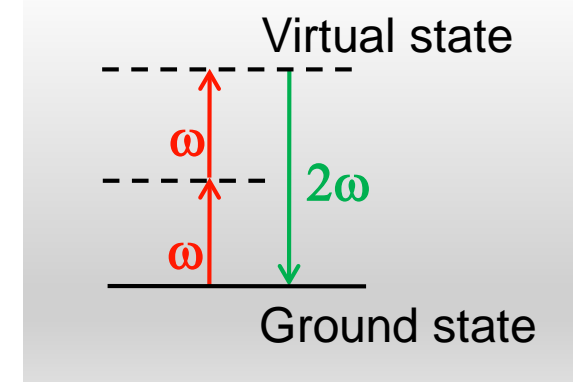
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R. W. Boyd, Nonlinear Optics (Academic Press, 2008)

Why the metal-oxides ?



Second Harmonic Generation

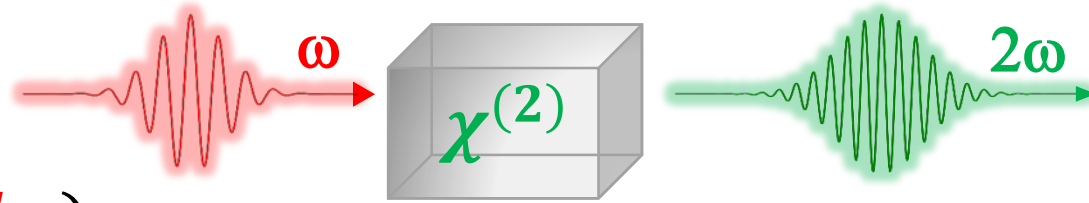


Advantages of metal-oxides ($\chi^{(2)}$ materials)

- Barium titanate, Lithium niobate as bulk
- High nonlinearities
- Large band gap: 3-4 eV
- Refractive index > 2
- Electro-optic, ferroelectric,...

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How is the interaction within the material ?

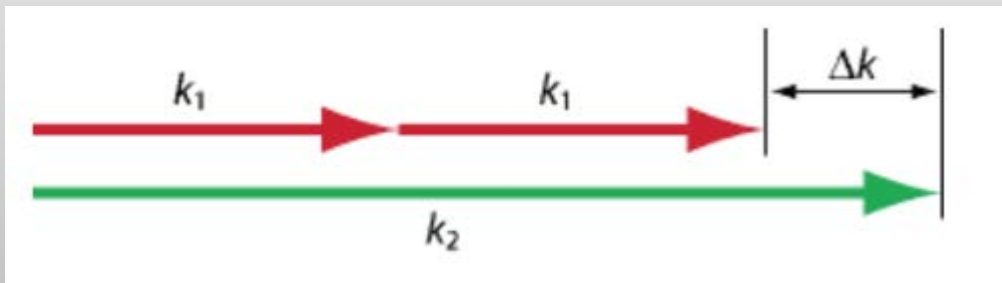


$$\vec{E}(t) = A \cos(\omega t + kz)$$

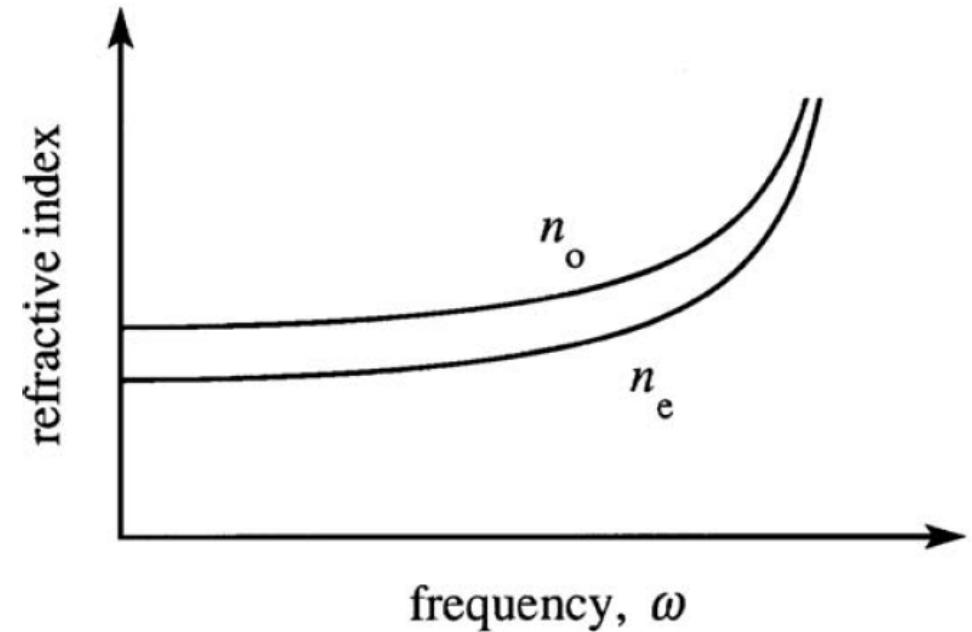
wavenumber

Phase mismatch due to dispersion

Phase mismatch $\Delta k \neq 0$

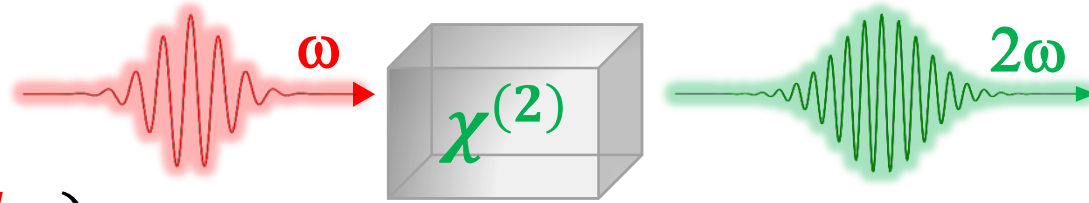


Chromatic dispersion : $n(\omega) \neq n(2\omega)$



R. W. Boyd, *Nonlinear Optics*, Academic Press, 2008).

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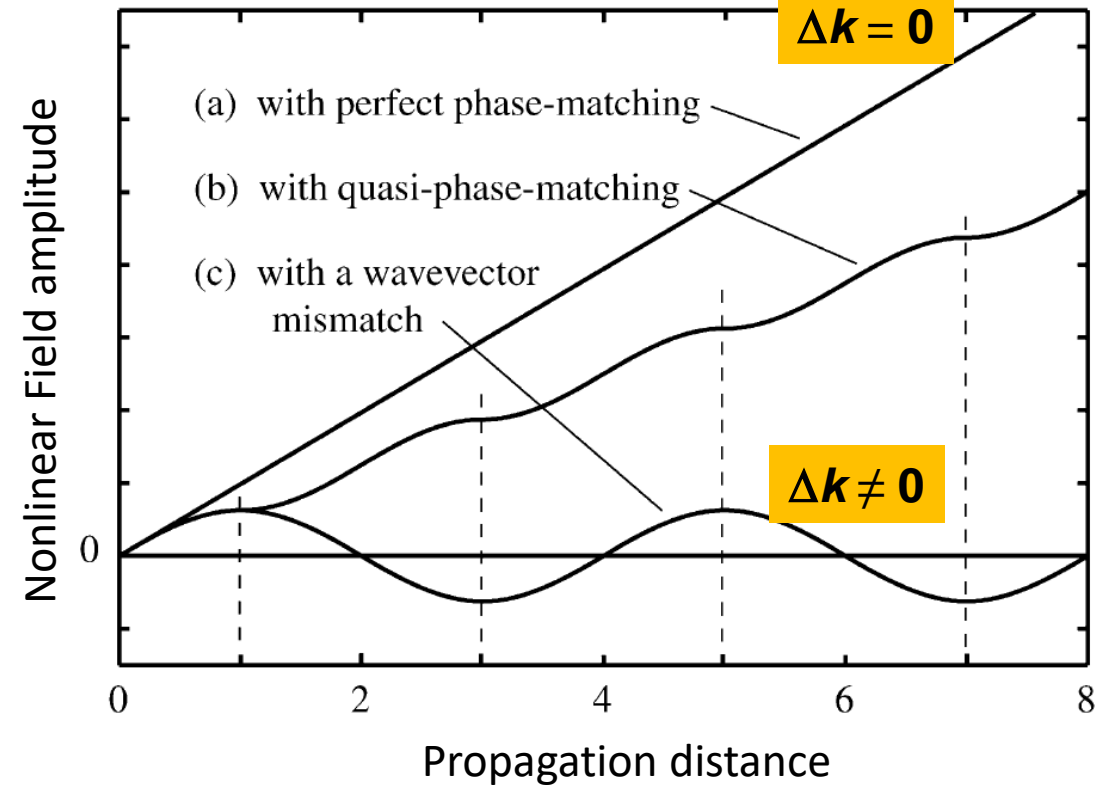
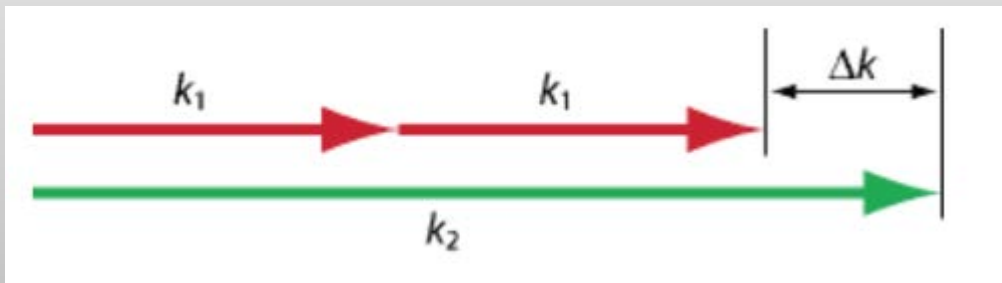


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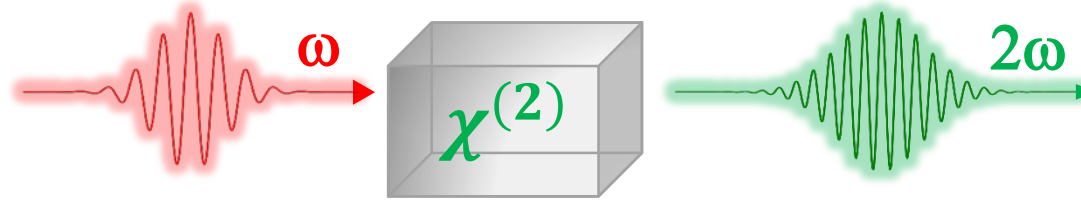
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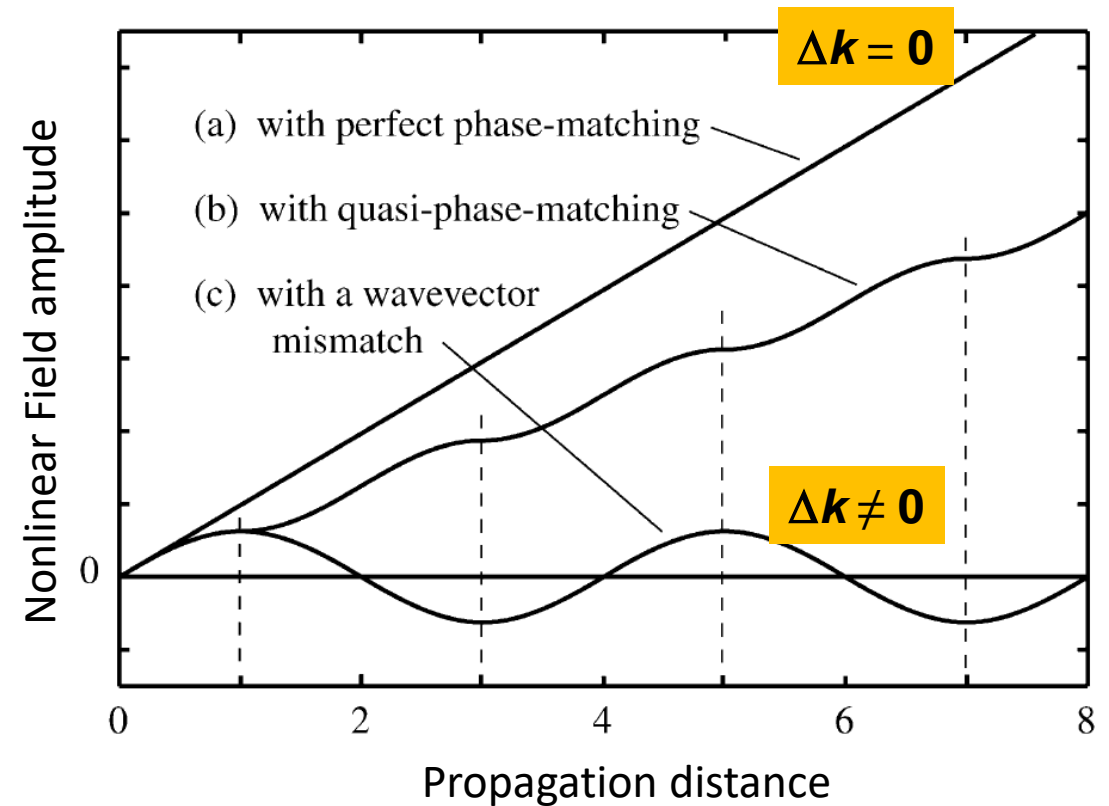
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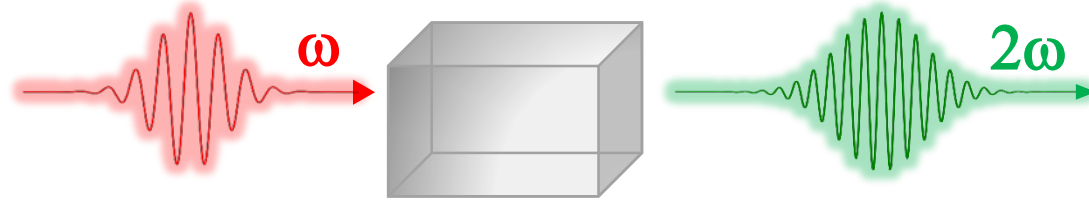


How to deal with phase matching ?

- Periodical poling
- Angle or temperature tuning
- In waveguides:
 - modal phase-matching
- Reducing the size of the material



Limitations of current nonlinear optics



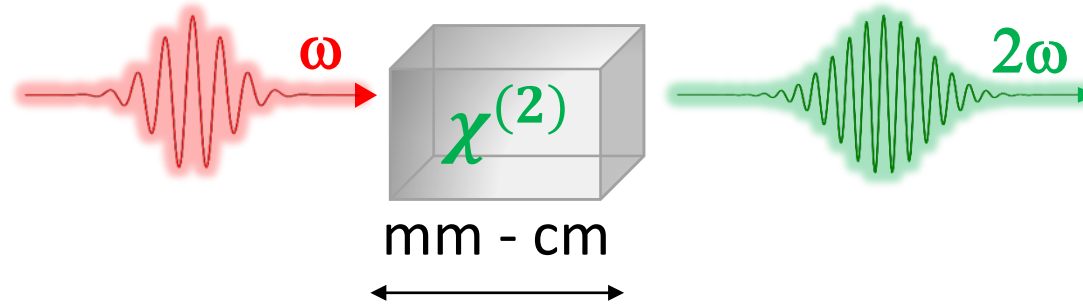
$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E} + \epsilon_0 \chi^{(2)} \vec{E}^2 + \dots$$

Limitations

- High power -> high energy cost and damage
- Bulky material -> no integration
- Phase mismatch -> low output signal
- Small sizes -> reduced signal output

My goal

Limitations of current nonlinear optics



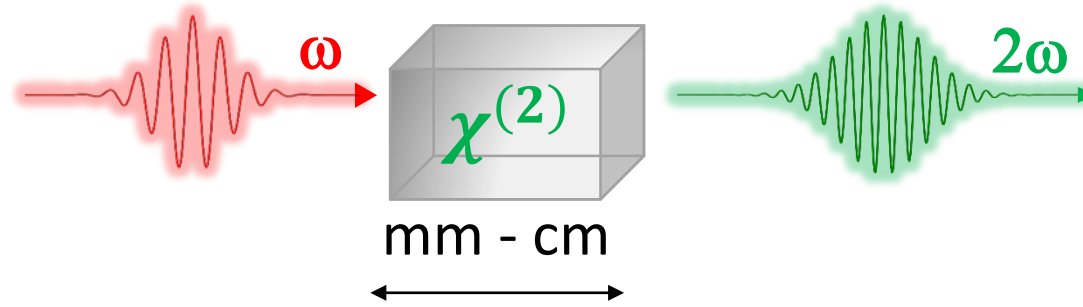
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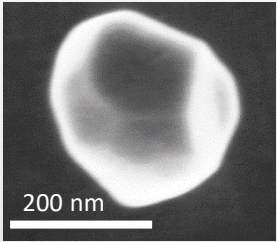
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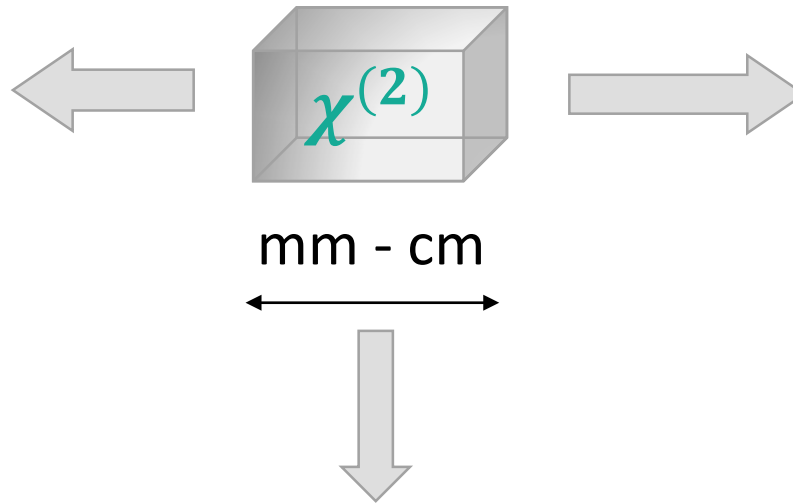
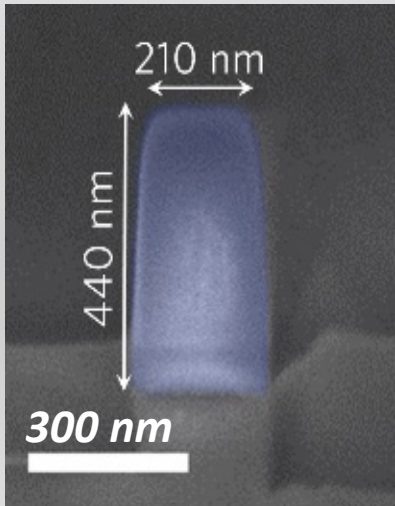
**Resonance or confinement
in particles and thin films
to achieve compact devices**

$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks

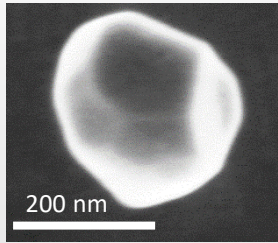


20 - 100s nm

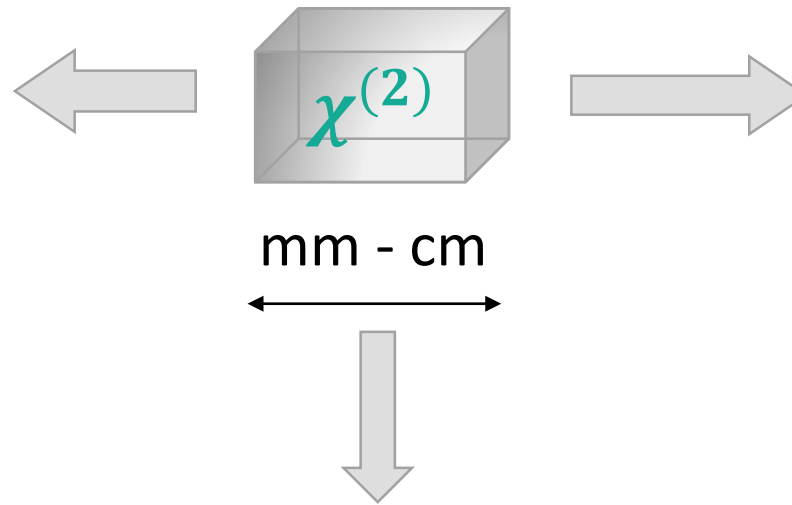
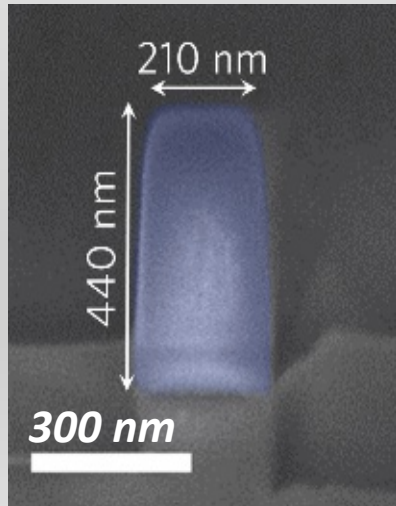


$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks

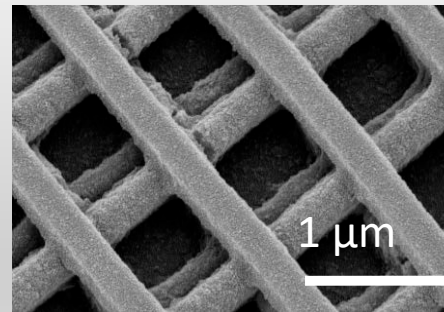


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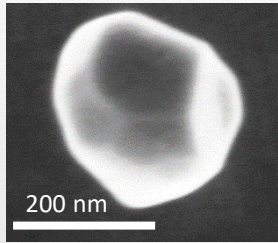
Bottom-up processing

Wide surface area photonic crystal

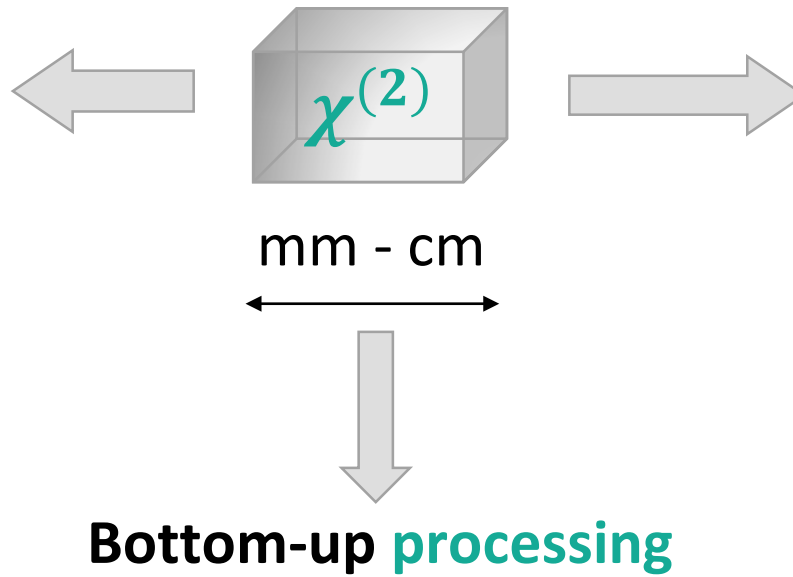
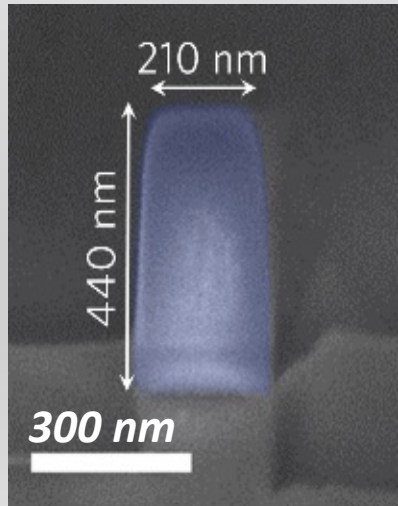


$\chi^{(2)}$ materials at small scale

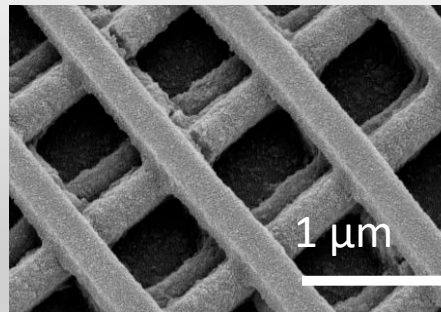
$\chi^{(2)}$ Building blocks



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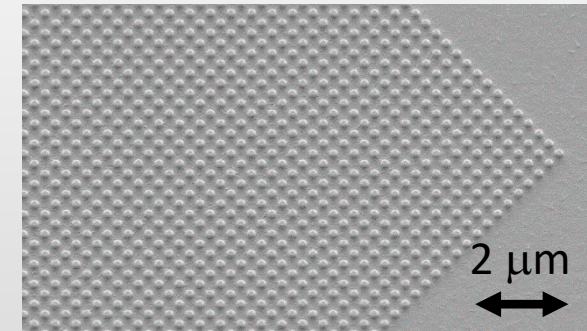


Wide surface area photonic crystal

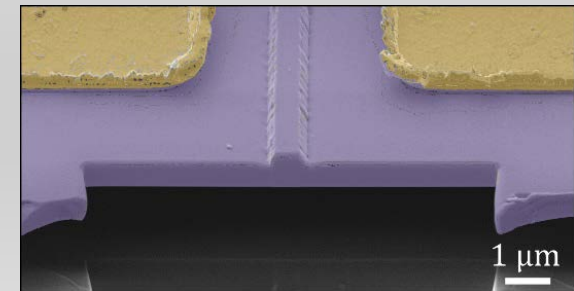


Top-down fabrication

Flat photonic structures

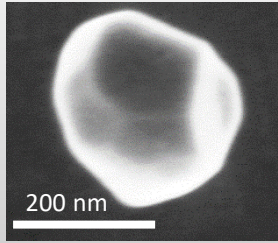


Integrated devices



A powerful characterization tool

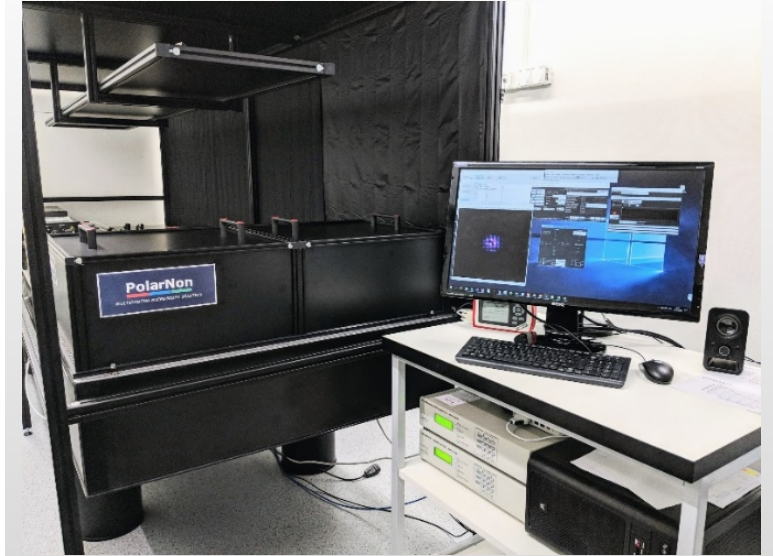
Understanding building blocks



20 - 100s nm

Diffraction limit

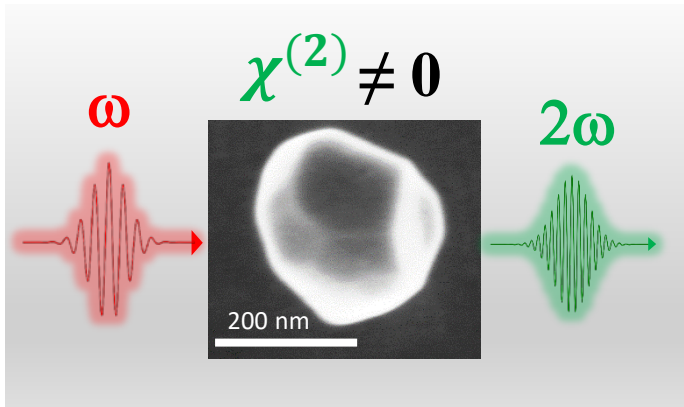
$$d = \frac{\lambda}{2NA}$$



Super-resolution multiphoton polarimetric microscope

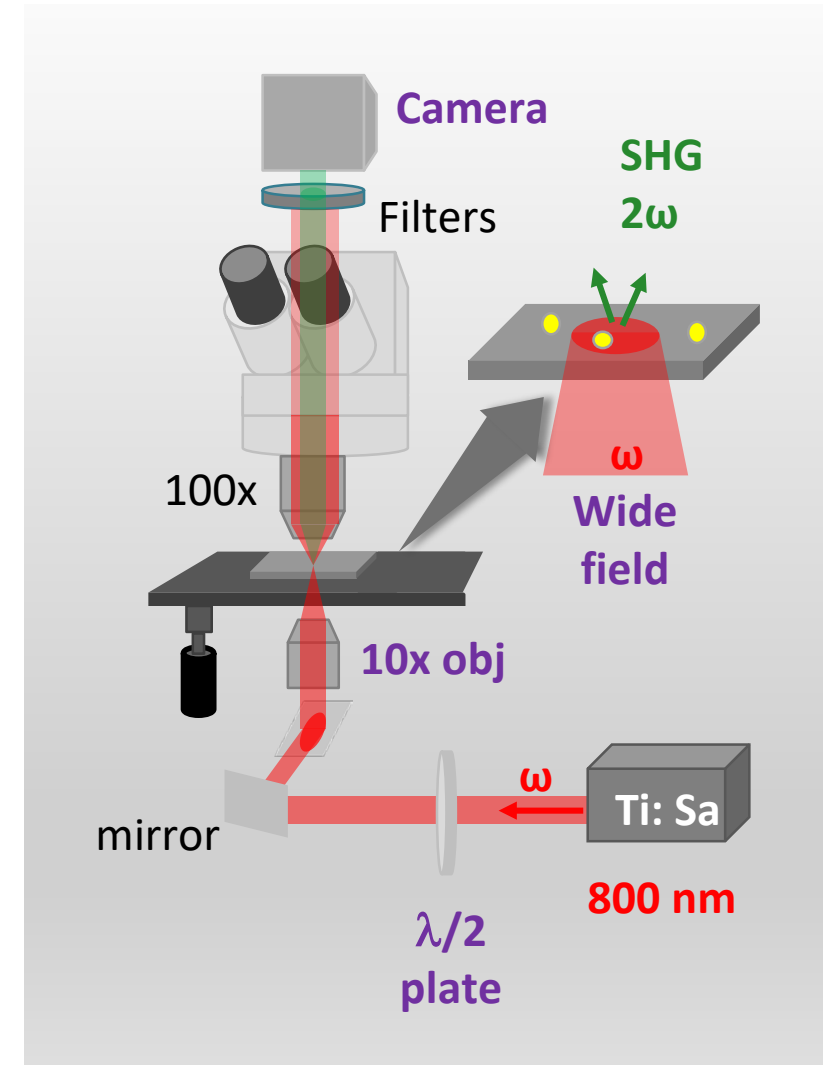
2018 competitors at the SPIE Europe Photonics
Village in Strasbourg / [Patent EP18193103.1](#)

A powerful characterization tool



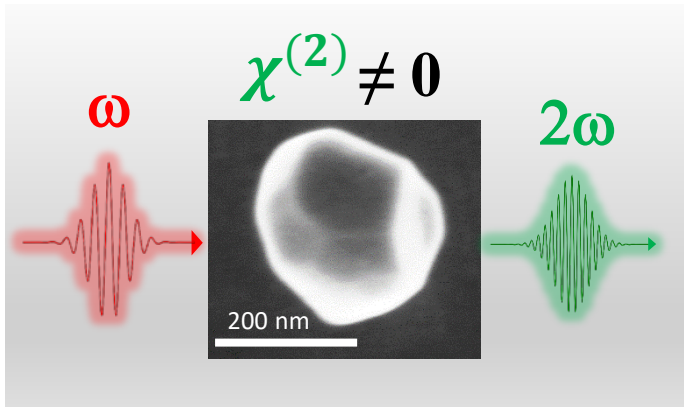
Diffraction limit

$$d = \frac{\lambda}{2NA}$$



2018 competitors at the SPIE Europe Photonics Village in Strasbourg / [Patent EP18193103.1](#)

A powerful characterization tool

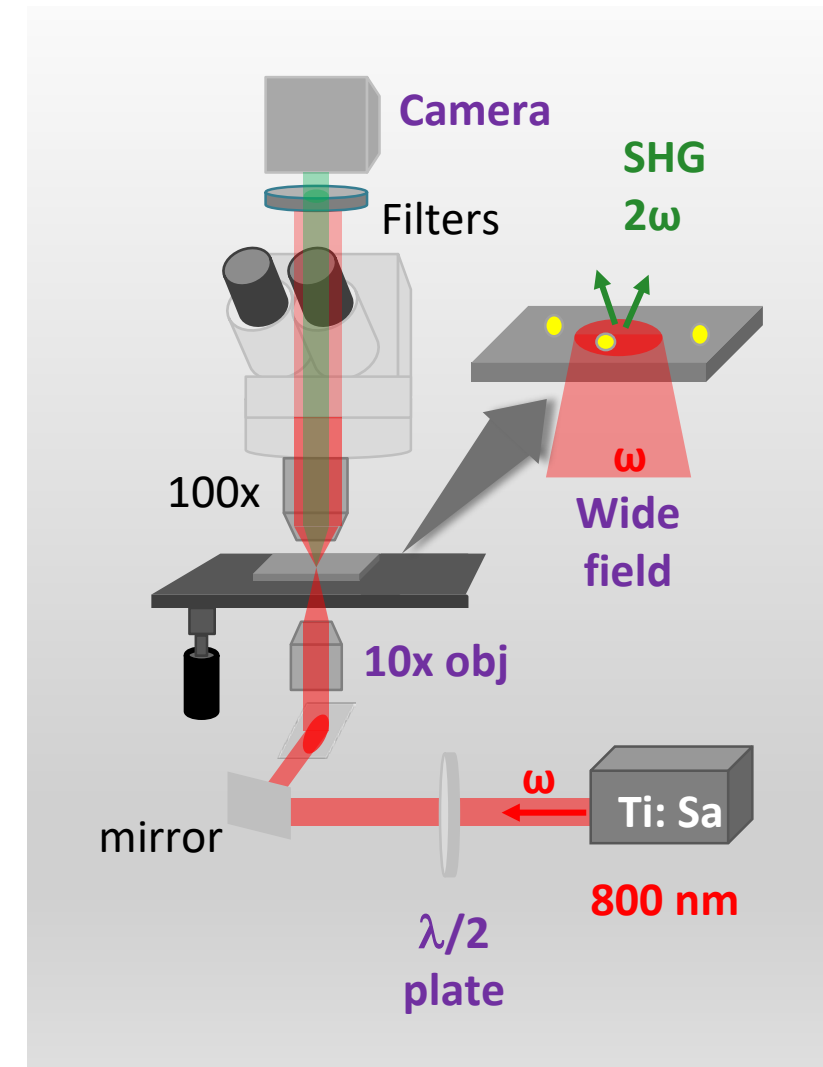


Diffraction limit

$$d = \frac{\lambda}{2NA}$$

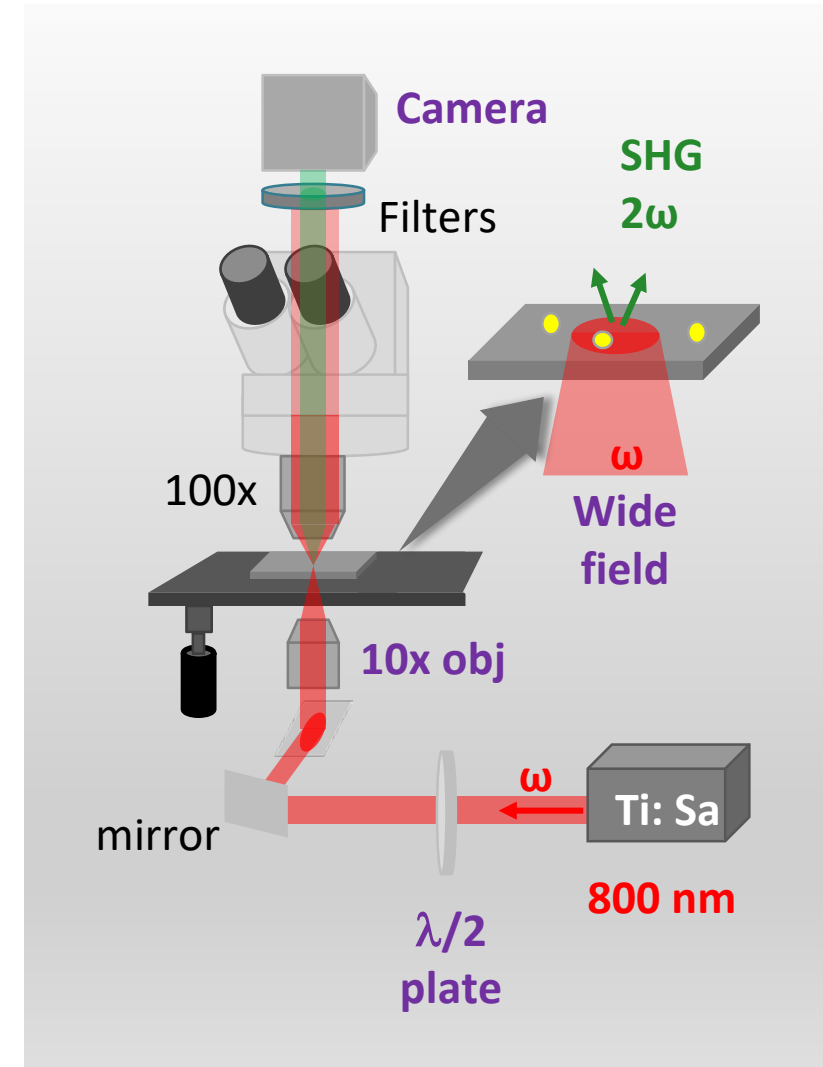
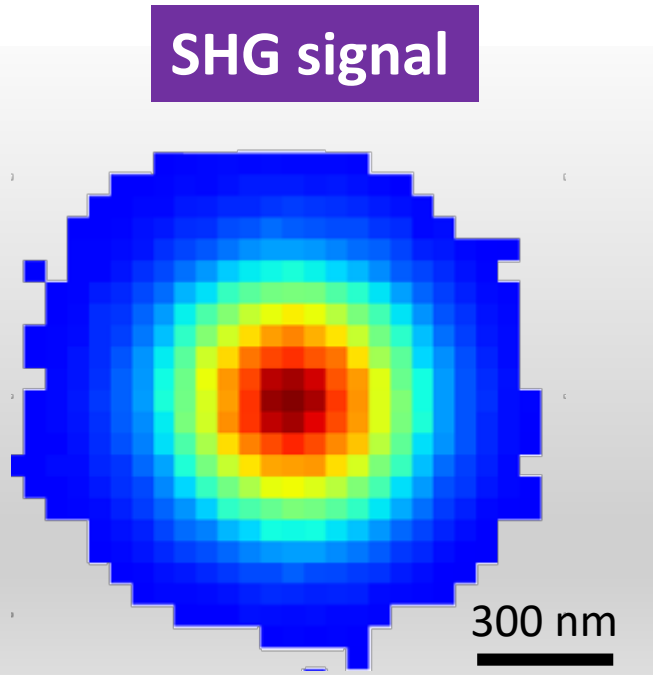
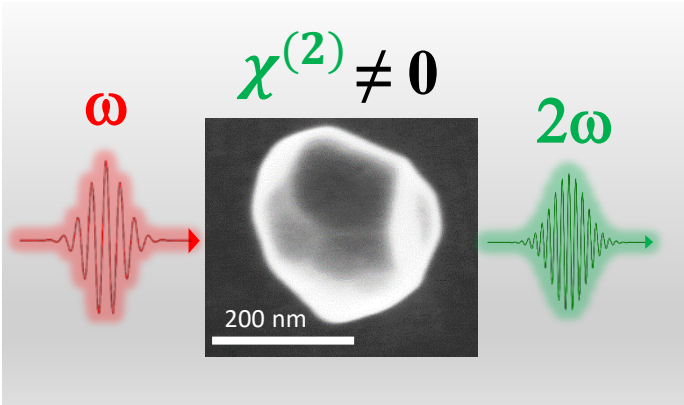
Special features

- Incident polarization rotation
- Wide field illumination
- Single pixel detection



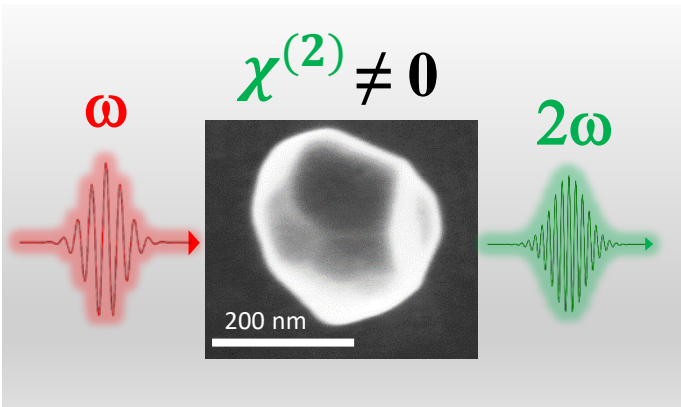
2018 competitors at the SPIE Europe Photonics Village in Strasbourg / [Bridge Funding](#)

A powerful characterization tool



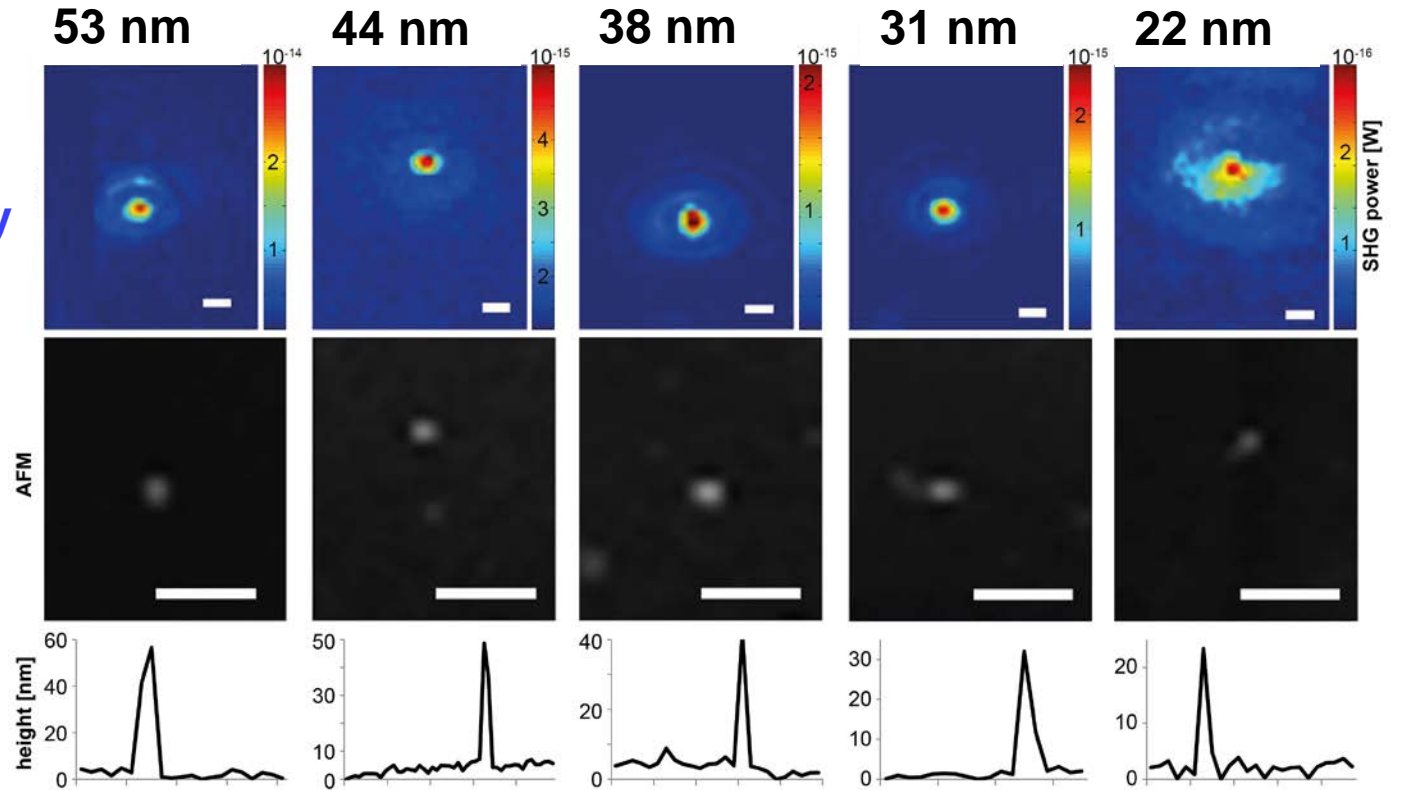
2018 competitors at the SPIE Europe Photonics Village in Strasbourg / [Bridge Funding](#)

Average SHG intensity of single particles

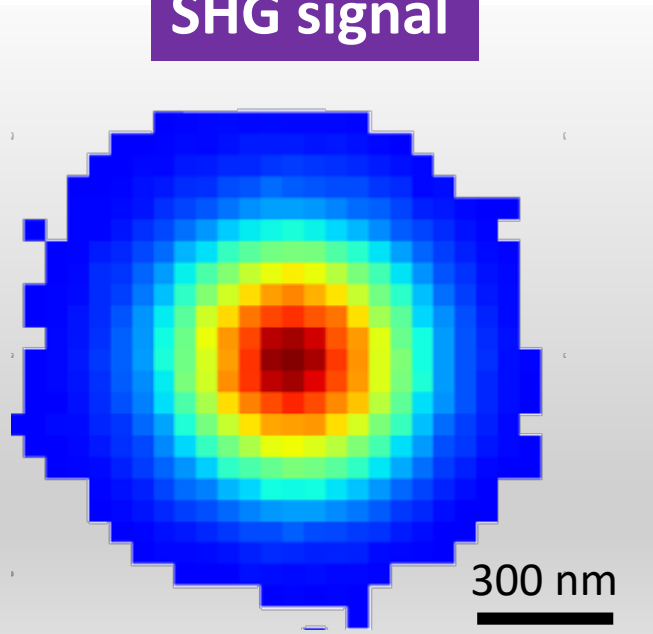


BaTiO₃ particles with diameters of :

SHG intensity

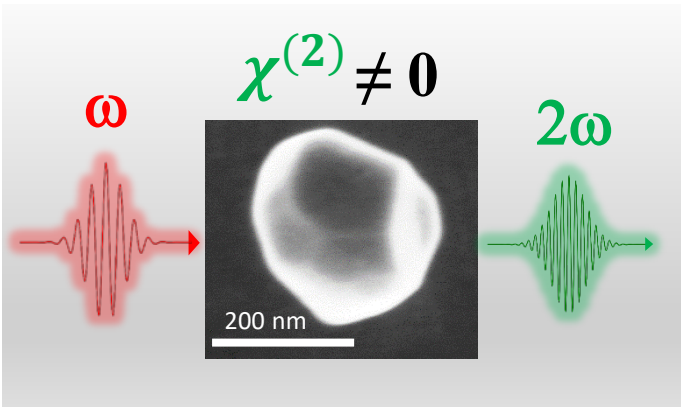


SHG signal



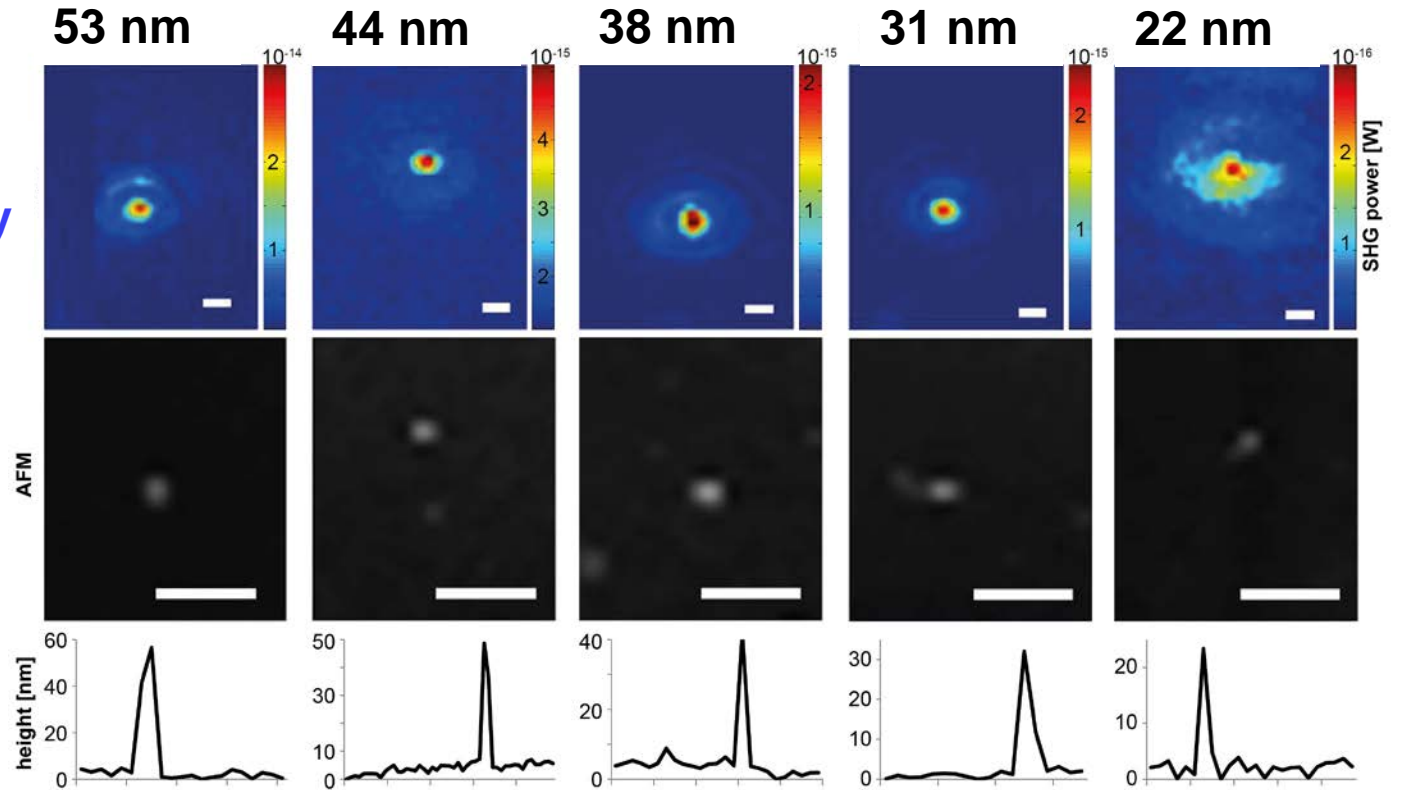
E. Kim, A. Steinbrück, M. T. Buscaglia, V. Buscaglia, T. Pertsch, and R. Grange, ACS Nano 7, 5343–5349, 2013.

Average SHG intensity of single particles

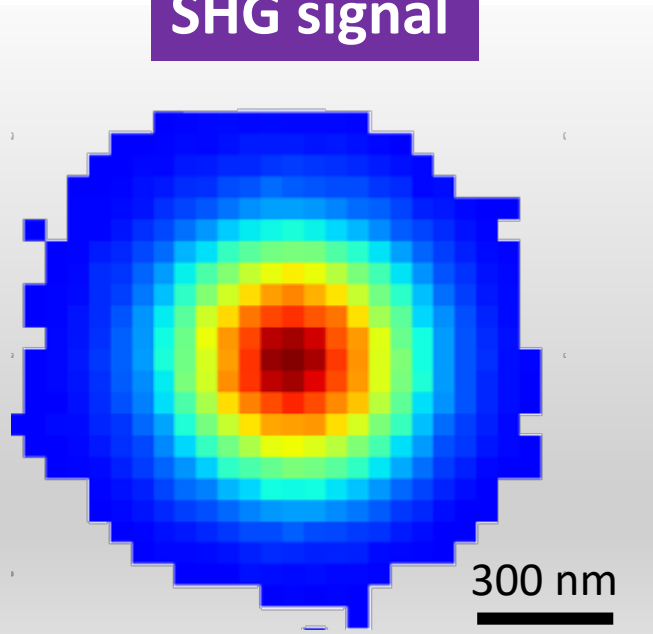


BaTiO₃ particles with diameters of :

SHG intensity



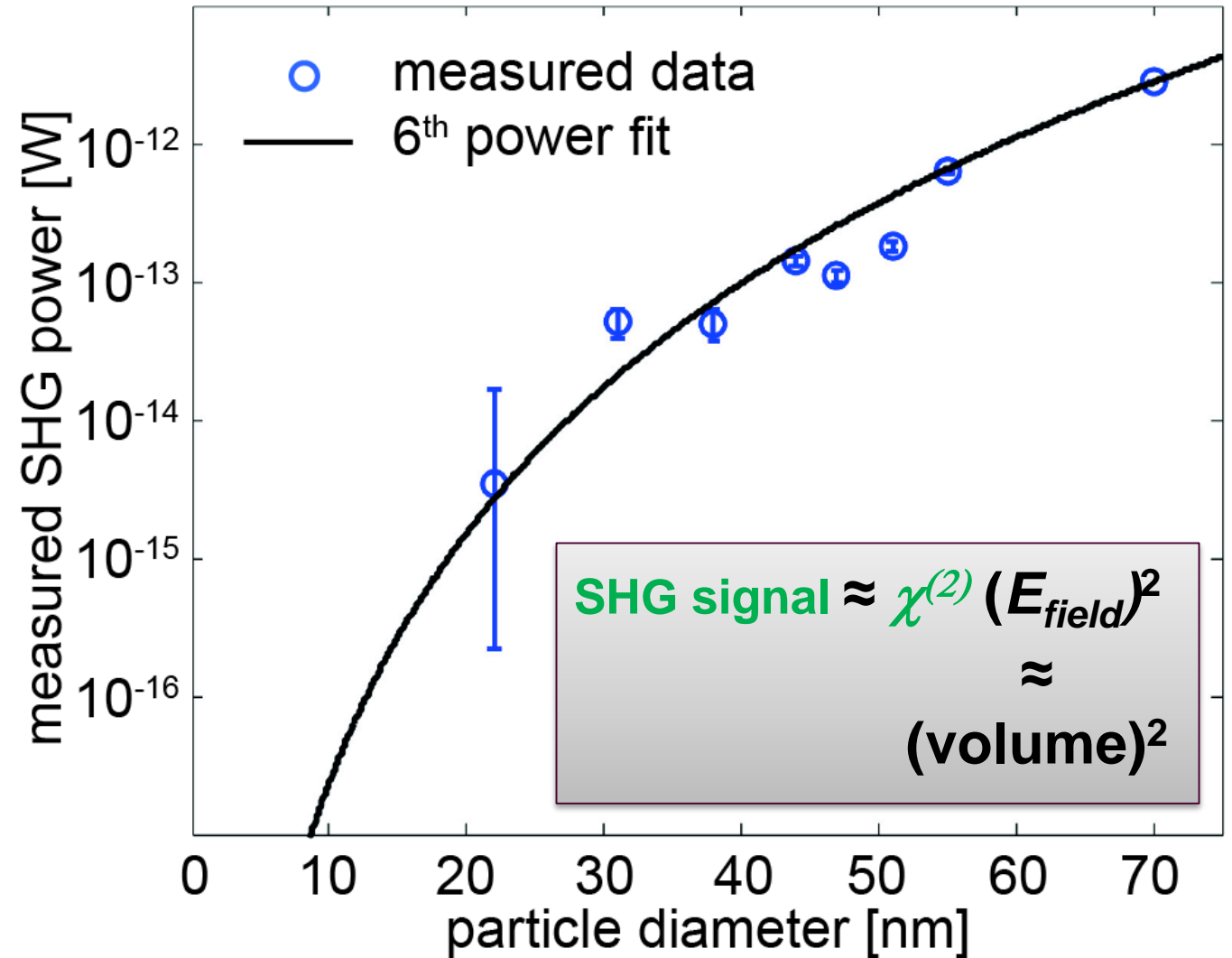
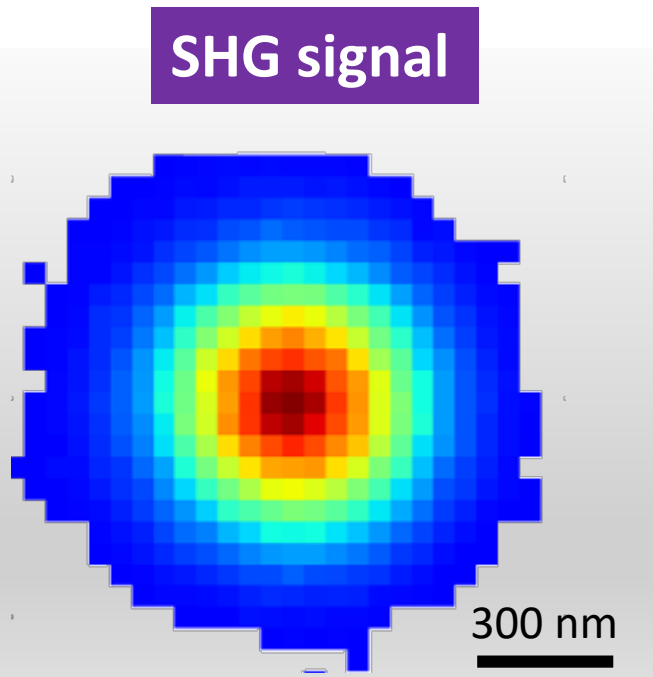
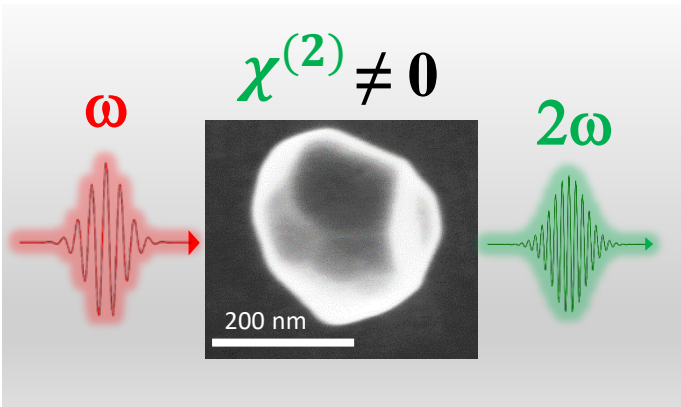
SHG signal



Particles are bright down to 20 nm...

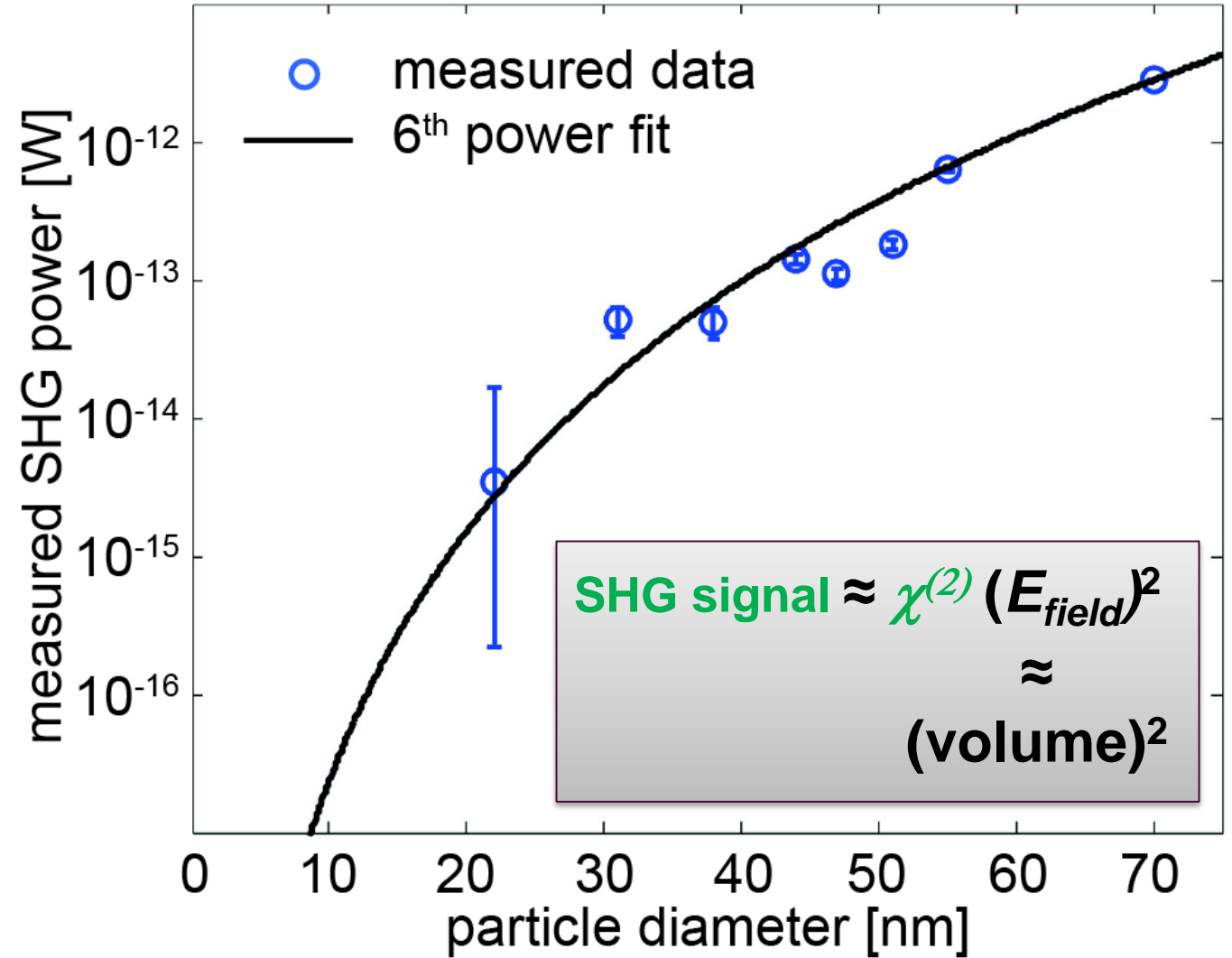
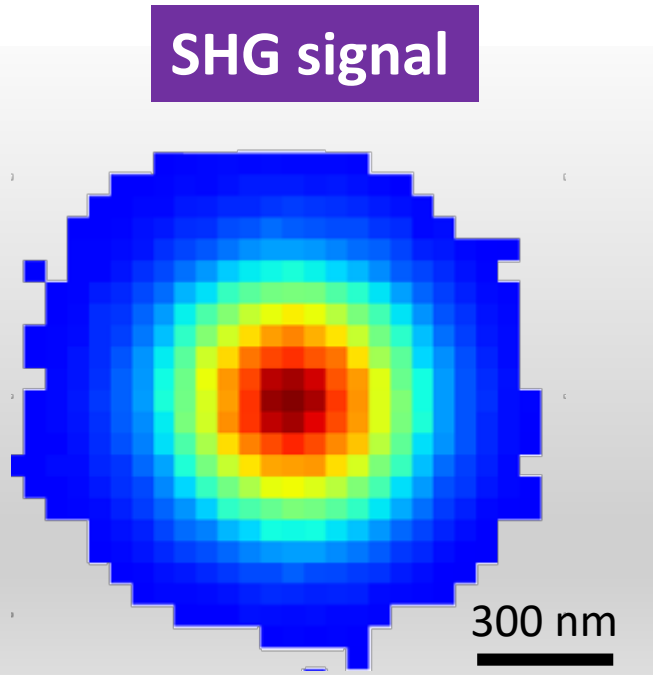
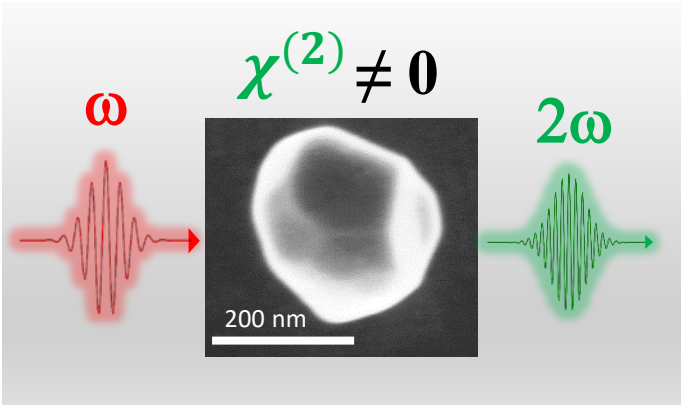
E. Kim,...,and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva,..., R. Grange Nano Lett.,16, 6290, 2016.

Average SHG intensity of single particles



E. Kim, A. Steinbrück, M. T. Buscaglia, V. Buscaglia, T. Pertsch, and R. Grange, ACS Nano 7, 5343–5349, 2013.

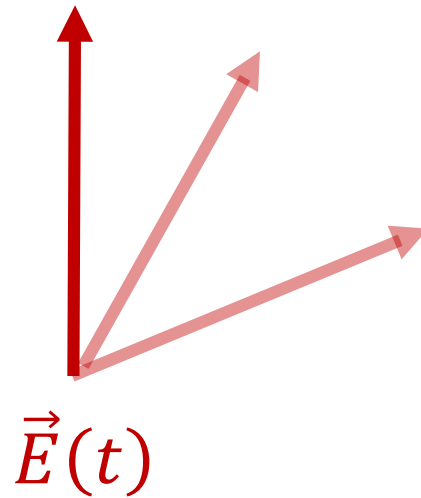
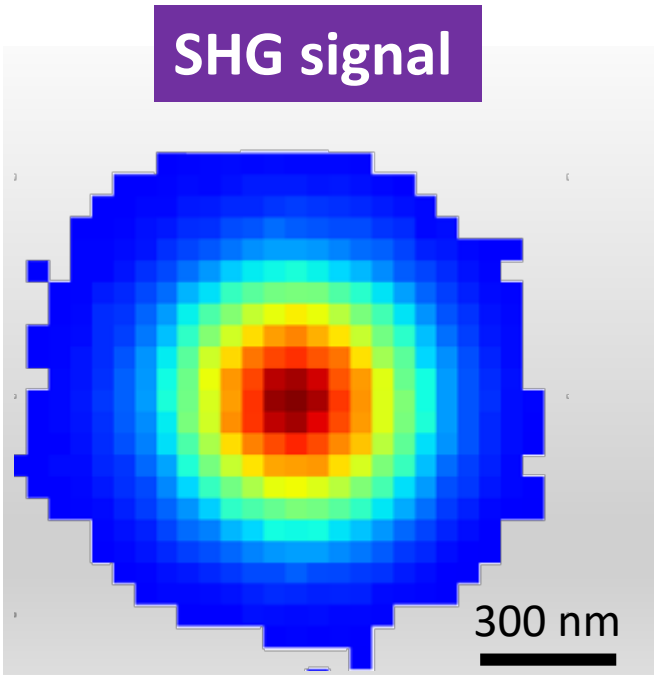
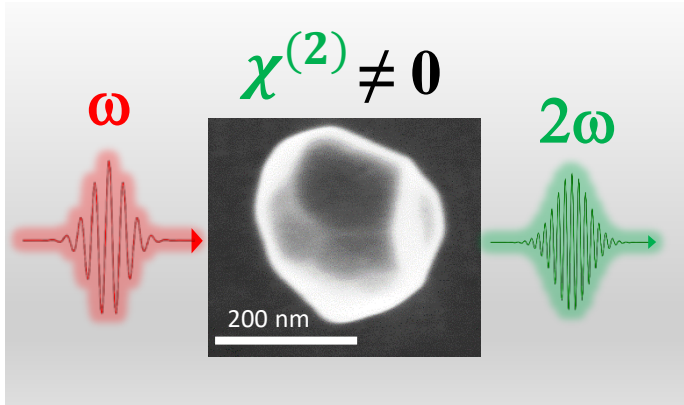
Average SHG intensity of single particles



... room for improvements

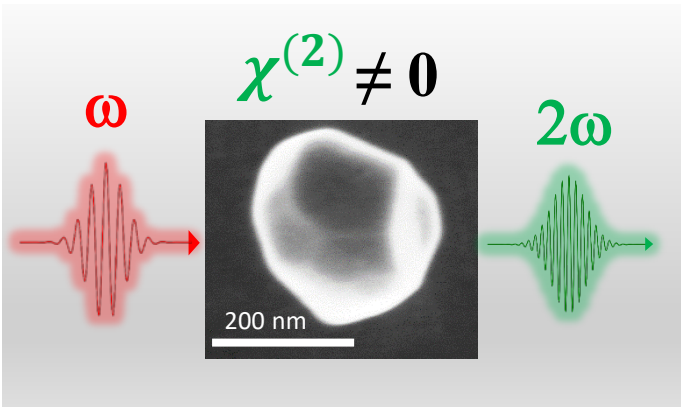
E. Kim,...,and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva,..., R. Grange Nano Lett.,16, 6290, 2016.

Average SHG intensity of single particles



E. Kim,...,and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva,..., R. Grange Nano Lett.,16, 6290, 2016.

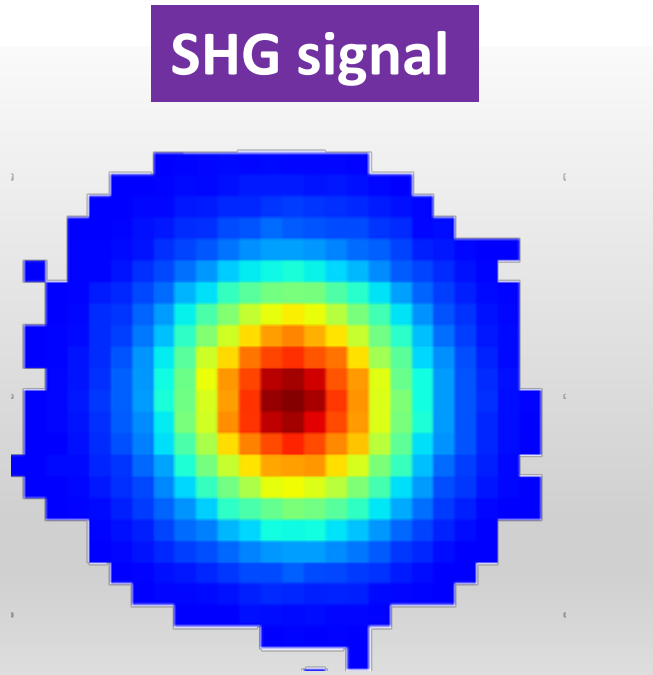
Average SHG intensity of single particles



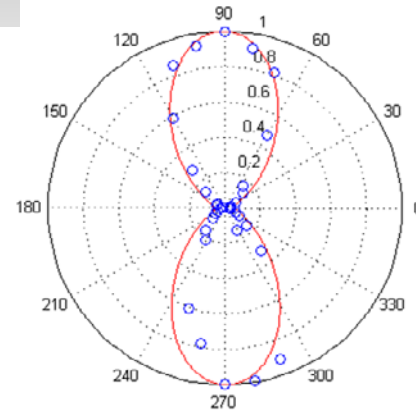
Average response over a single particle

$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E} + \epsilon_0 \chi^{(2)} \vec{E}^2 + \dots$$

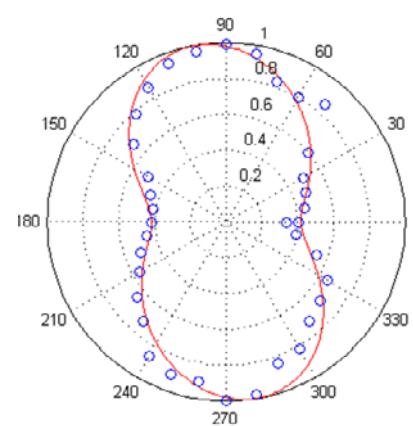
— Bulk fit
○ Measured value



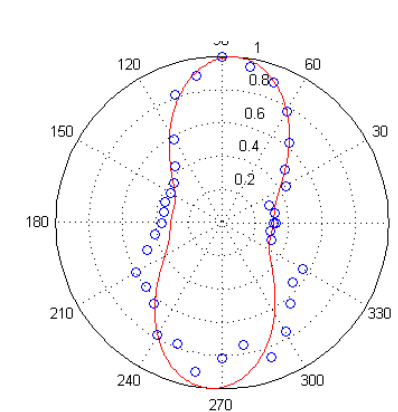
53 nm



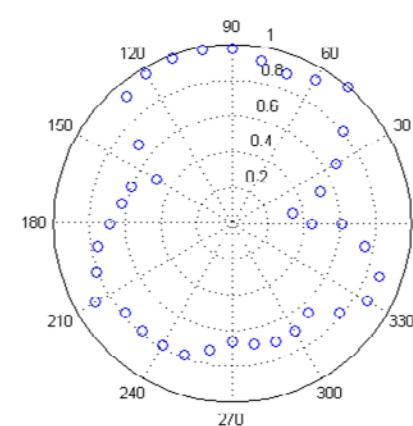
D = 44 nm



D = 31 nm

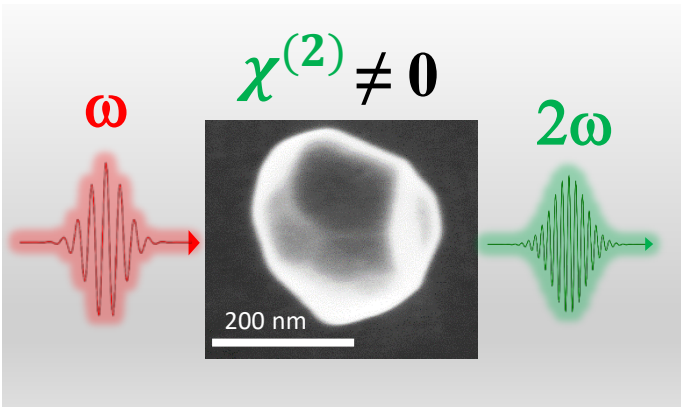


D = 22 nm



E. Kim, ..., and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva, ..., R. Grange Nano Lett., 16, 6290, 2016.

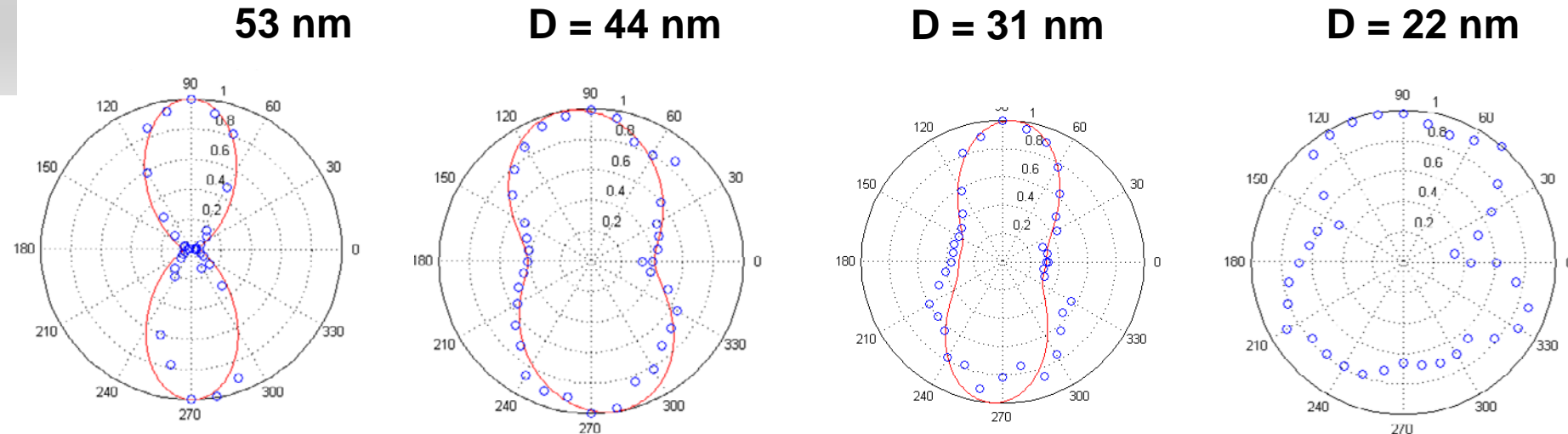
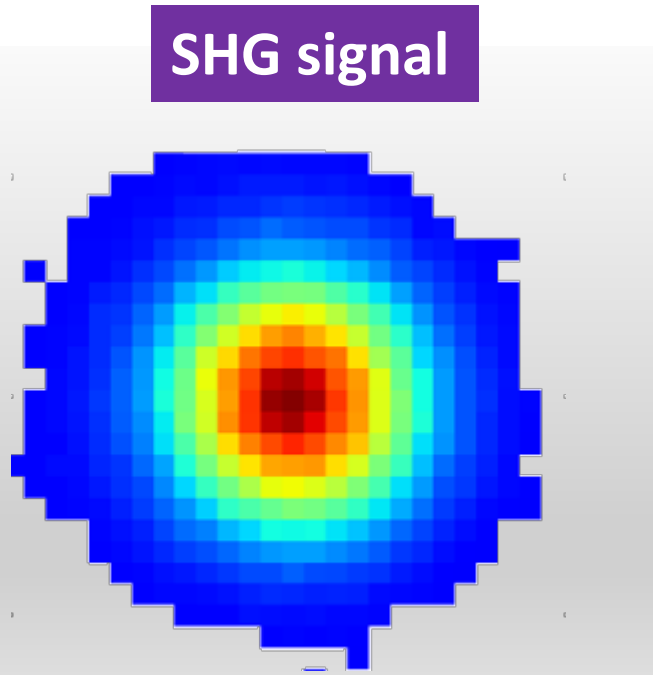
Average SHG intensity of single particles



Average response over a single particle

$$\vec{P} = \epsilon_0 \chi^{(1)} \vec{E} + \epsilon_0 \chi^{(2)} \vec{E}^2 + \dots$$

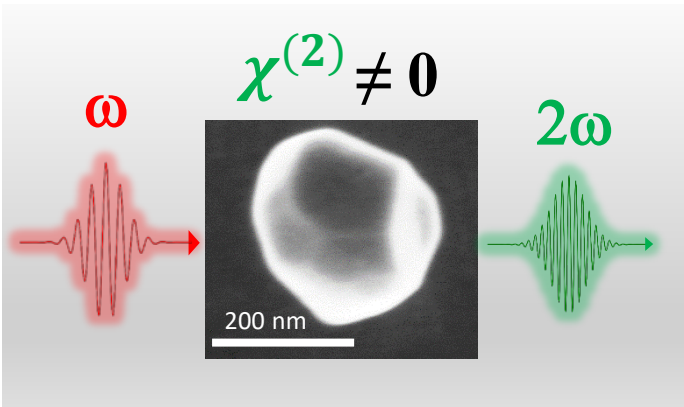
— Bulk fit
○ Measured value



Uncertain crystallography
for small particle sizes

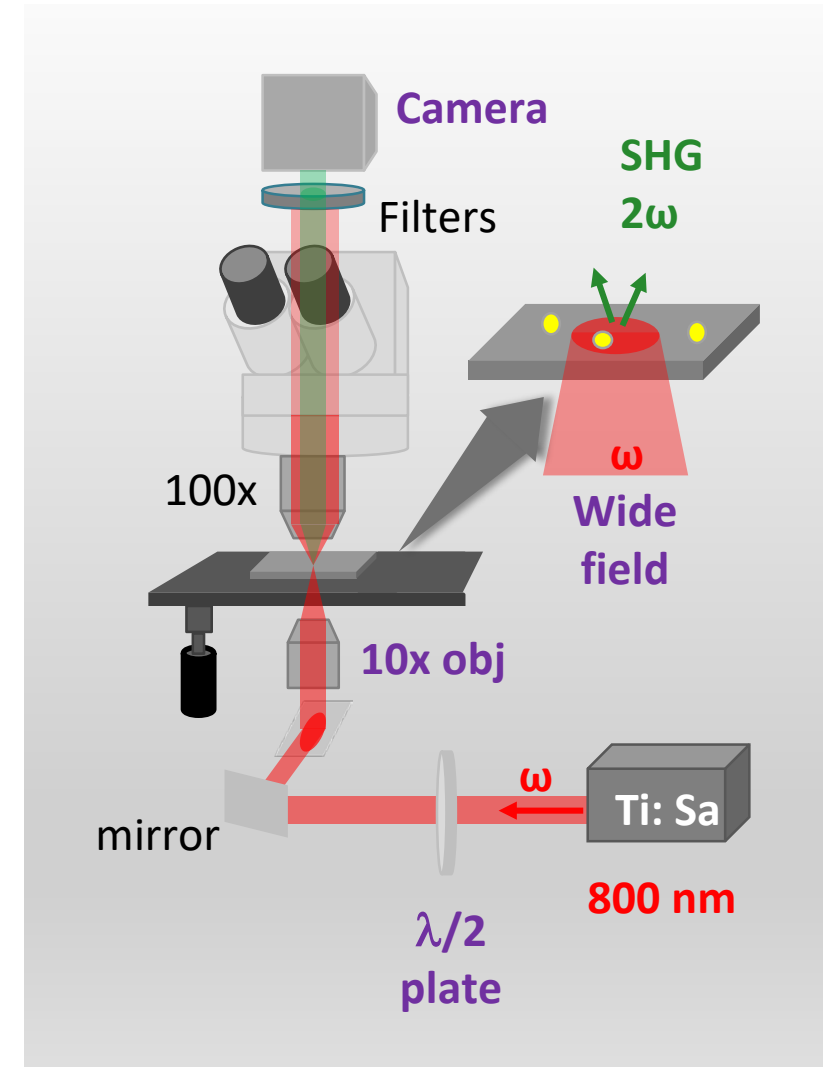
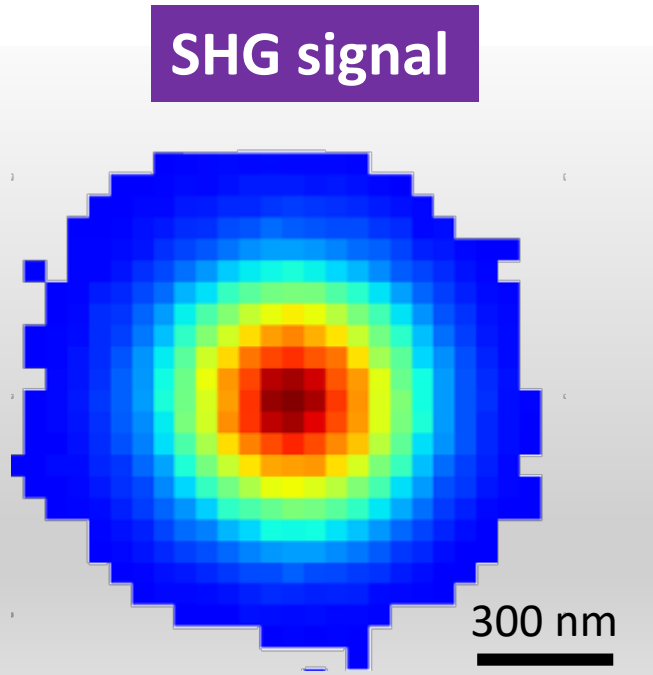
E. Kim, ..., and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva, ..., R. Grange Nano Lett., 16, 6290, 2016.

A powerful characterization tool

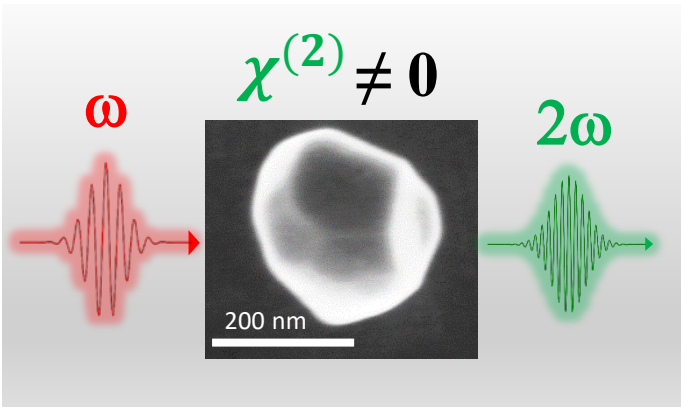


SHG intensity is a fourth power dependence

$$I(\alpha) \propto |\chi^{(2)}: E(\alpha)E(\alpha)|^2$$



A powerful characterization tool



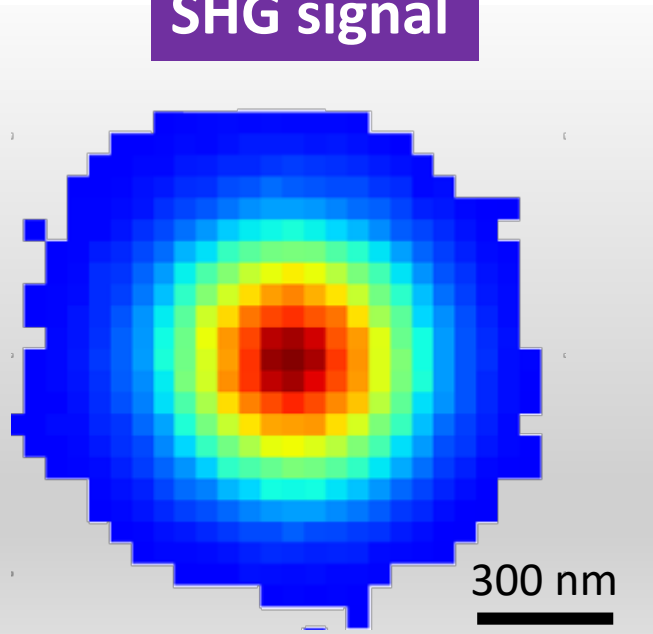
SHG intensity is a fourth power dependence

$$I(\alpha) \propto |\chi^{(2)}: E(\alpha)E(\alpha)|^2$$

Decomposition into functions

$$I(\alpha)/I_0 \propto 1 + a_2 \cos 2\alpha + b_2 \sin 2\alpha + a_4 \cos 4\alpha + b_4 \sin 4\alpha$$

SHG signal



Amplitudes

$$I_2 = \sqrt{a_2^2 + b_2^2}, \quad I_4 = \varepsilon_4 \sqrt{a_4^2 + b_4^2}$$

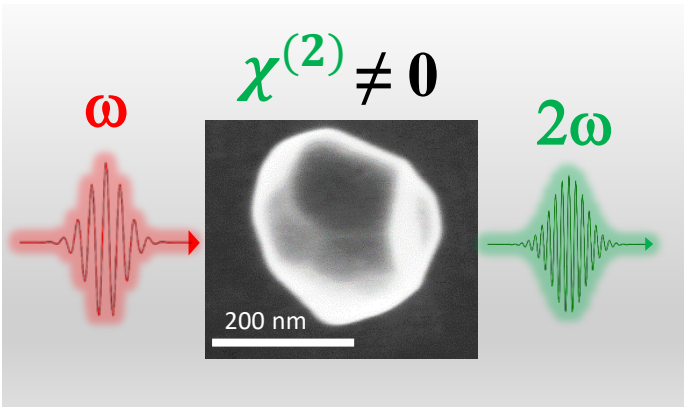
Orientations

$$\varphi_2 = \frac{1}{2} \arctan \frac{b_2}{a_2}, \quad \varphi_4 = \frac{1}{4} \arctan \frac{b_4}{a_4}$$

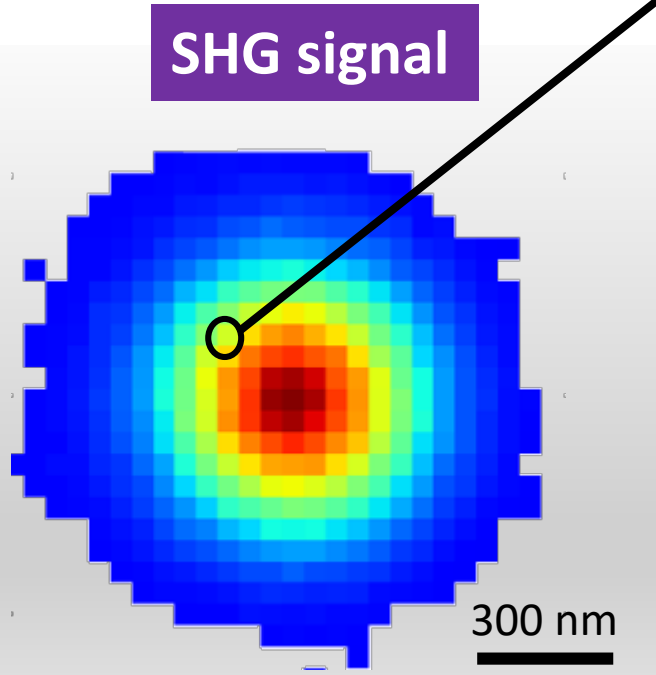
Collaboration with Sophie Brasselet, Inst. Fresnel

Rendon, Timpu, Grange, Brasselet, Scientific report, Accepted, 2018

Super-resolution multiphoton polar microscope



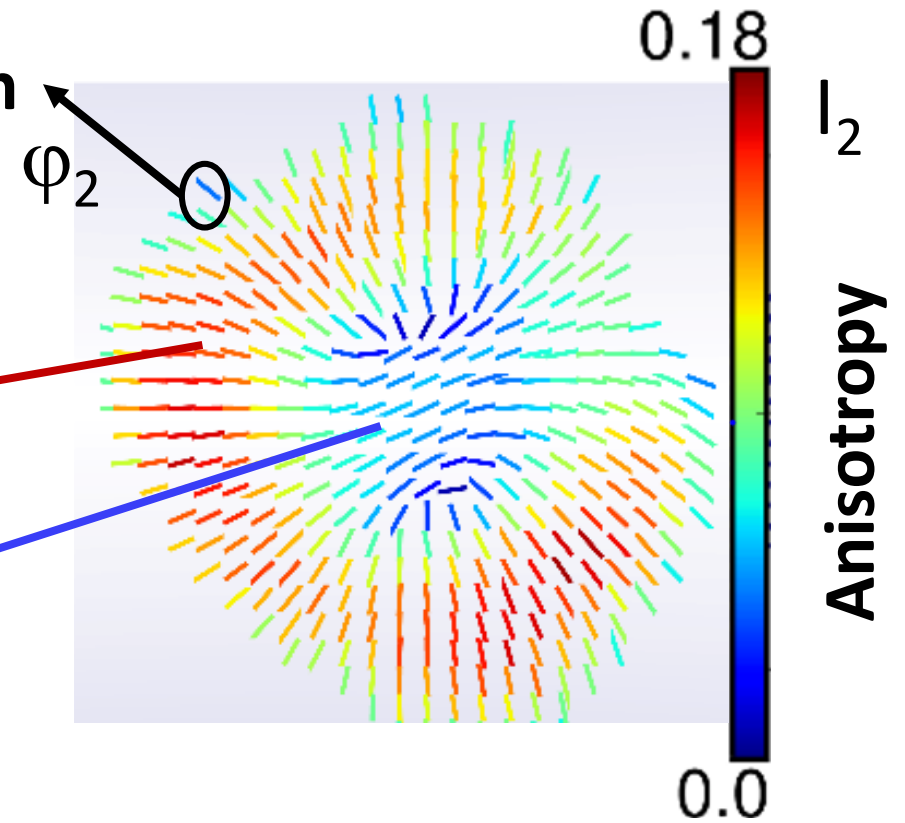
Single pixel analysis



1 pixel = 50 nm

Cubic at the surface

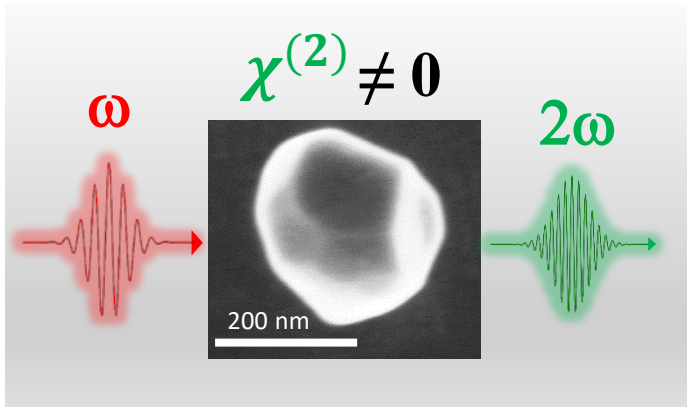
Tetragonal core



Collaboration with Sophie Brasselet, Inst. Fresnel

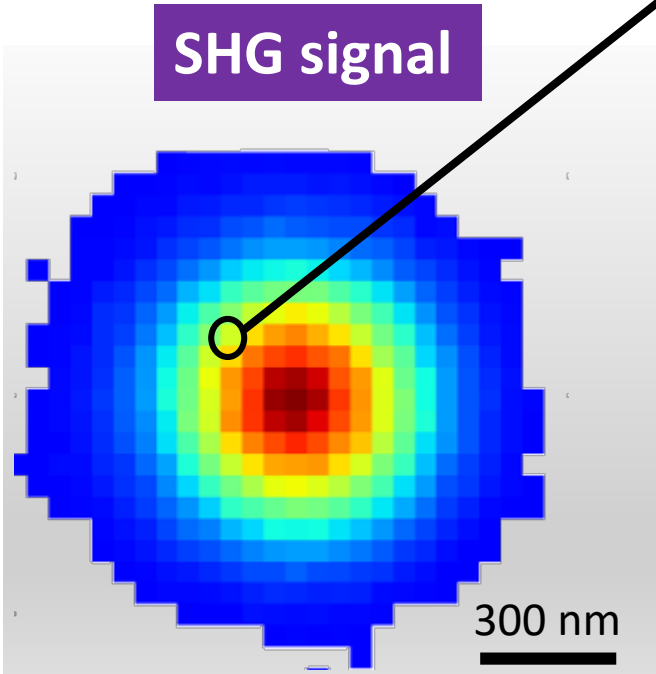
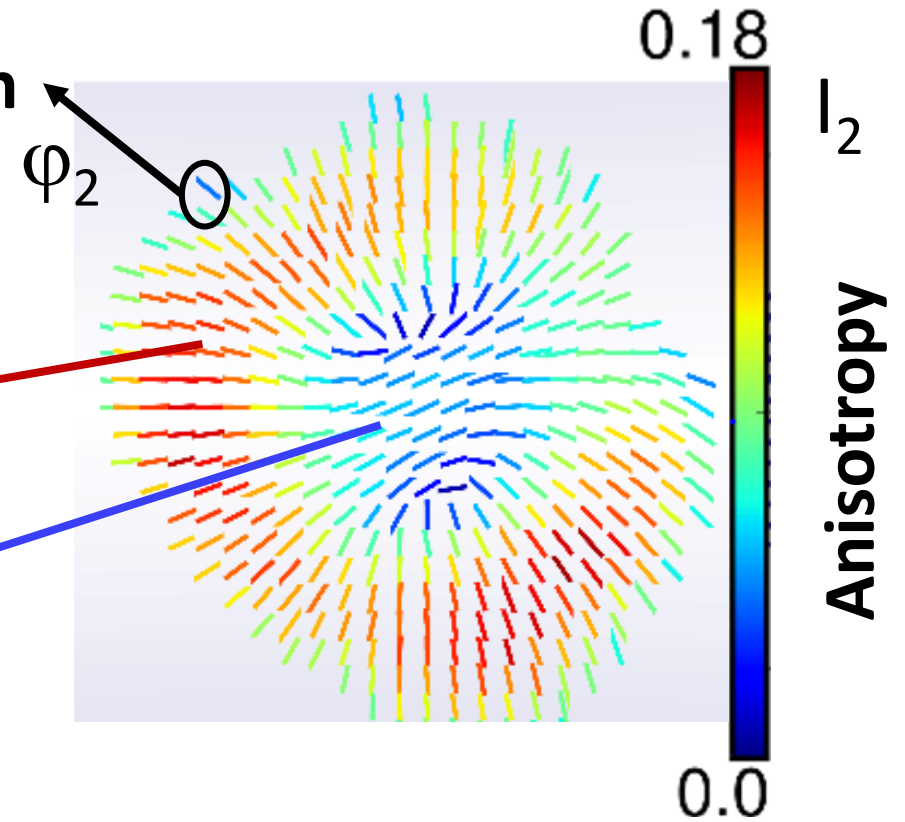
Rendon, Timpu, Grange, Brasselet, Scientific report, Accepted, 2018

Super-resolution multiphoton polar microscope



Single pixel analysis

1 pixel = 50 nm



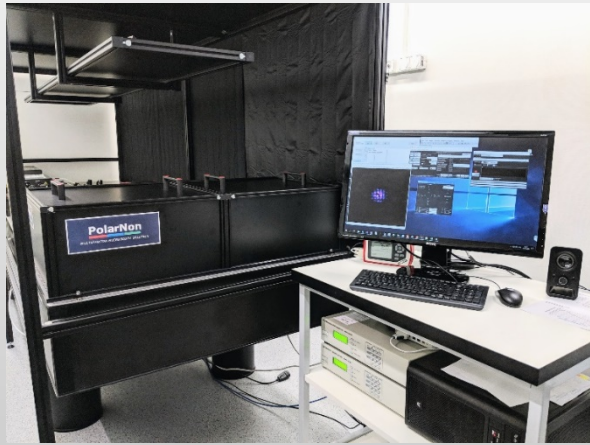
Solve the crystal uncertainty with an optical microscope

Collaboration with Sophie Brasselet, Inst. Fresnel

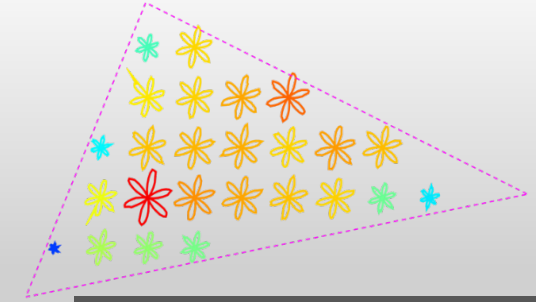
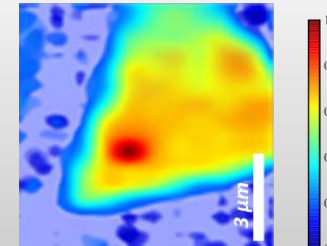
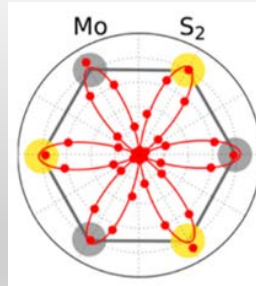
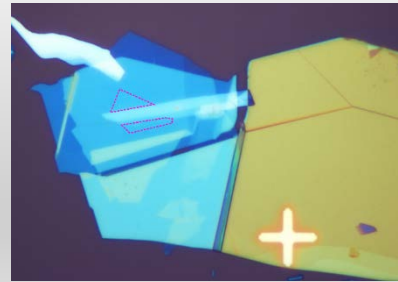
Rendon, Timpu, Grange, Brasselet, Scientific report, Accepted, 2018

Super-resolution polar multiphoton microscope

Super-resolution polar multiphoton microscope

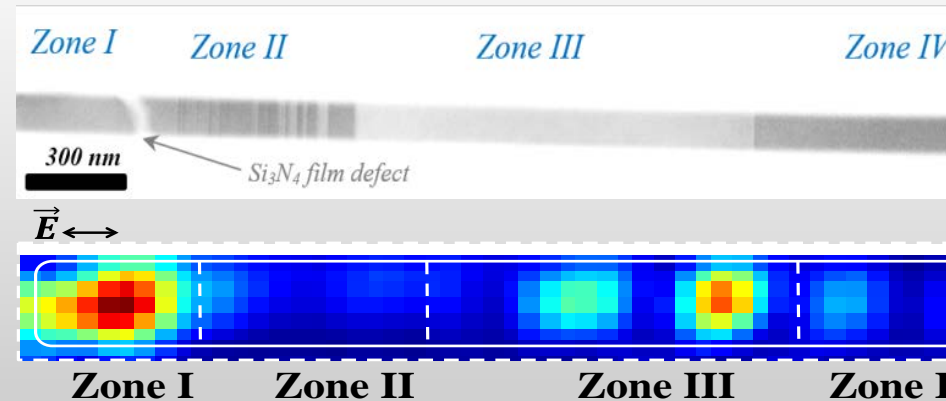


2D materials layers orientation

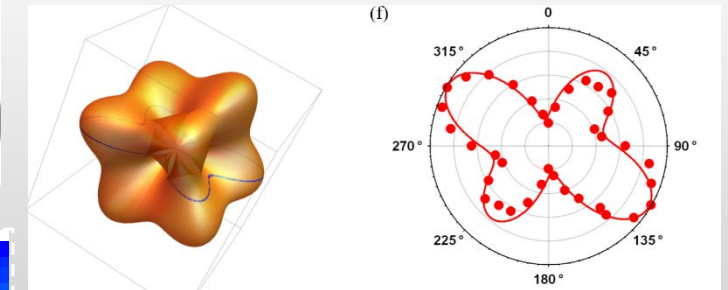


MoS2 sample from K. Ensslin

Crystal structures along III-V nanowires



Zone III

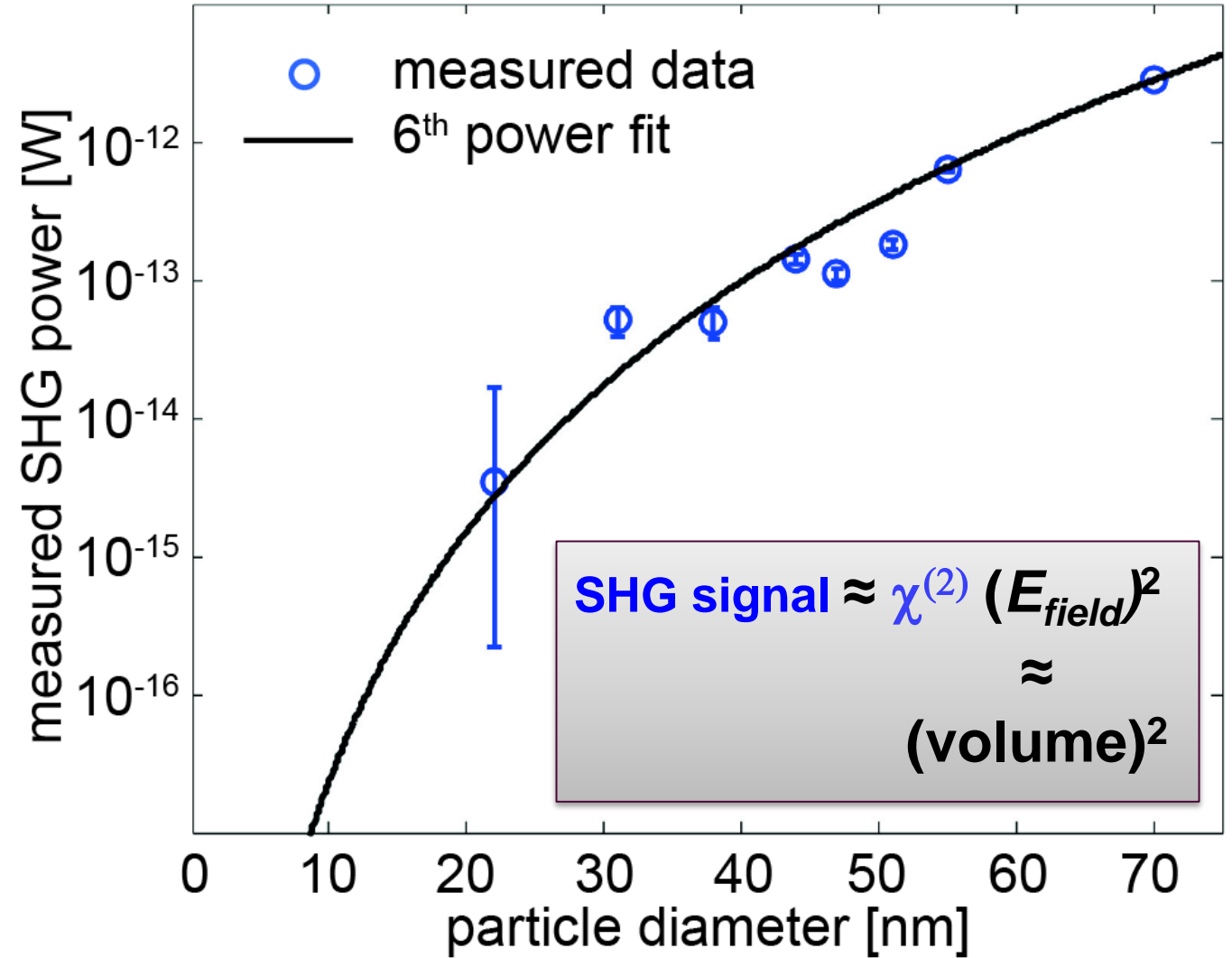
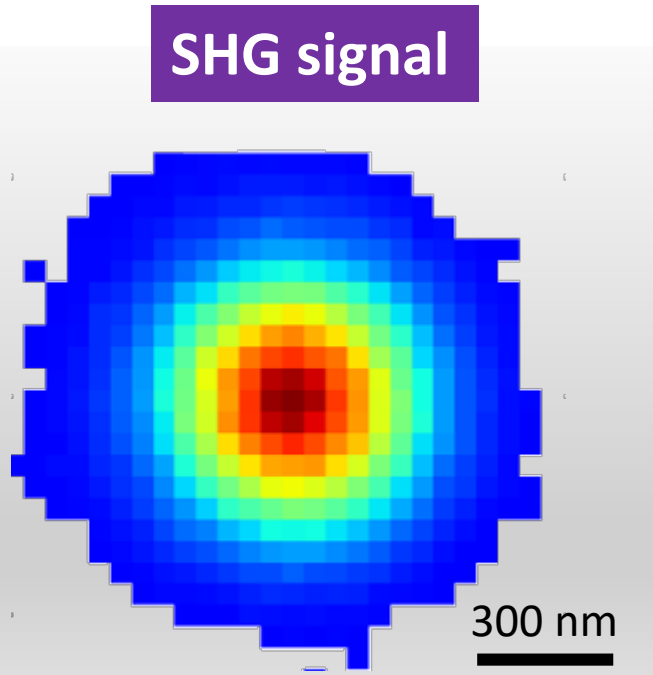
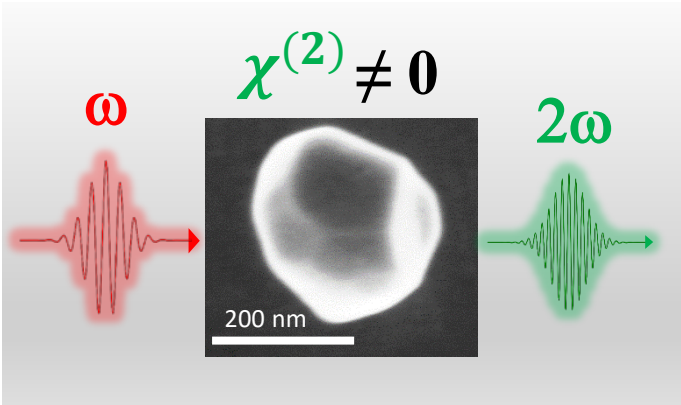


Pure Zinc Blend

Grange, R et al. *Nano letters* 12, 2012

M. Timofeeva, ..., R. Grange *Nano Lett.*, 16, 6290, 2016.

Average SHG intensity of single particles



... room for improvements

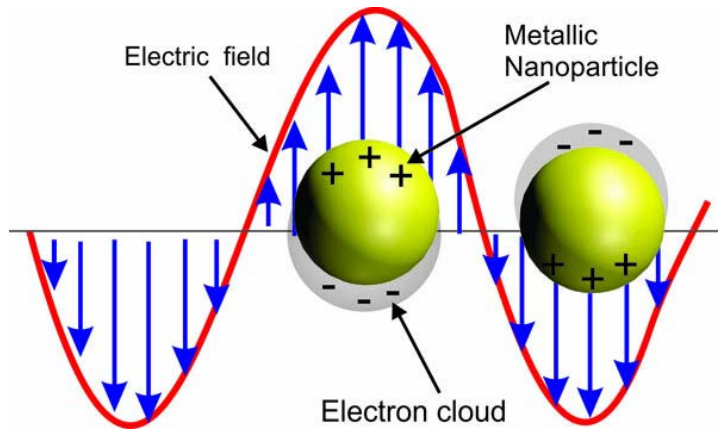
E. Kim,...,and R. Grange, ACS Nano 7, 5343–5349, 2013.
M. Timofeeva,..., R. Grange Nano Lett.,16, 6290, 2016.

Strategies to enhance nonlinear optical signals

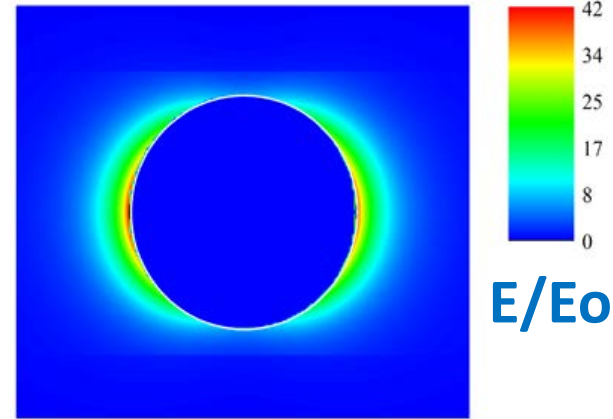
- **with plasmonics**
- with the $\chi^{(2)}$ material itself

What is plasmonics ?

60 nm gold nanoparticles



Conduction electrons can oscillate at the surface of the metal...



...generating strong resonance in the vicinity of the metal...

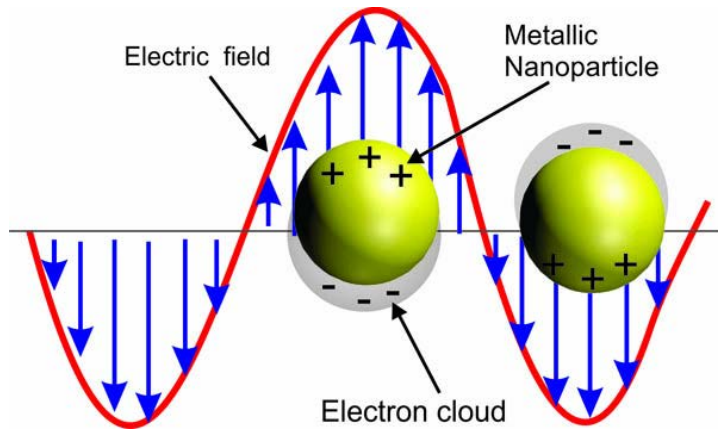


Stained glass, Notre Dame de Paris , 1250

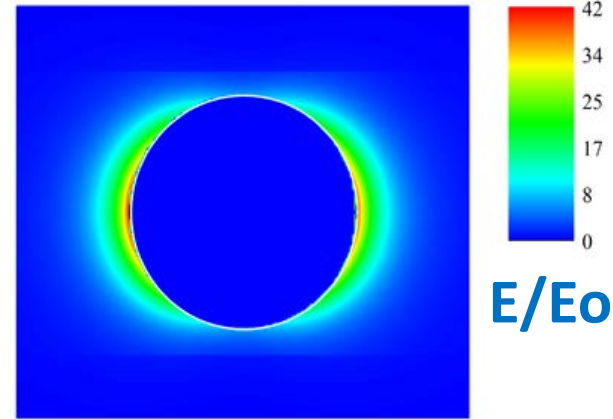
... at visible frequencies

What is plasmonics ?

60 nm gold nanoparticles

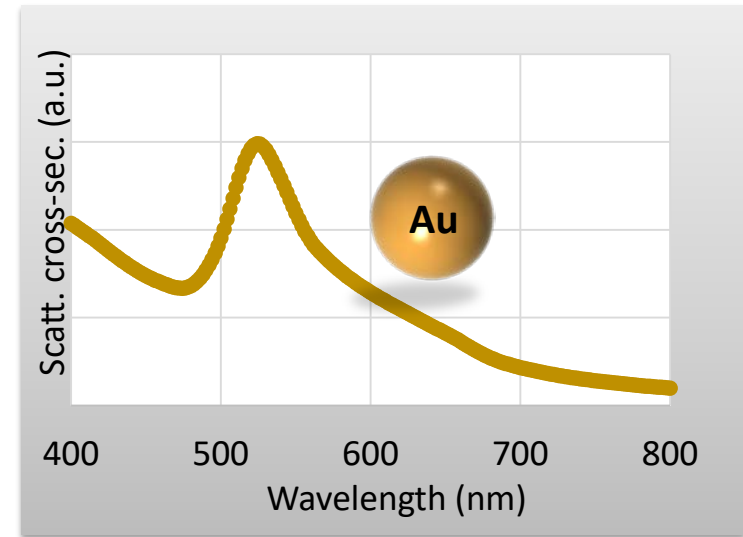


Conduction electrons can oscillate at the surface of the metal...



...generating strong resonance in the vicinity of the metal...

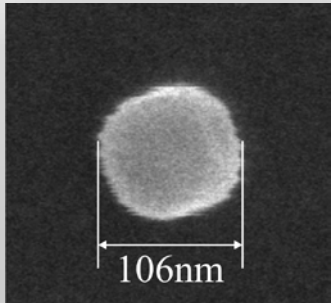
Scattering cross-section



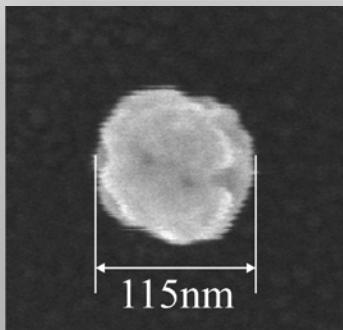
Single electric dipole resonance

$\chi^{(2)}$ material core + gold shell

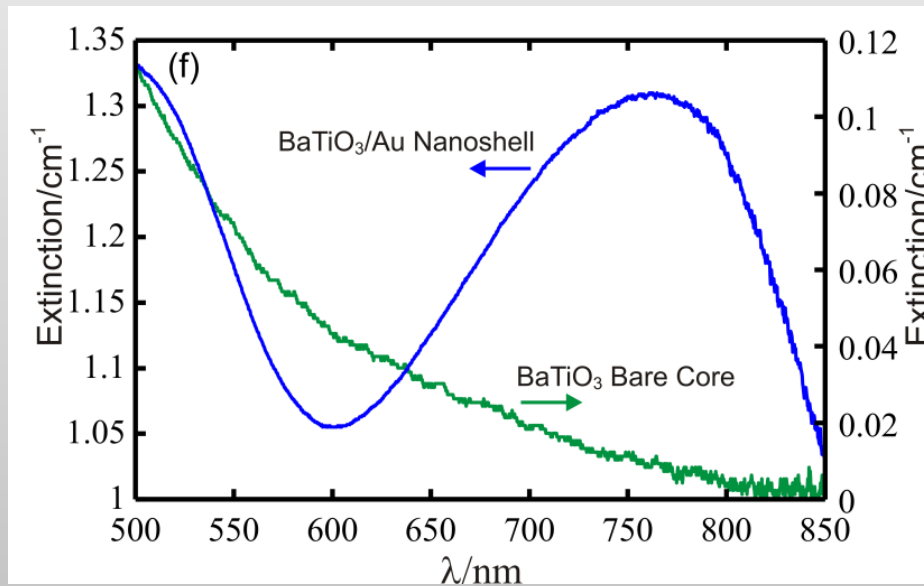
Bare BaTiO₃



Core-shell BaTiO₃

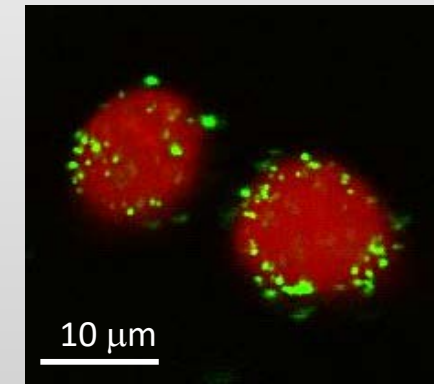


Extinction measurement in suspension



Y. Pu, R. Grange, Ch.-L. Hsieh, and D. Psaltis, Phys. Rev. Lett., 104, 2010.

Bioimaging

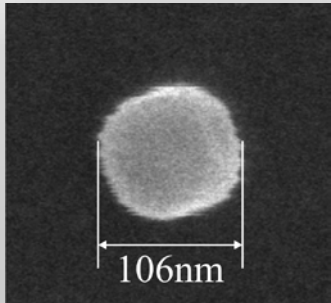


Hsieh, Grange et al, Opt. Exp. 2009
Hsieh, Grange et al, Biomaterials, 2010
Grange et al., Biomedical Opt Exp, 2011

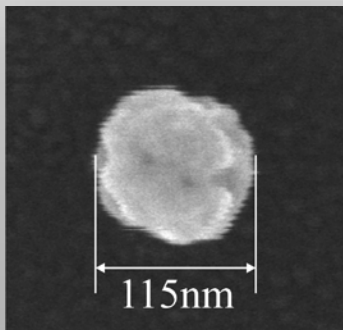
SHG enhancement factor > 500

$\chi^{(2)}$ material core + gold shell

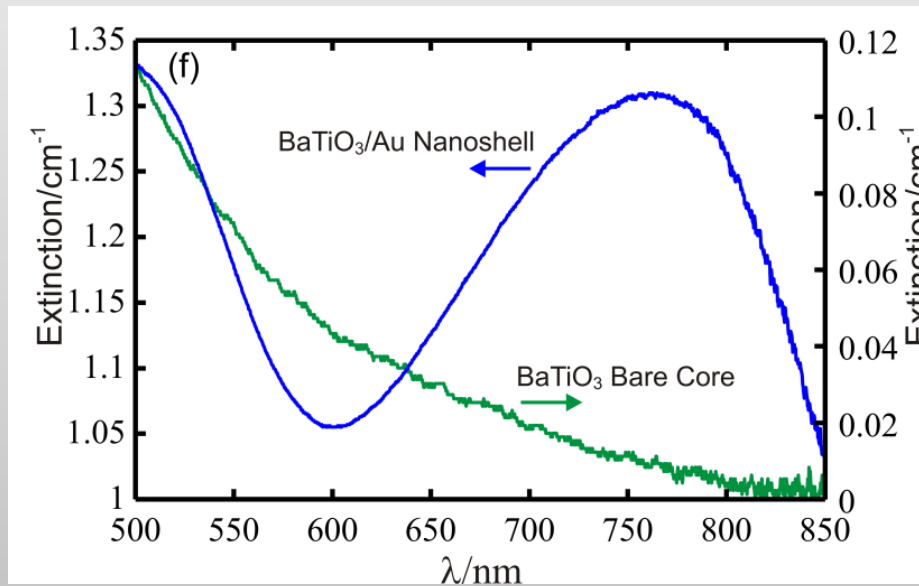
Bare BaTiO₃



Core-shell BaTiO₃

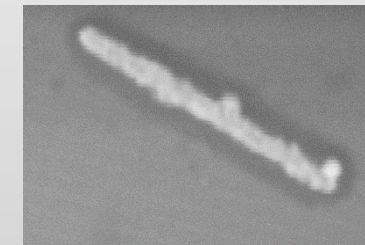


Extinction measurement in suspension



Y. Pu, R. Grange, Ch.-L. Hsieh, and D. Psaltis, Phys. Rev. Lett., 104, 2010.

KNbO₃ nanowires

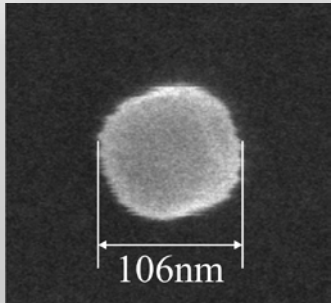


J. Richter, ..., R. Grange
Nanoscale, 6, 5200, 2014.

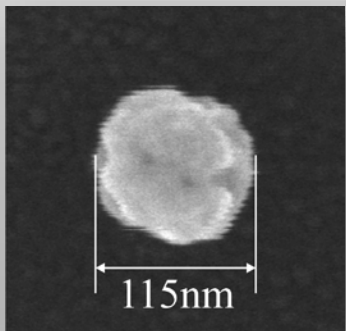
SHG enhancement factor > 250

$\chi^{(2)}$ material core + gold shell

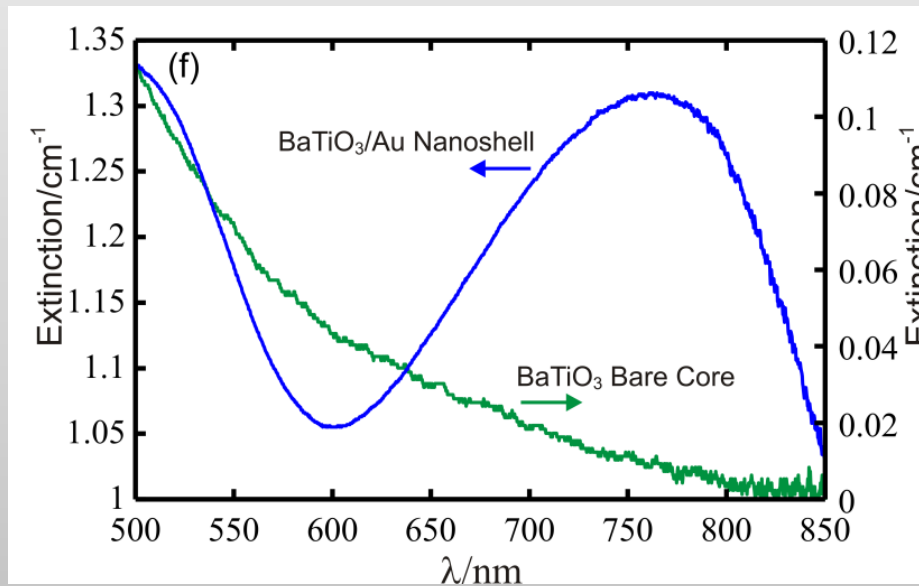
Bare BaTiO₃



Core-shell BaTiO₃

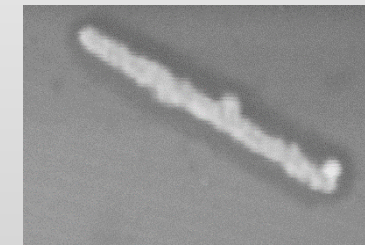


Extinction measurement in suspension



Y. Pu, R. Grange, Ch.-L. Hsieh, and D. Psaltis, Phys. Rev. Lett., 104, 2010.

KNbO₃ nanowires



J. Richter, ..., R. Grange
Nanoscale, 6, 5200, 2014.

SHG enhancement factor > 250

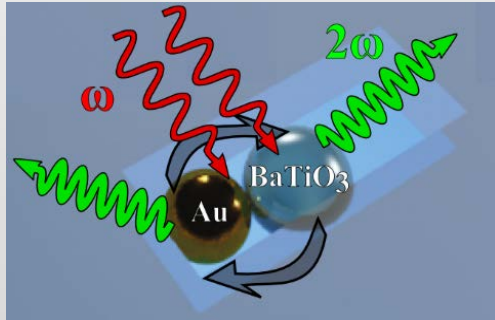
Markers as bright as quantum dots, no bleaching, and biocompatible

Combination of plasmonics with

$\chi^{(2)}$ materials

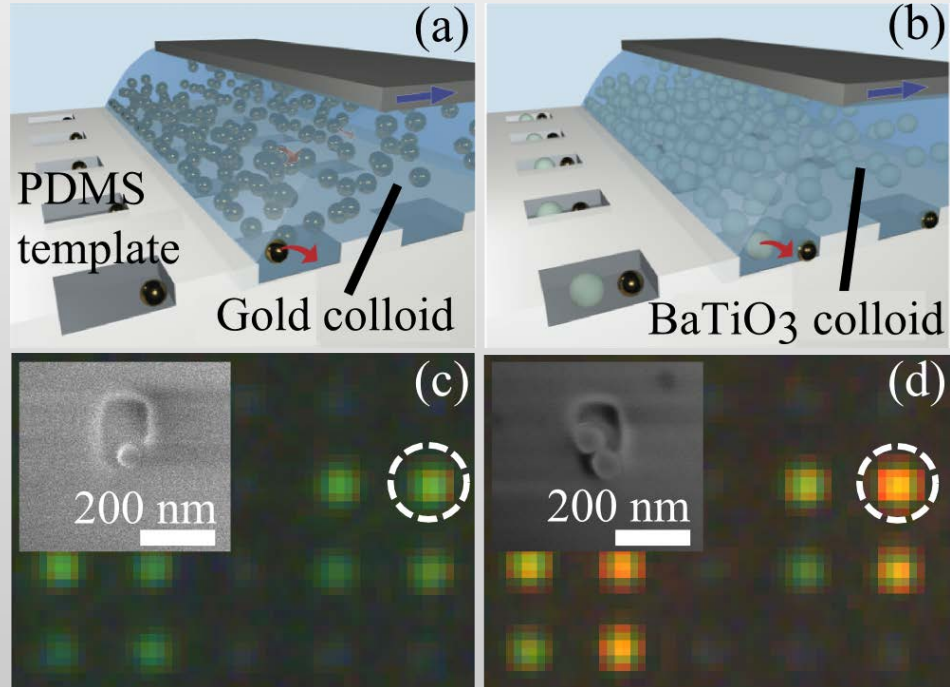
Sequential Capillarity-Assisted Particle Assembly

Hybrid Nanodimers



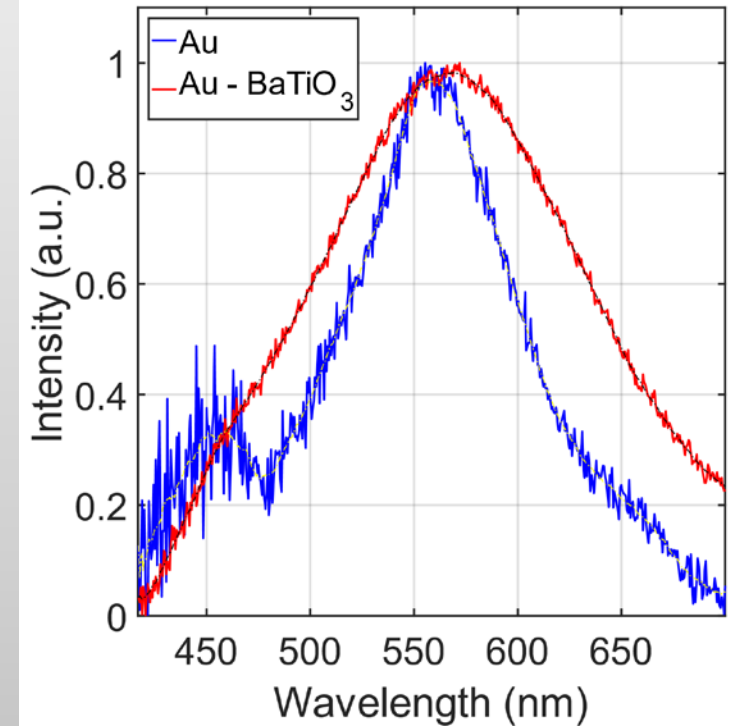
Both 100 nm \varnothing

Collaboration with Lucio Isa (ETH) and Heiko Wolf (IBM)



Timpu, F.; Hendricks, N. R.; Petrov, M.; Ni, S.; Renaut, C.; Wolf, H.; Isa, L.; Kivshar, Y.; Grange, R. *Nano Letters*, 17, 5381, 2017.

Extinction

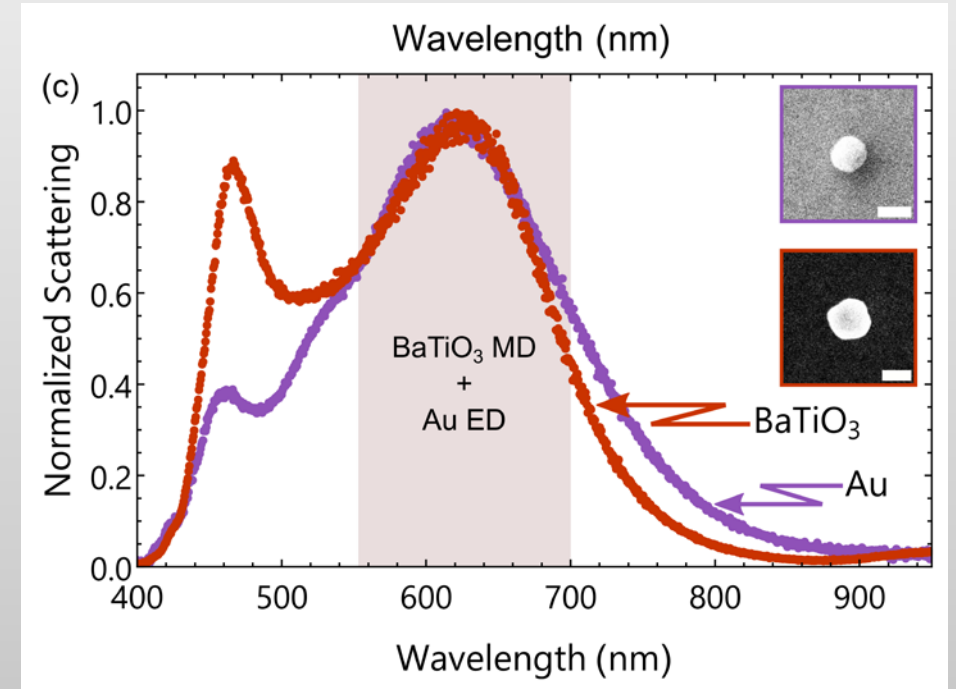
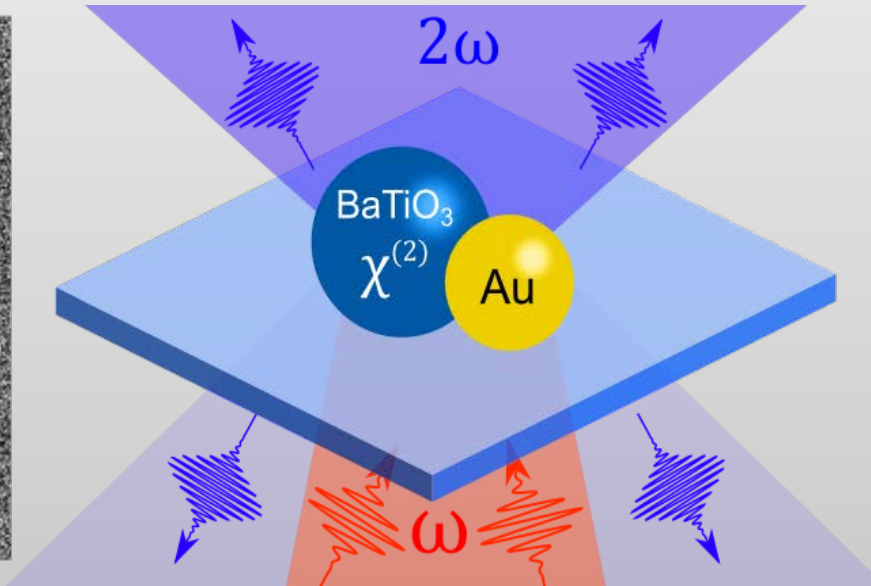
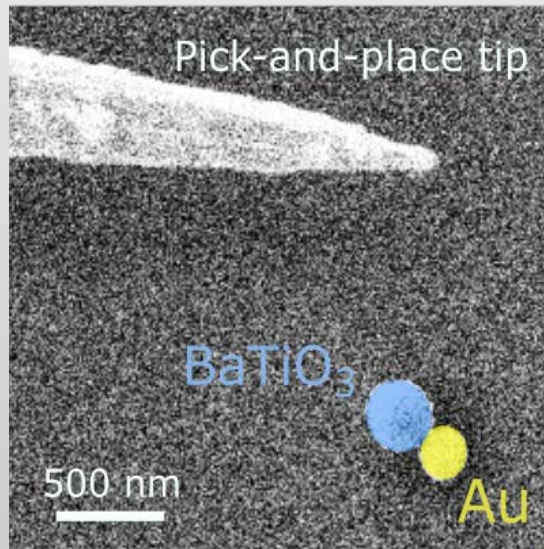


SHG enhancement : 5-15

High-throughput and reduced fabrication complexity

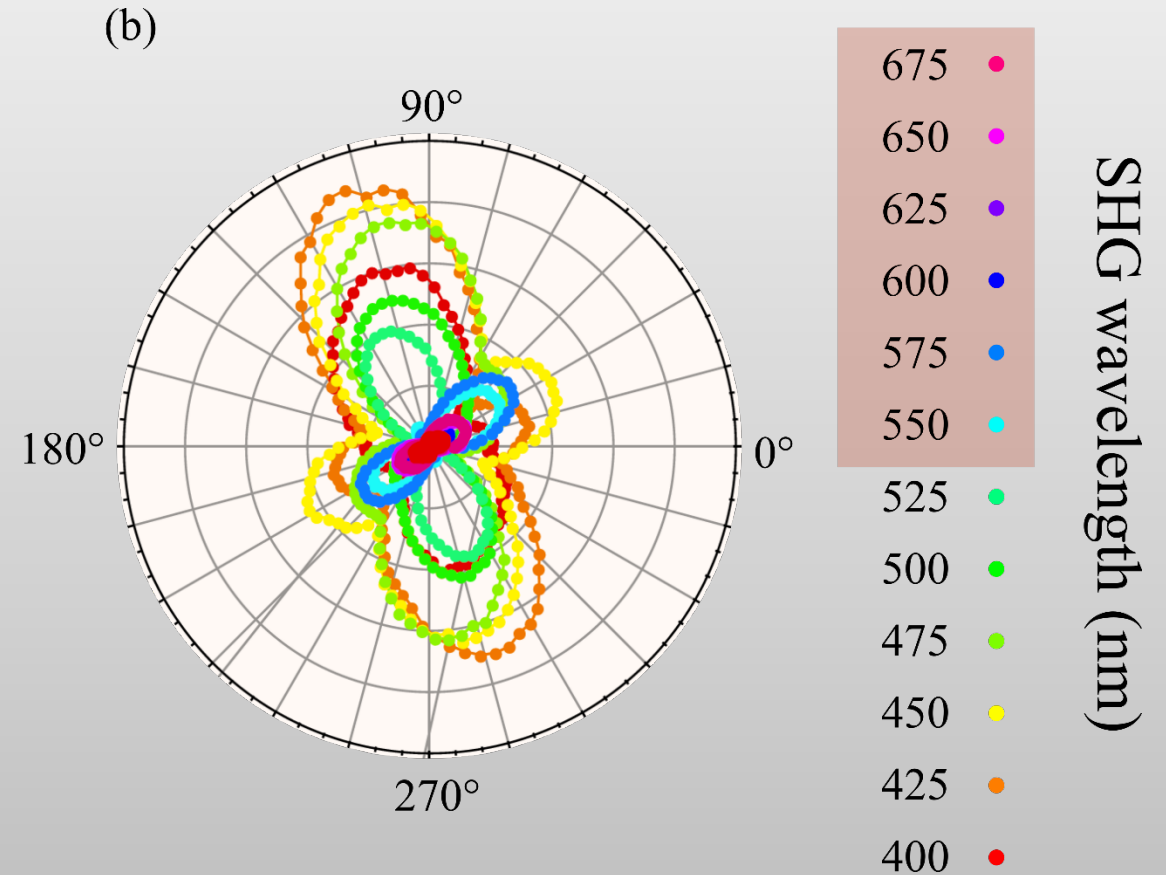
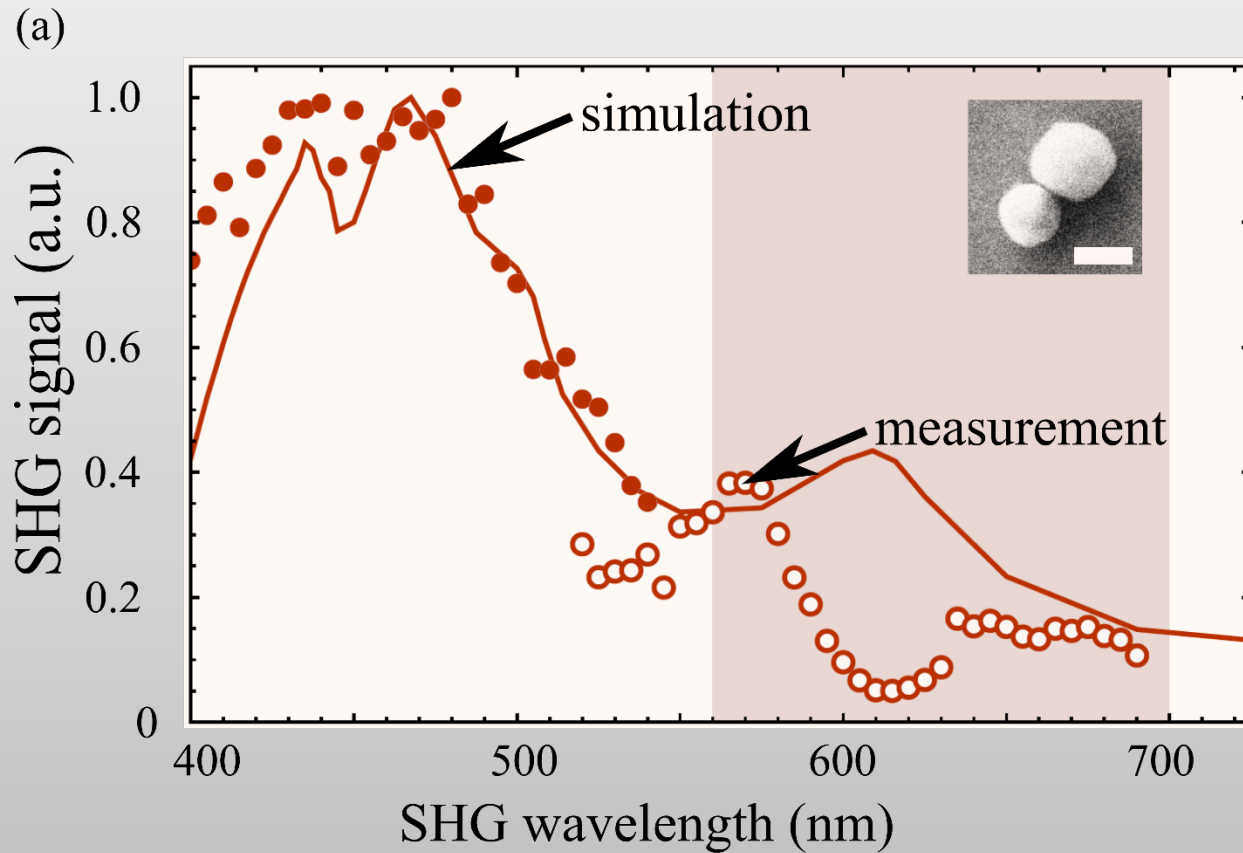
Reshaping the Second-Order Polar Response

Hybrid Nanodimers



Renaut et al., *Nano Letters*, just accepted, 2019

Reshaping the Second-Order Polar Response



Renaut et al., *Nano Letters*, just accepted, 2019

Drawback of metals : thermal losses

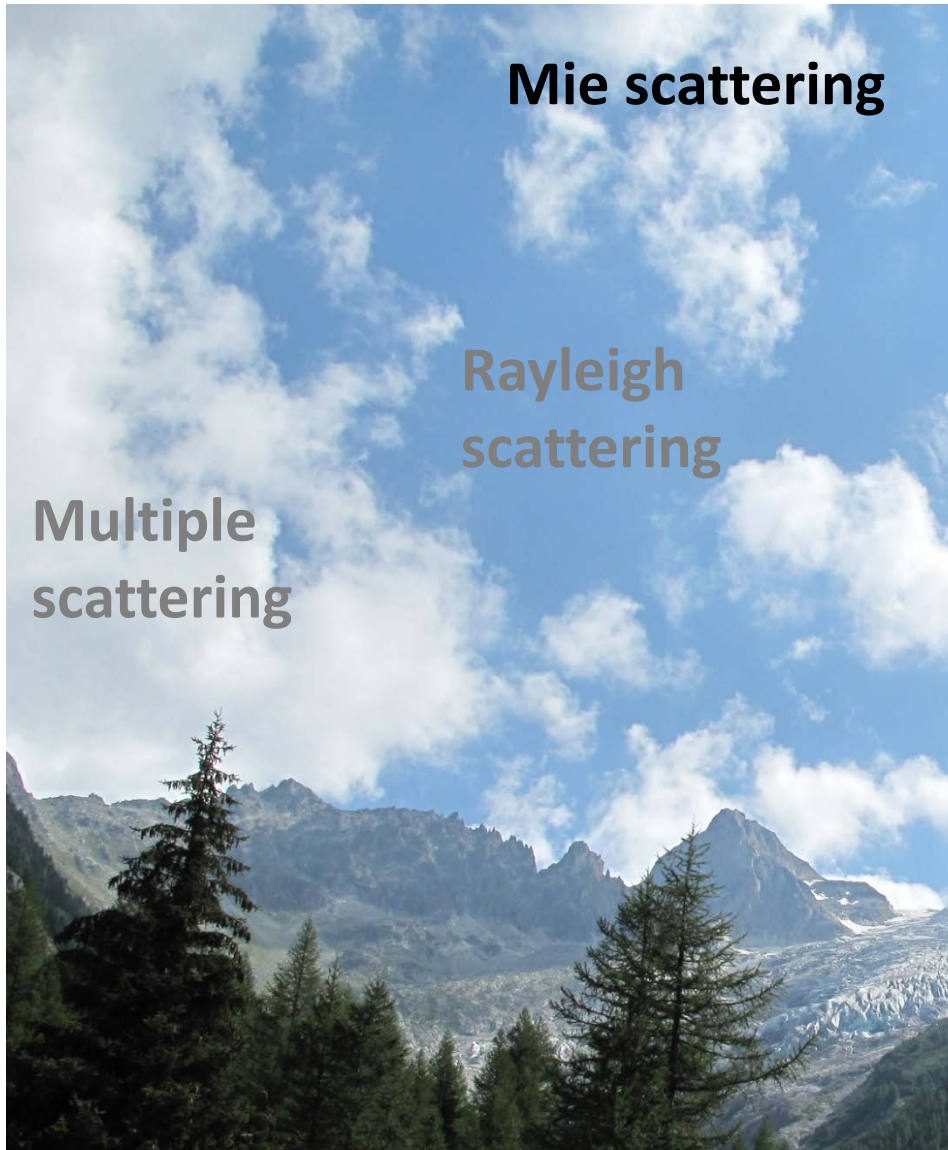


J. B. Khurgin, How to deal with loss in plasmonics and metamaterials, Nature Nanotechnology, 10, 2015.

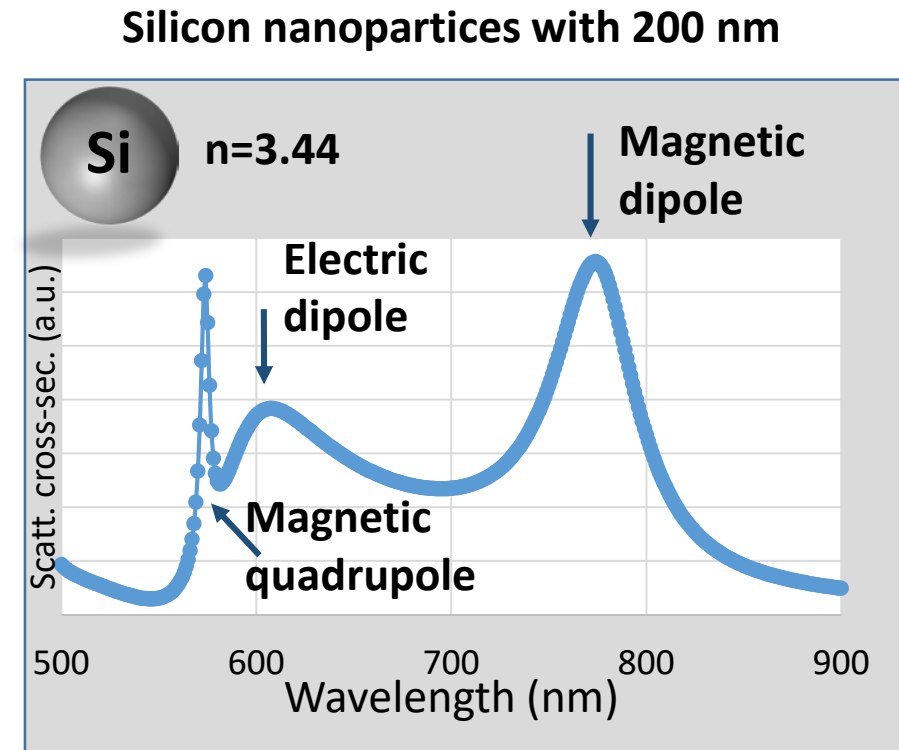
Strategies to enhance nonlinear optical signals

- with plasmonics
- with the $\chi^{(2)}$ material itself

Mie resonances of the $\chi^{(2)}$ material itself



Mie resonances in material with $n > 2$ and size $\sim \lambda$:
electric and magnetic resonance have comparable strength

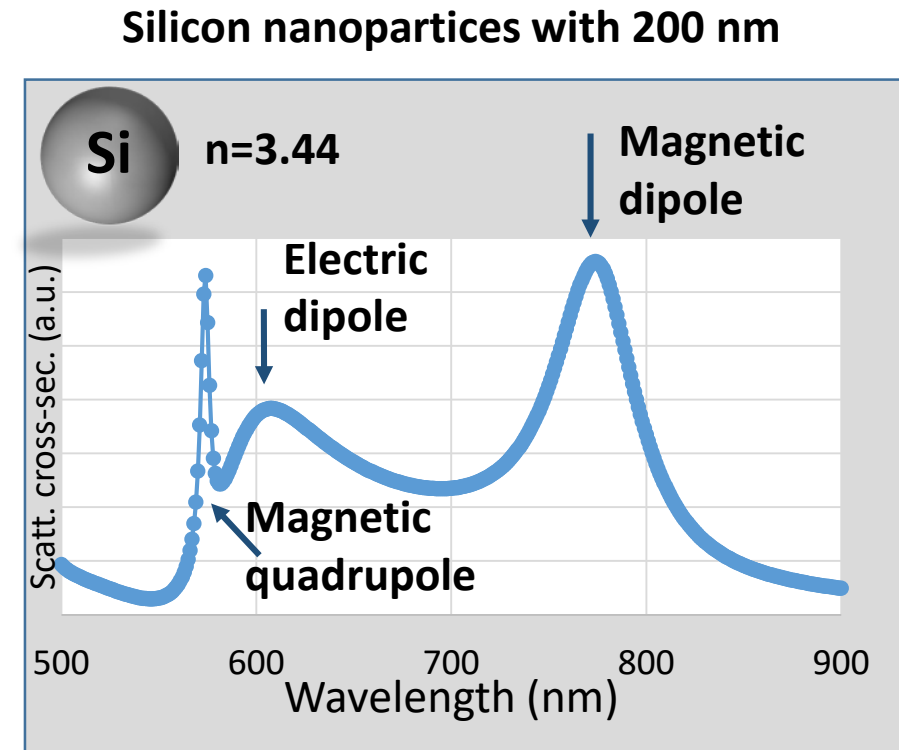
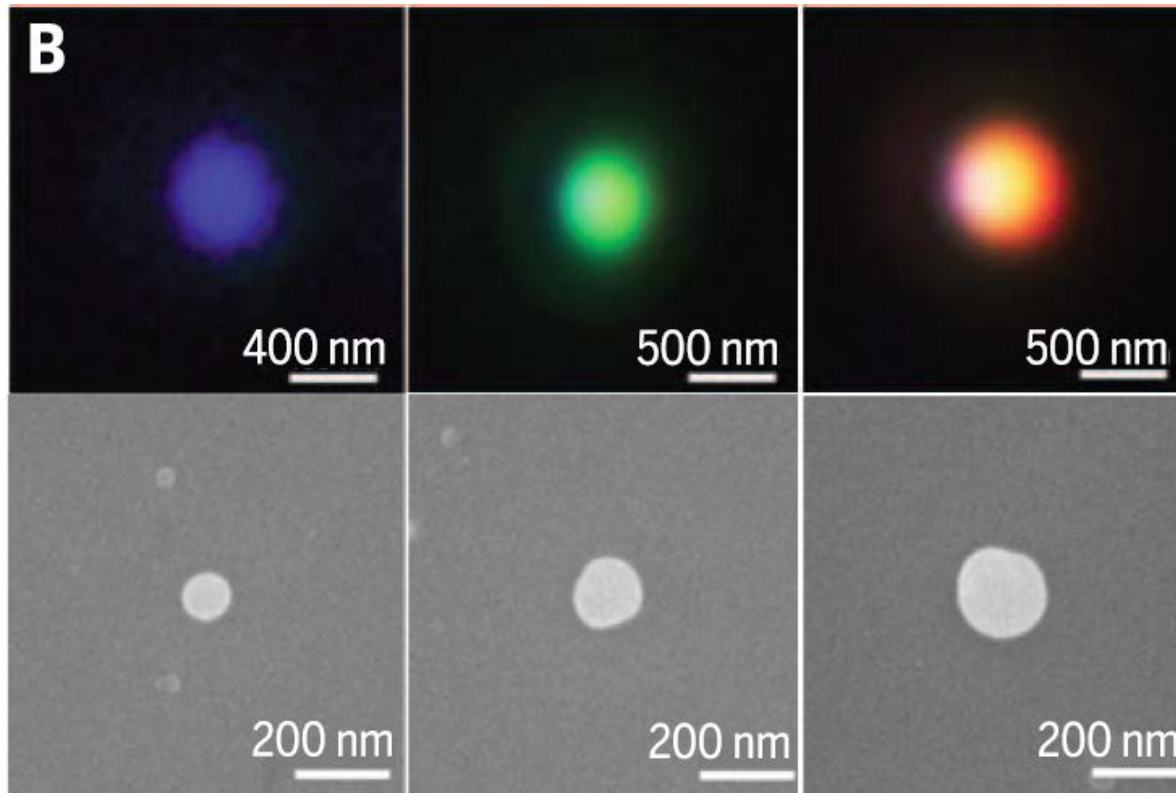


- A. García-Etxarri et al., *Opt. Express*, 19, (2011)
- Krasnok, A. et al., *Opt. Express* 20, 20599 (2012)
- A. B. Evlyukhin et al. , *Nano Lett.* 12,3749, (2012)
- Zywiets, U. et al. *Nat. Commun.* 5:3402 (2014)

Mie resonances of the $\chi^{(2)}$ material itself

Mie resonances in material with $n > 2$ and size $\sim \lambda$:
electric and magnetic resonance have comparable strength

Colors vary with sizes of the particle



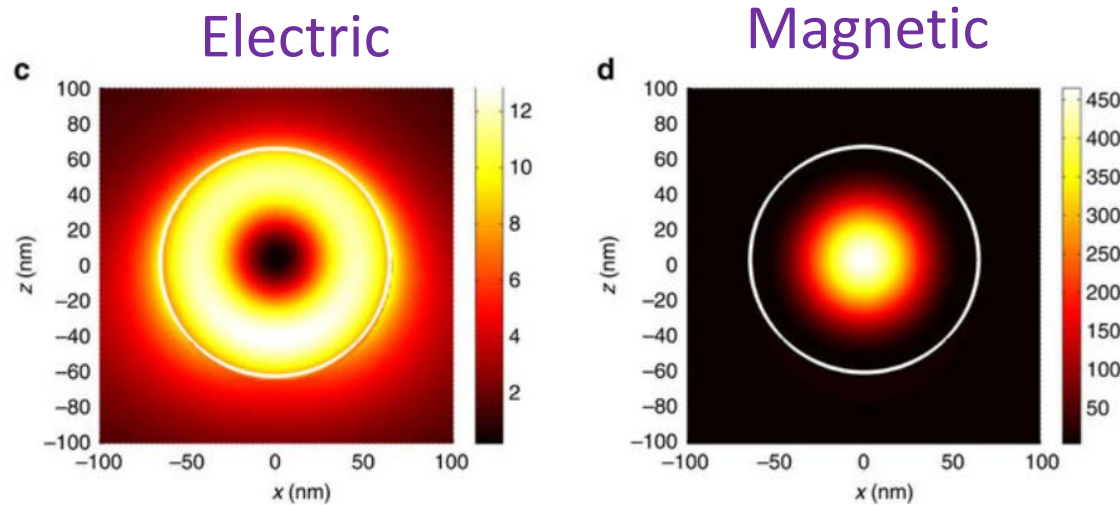
Arseniy I. Kuznetsov et al. Science 2016, 354, 2472

- A. García-Etxarri et al., *Opt. Express*, 19, (2011)
- Krasnok, A. et al., *Opt. Express* 20, 20599 (2012)
- A. B. Evlyukhin et al. , *Nano Lett.* 12,3749, (2012)
- Zywiets, U. et al. *Nat. Commun.* 5:3402 (2014)

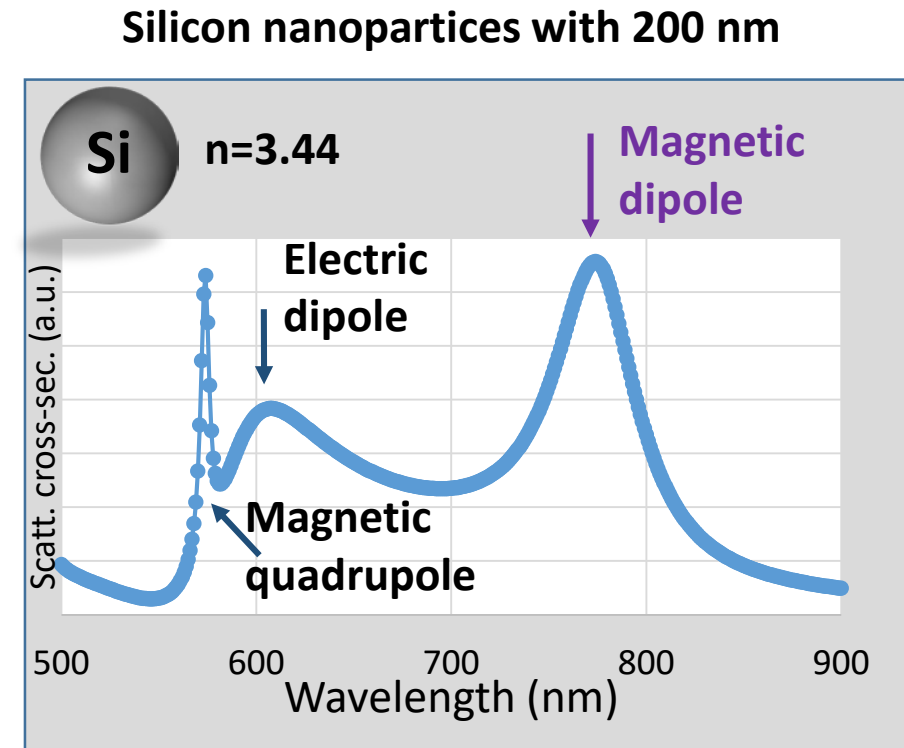
Mie resonances of the $\chi^{(2)}$ material itself

Mie resonances in material with $n > 2$ and size $\sim \lambda$:
electric and magnetic resonance have comparable strength

Calculated field intensities at the
magnetic dipole resonances



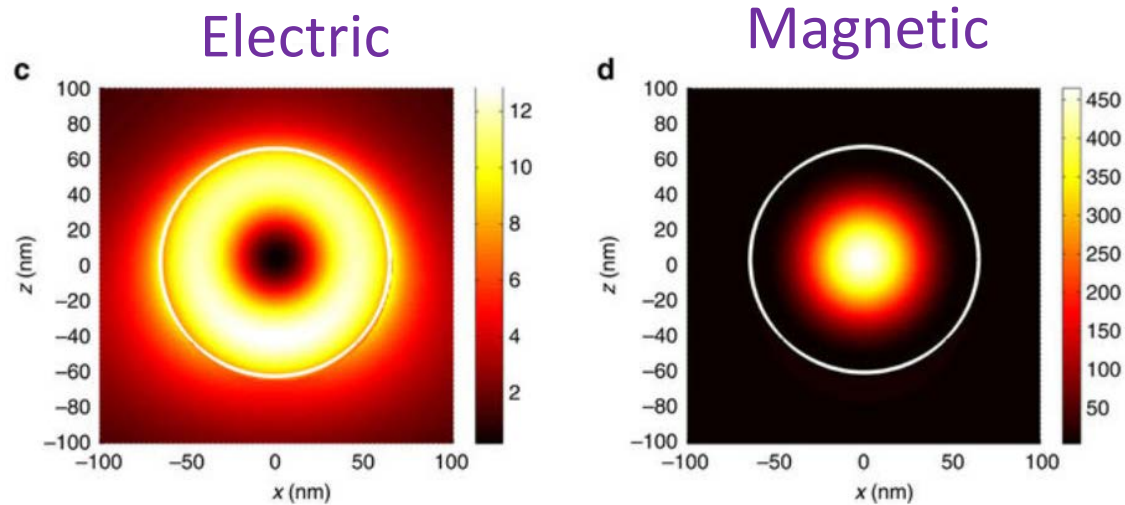
'Dielectric' material: electromagnetic field amplified in the volume of the nanoparticle



- A. García-Etxarri et al., *Opt. Express*, 19, (2011)
- Krasnok, A. et al., *Opt. Express* 20, 20599 (2012)
- A. B. Evlyukhin et al., *Nano Lett.* 12,3749, (2012)
- Zywietz, U. et al. *Nat. Commun.* 5:3402 (2014)

Mie resonances of the $\chi^{(2)}$ material itself

Calculated field intensities at the
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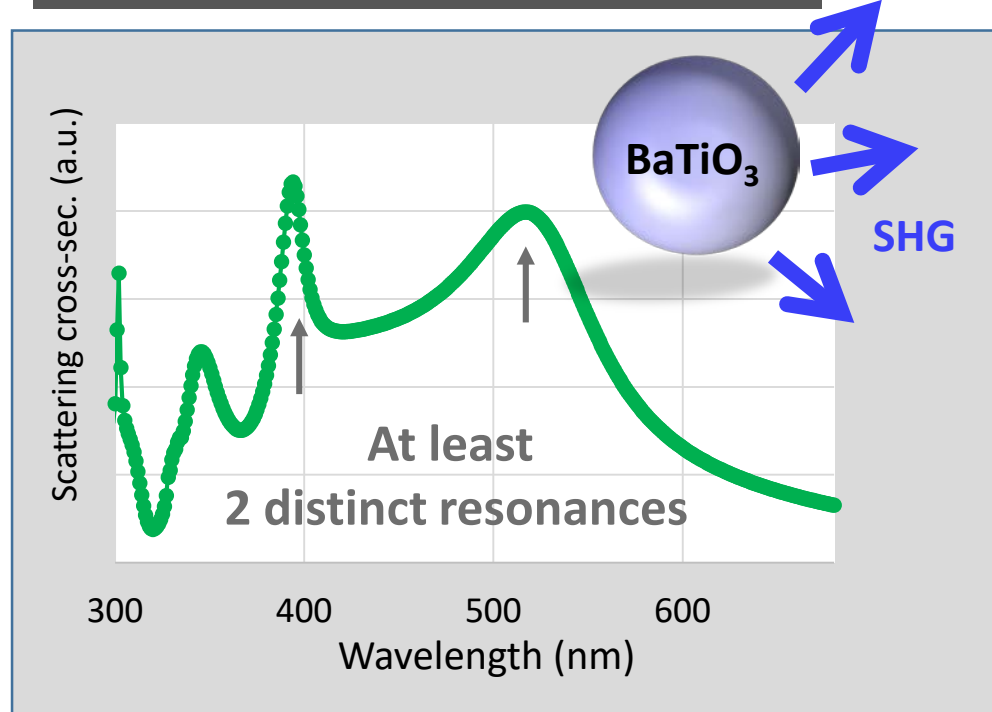


‘Dielectric’ material: electromagnetic field amplified in the volume of the nanoparticle

Material	$\chi^{(2)}$ coeff. (pm/V)
Silicon	0
SiO ₂ (Quartz)	0.335
BaTiO ₃	6.8 to 17
LiNbO ₃	6 to 34
GaAs	134 to 256

Mie resonances of the $\chi^{(2)}$ material itself

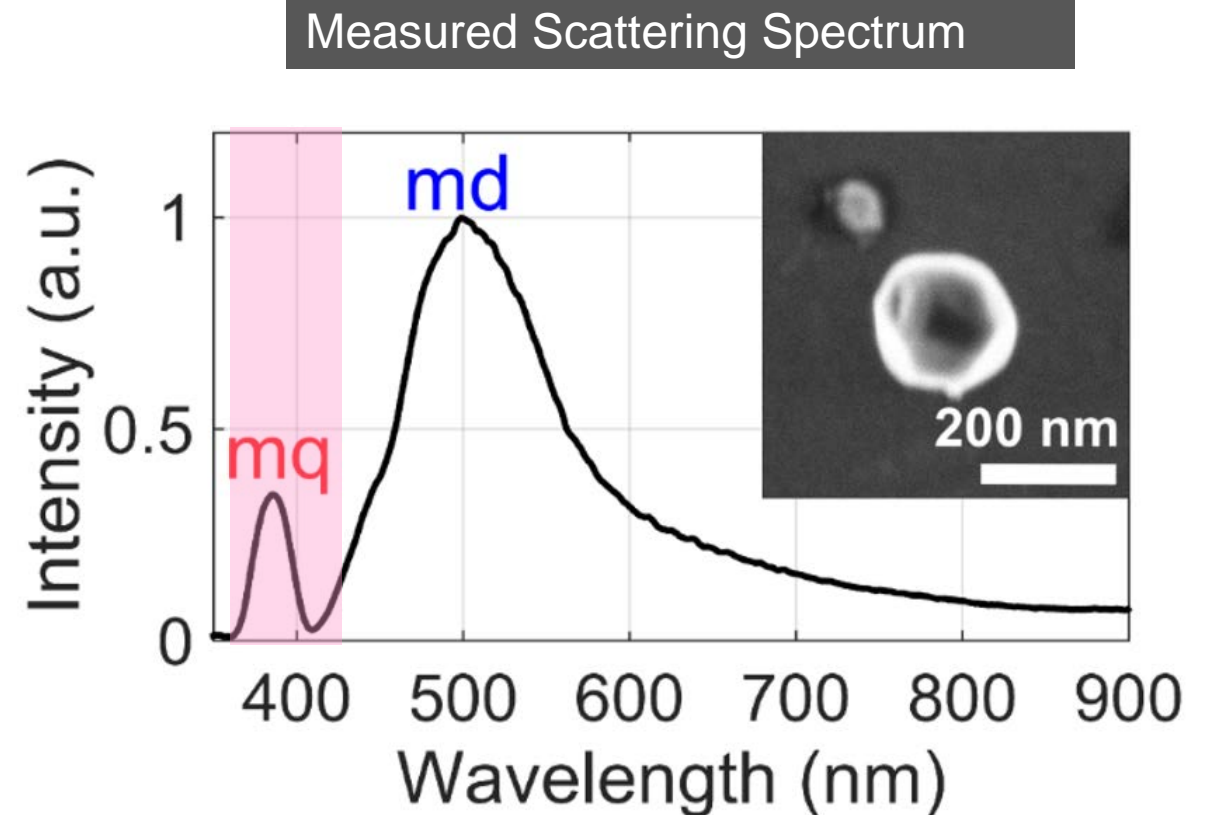
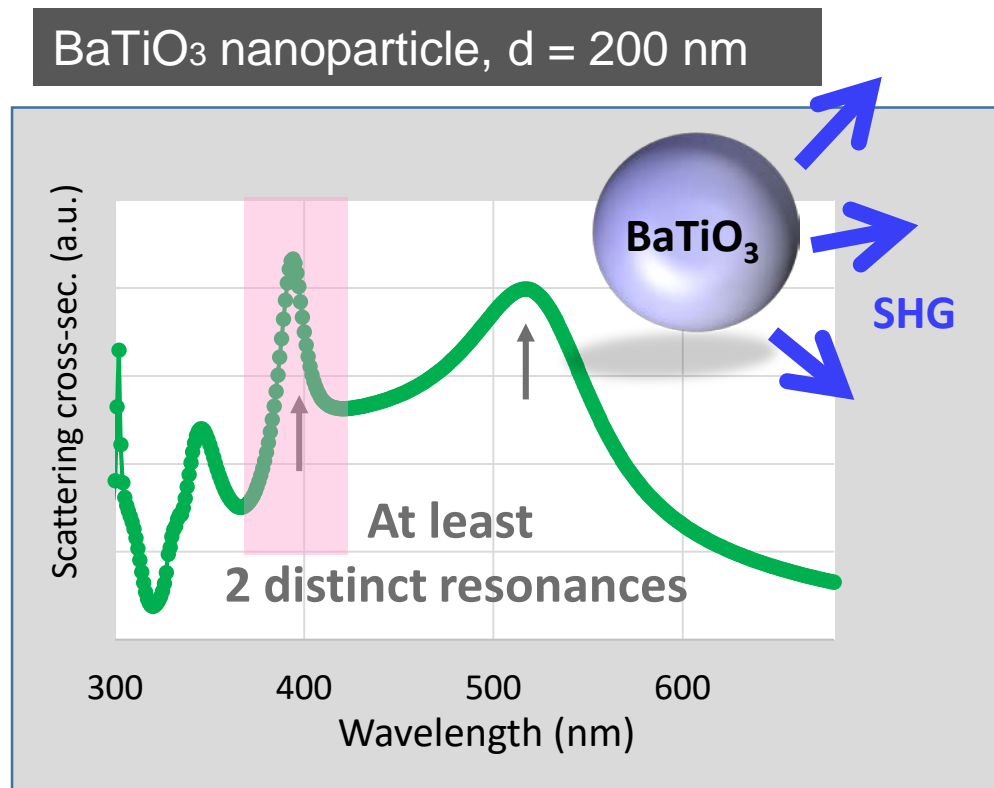
BaTiO₃ nanoparticle, d = 200 nm



Material	$\chi^{(2)}$ coeff. (pm/V)
Silicon	0
SiO ₂ (Quartz)	0.335
BaTiO ₃	6.8 to 17
LiNbO ₃	6 to 34
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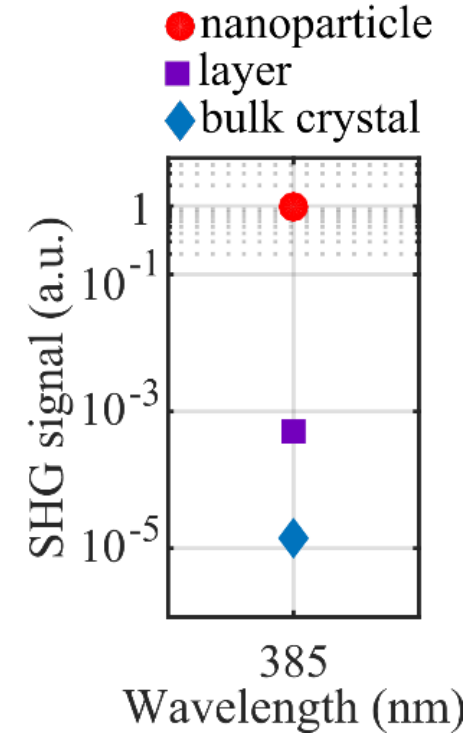
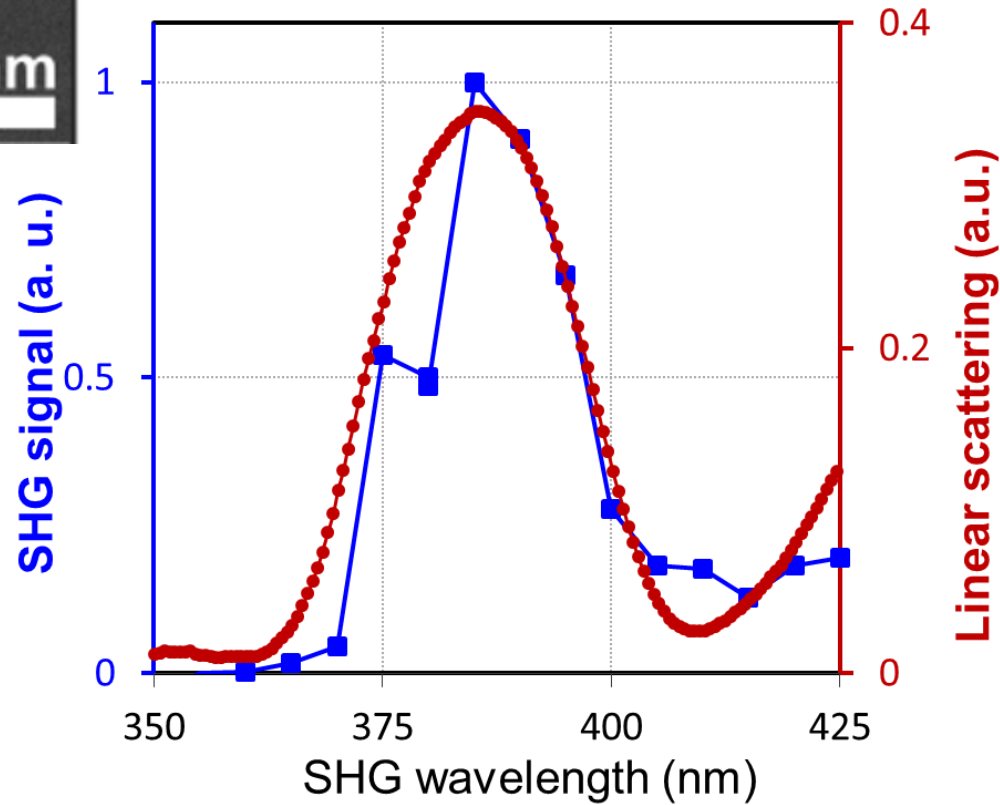
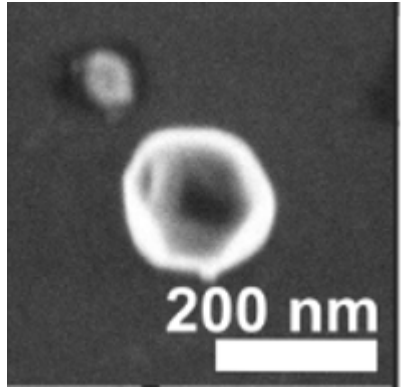
F. Timpu, A. Sergeyev, N. Hendricks, and R. Grange, ACS Photonics, 4, 2017.

Mie resonances of the $\chi^{(2)}$ material itself



F. Timpu, A. Sergeyev, N. Hendricks, and R. Grange, ACS Photonics, 4, 2017.

Mie resonances enhance SHG from single nanoparticle



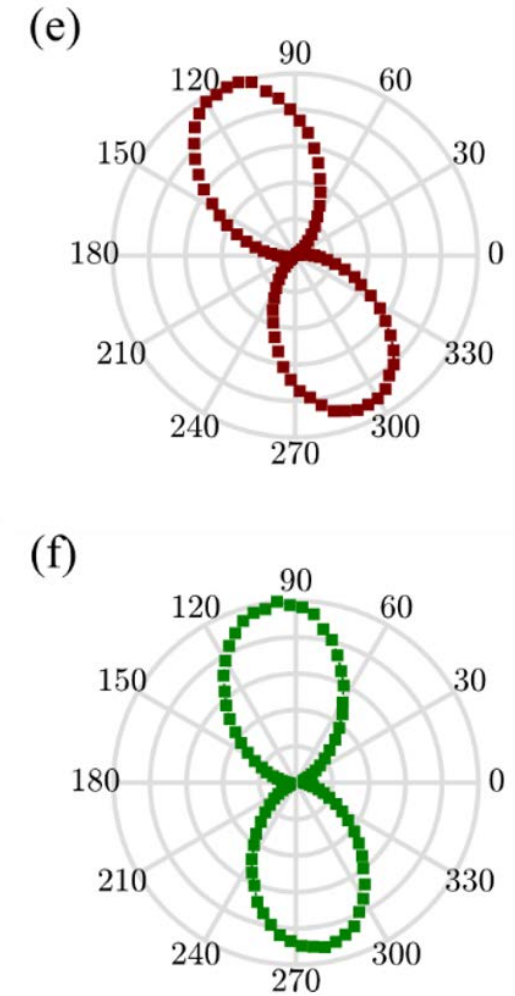
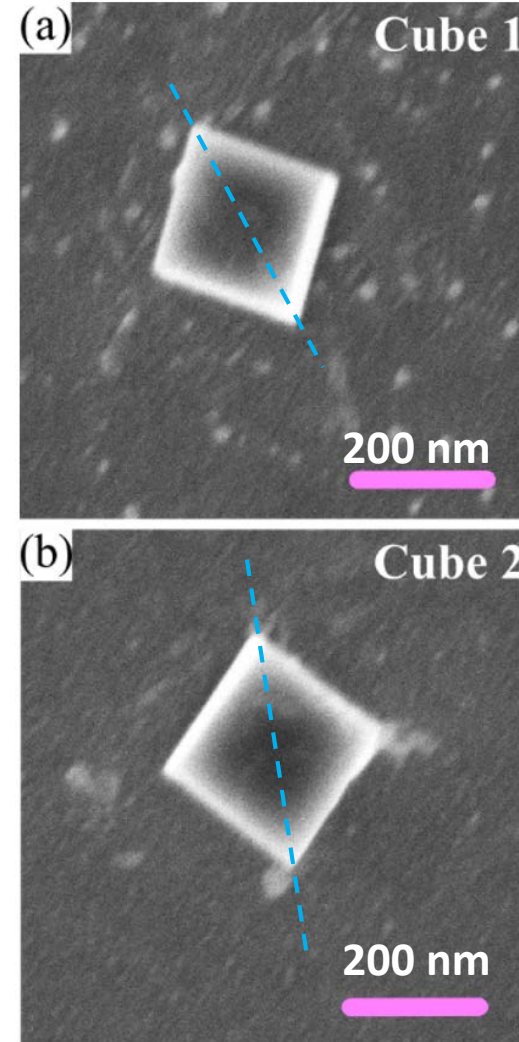
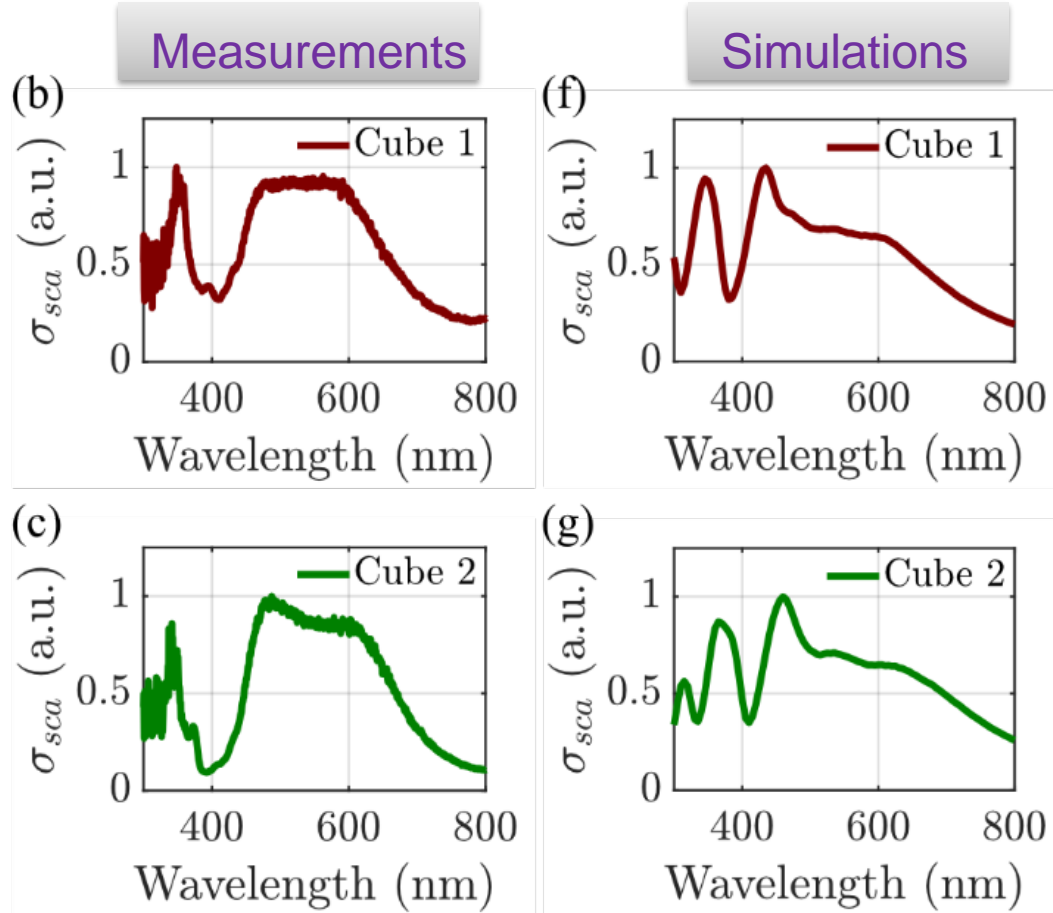
**10⁴
enhancement
with a single
resonance !**

F. Timpu, A. Sergeev, N. Hendricks, and R. Grange, ACS Photonics, 4, 2017.

Mie resonances in Lithium Niobate Nanocubes

Linear scattering

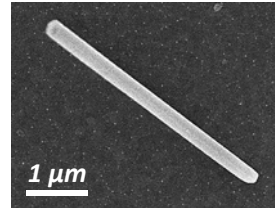
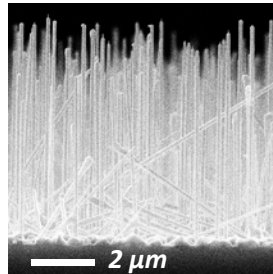
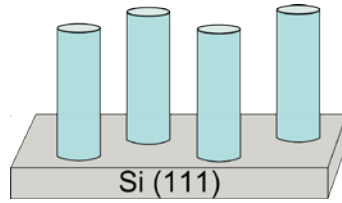
SHG Polar response



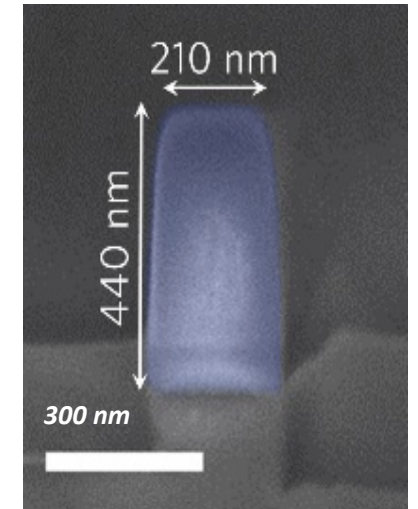
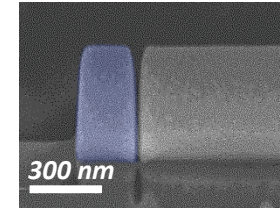
Mie resonances in III-V nanowires

Anapoles in Free-Standing III-V Nanodisks

$\text{Al}_{0.3}\text{Ga}_{0.7}\text{As}$
nanowires
grown on Si



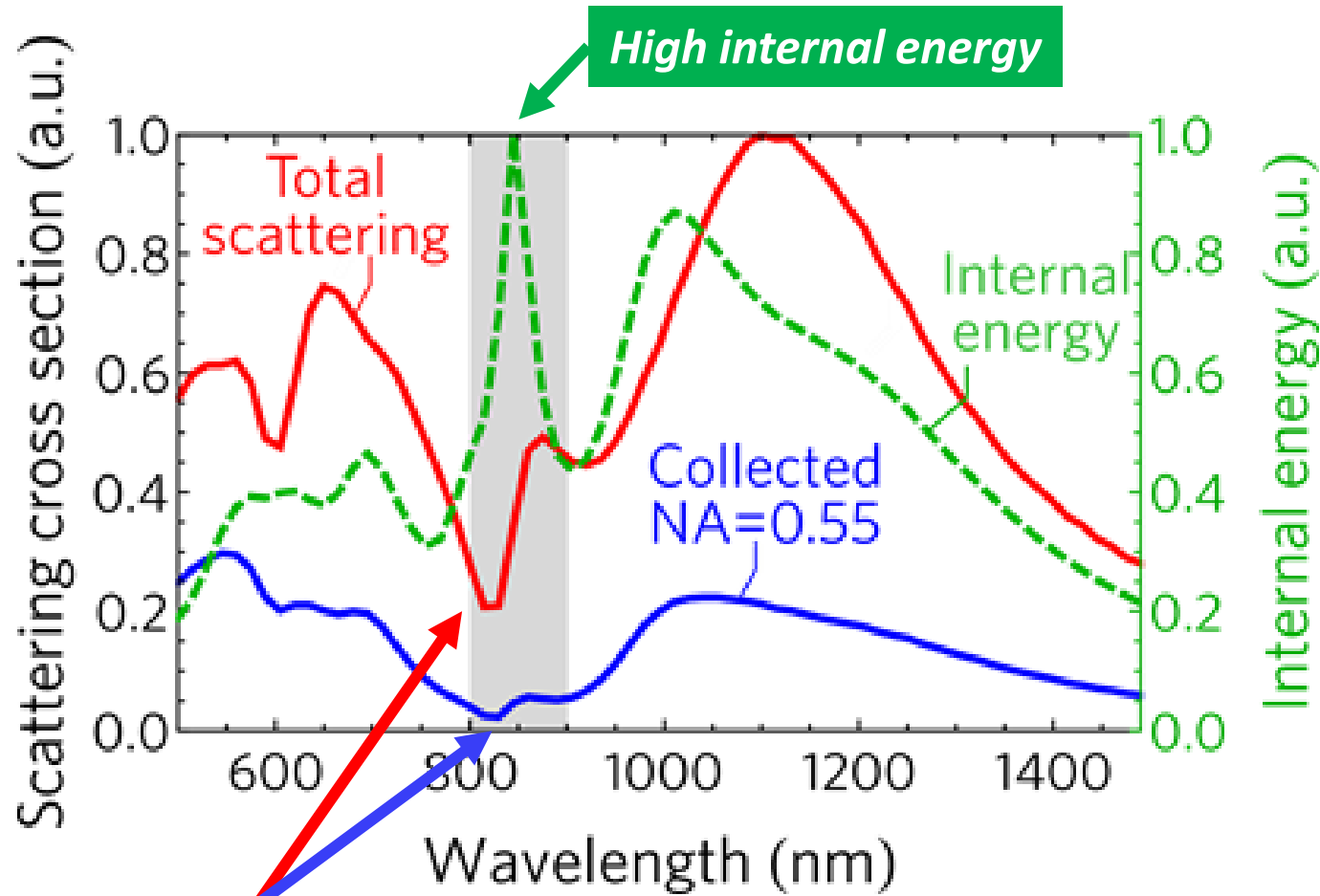
Focus Ion Beam milling



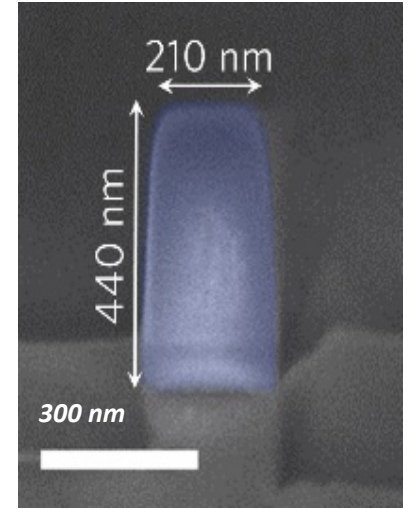
M. Timofeeva, ..., R. Grange, Nano Letters, 18,6, 2018

Mie resonances in III-V nanowires

Anapoles in Free-Standing III-V Nanodisks

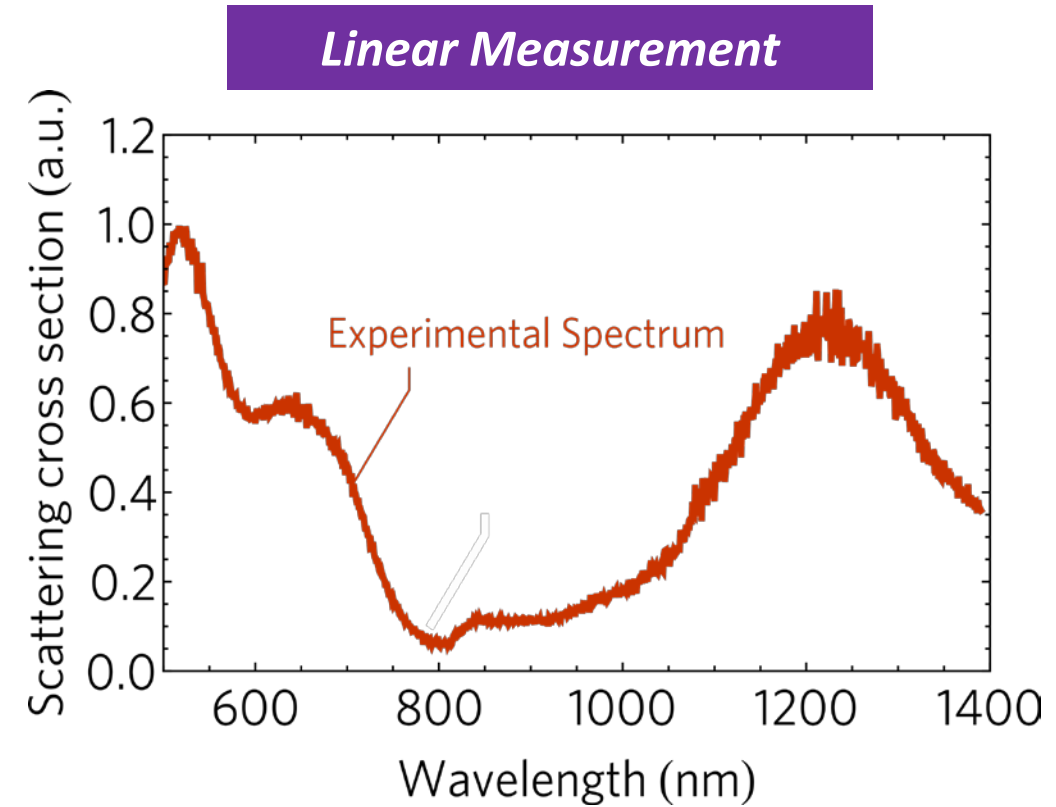
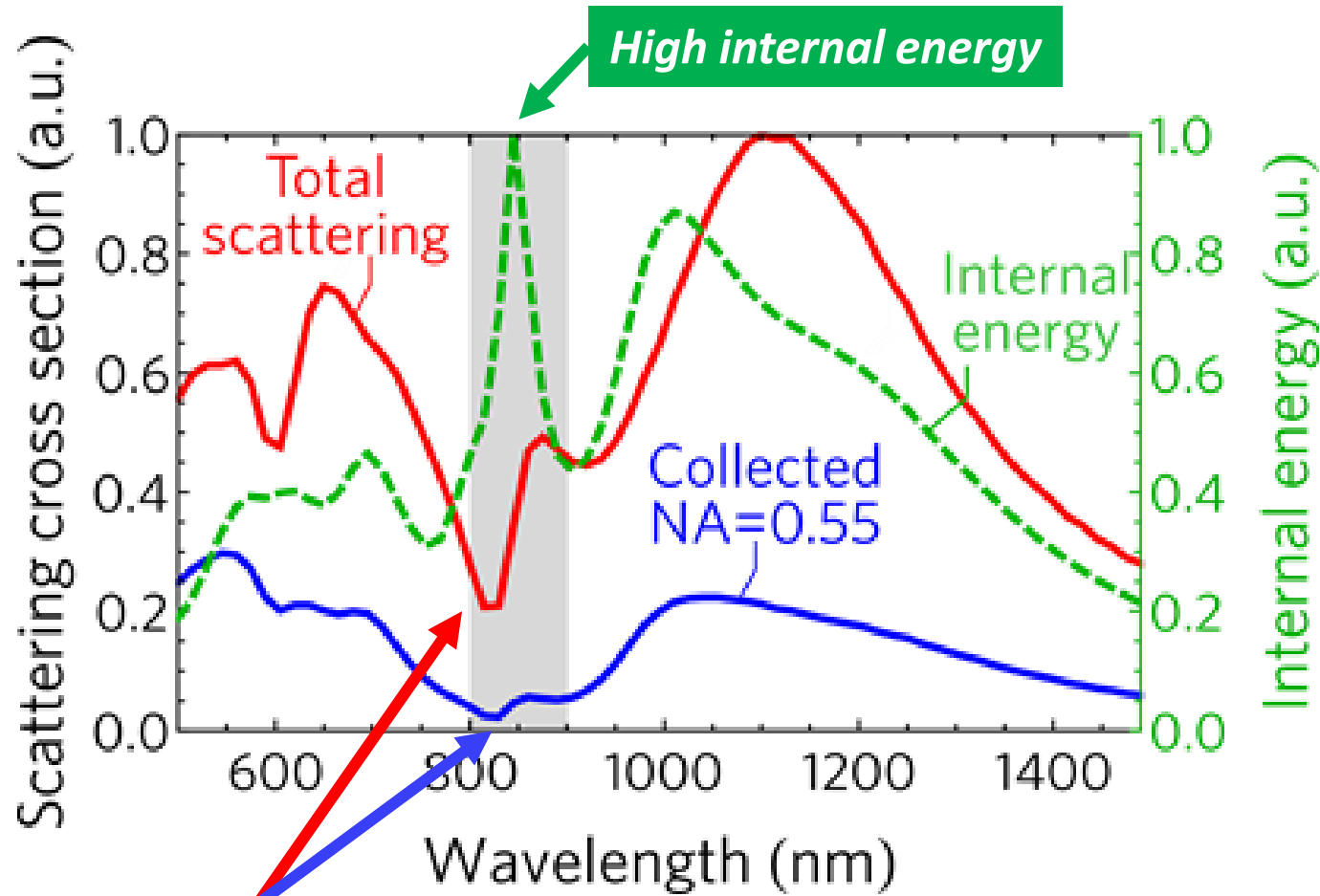


Suppression of far field radiation



M. Timofeeva, ..., R. Grange, Nano Letters, 18,6, 2018

Anapoles in Free-Standing III-V Nanodisks

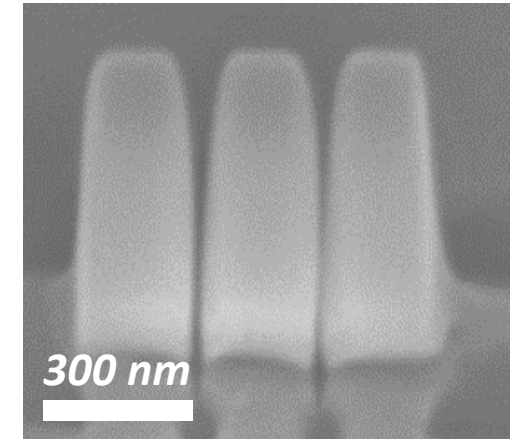
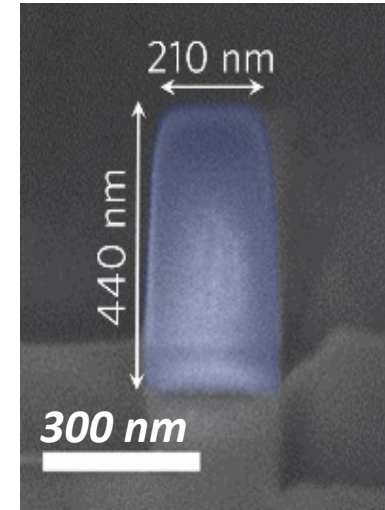
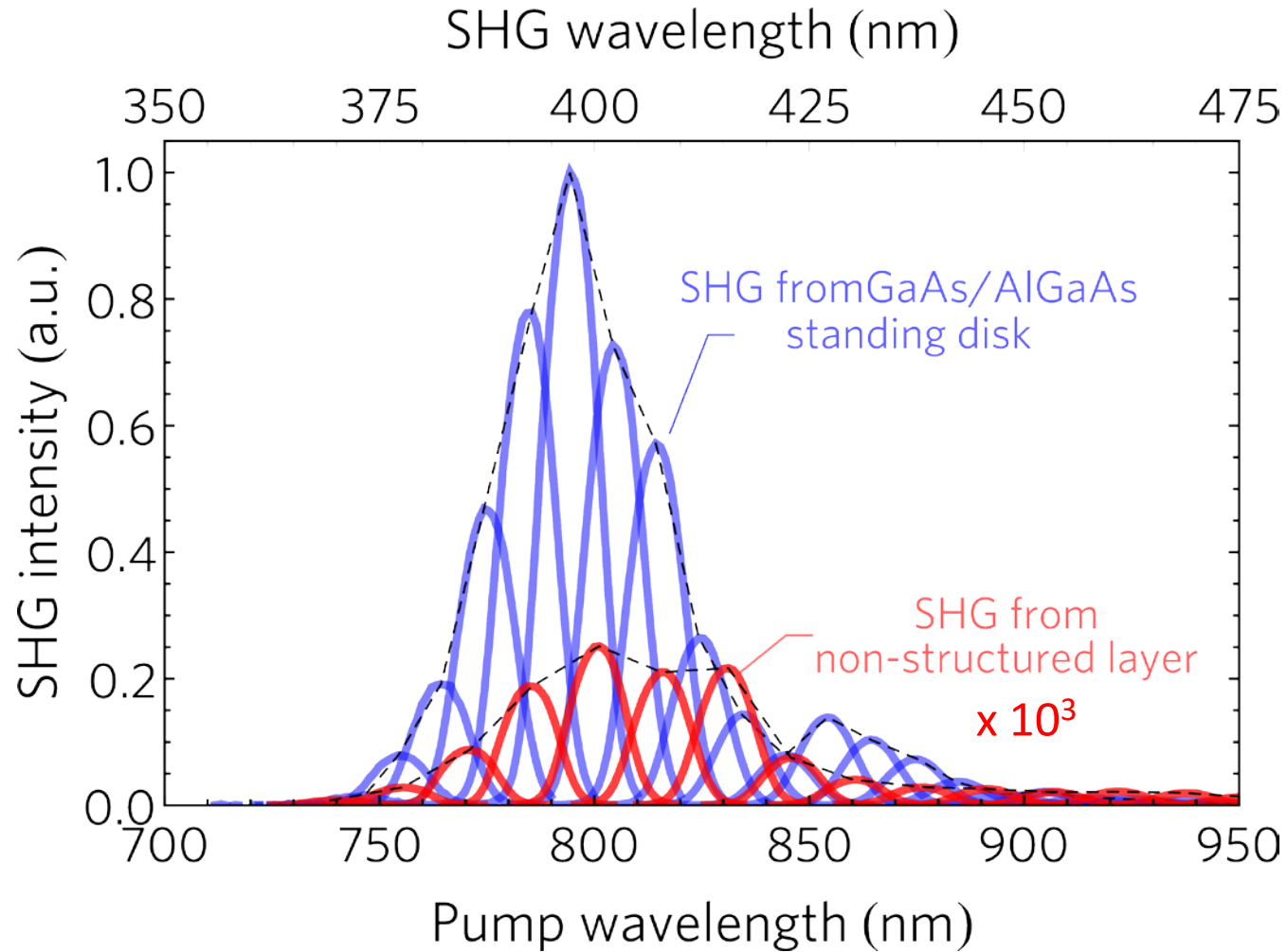


Suppression of far field radiation

M. Timofeeva, ..., R. Grange, Nano Letters, 18,6, 2018

Anapoles in Free-Standing III-V Nanodisks

Nonlinear Measurements



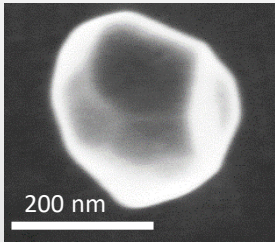
Strong nonlinear response:

>10³ SHG enhancement

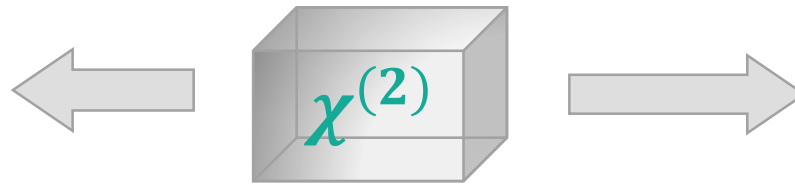
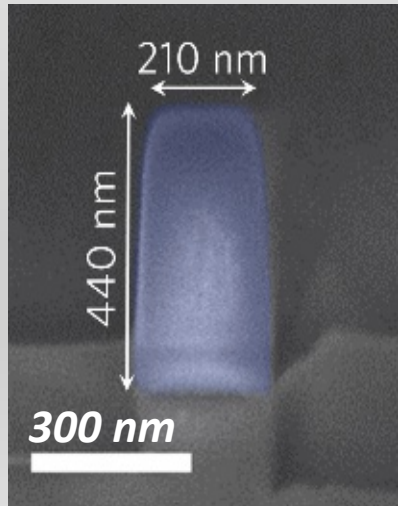
M. Timofeeva, ..., R. Grange, Nano Letters, 18,6, 2018

$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks

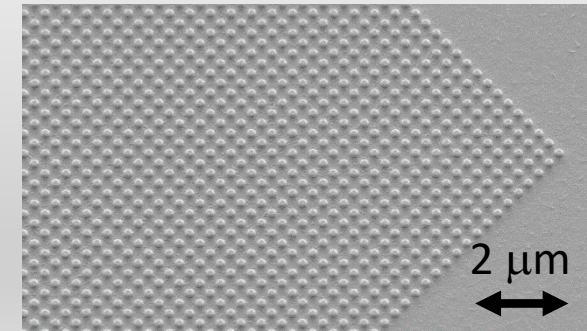


20 - 100s nm



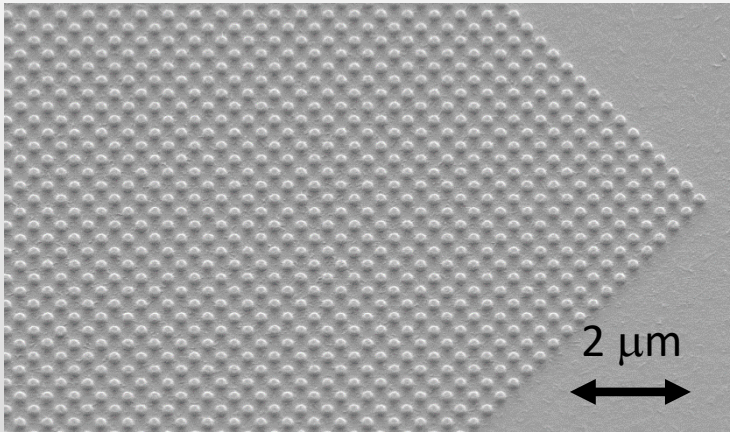
Top-down fabrication

Flat photonic structures



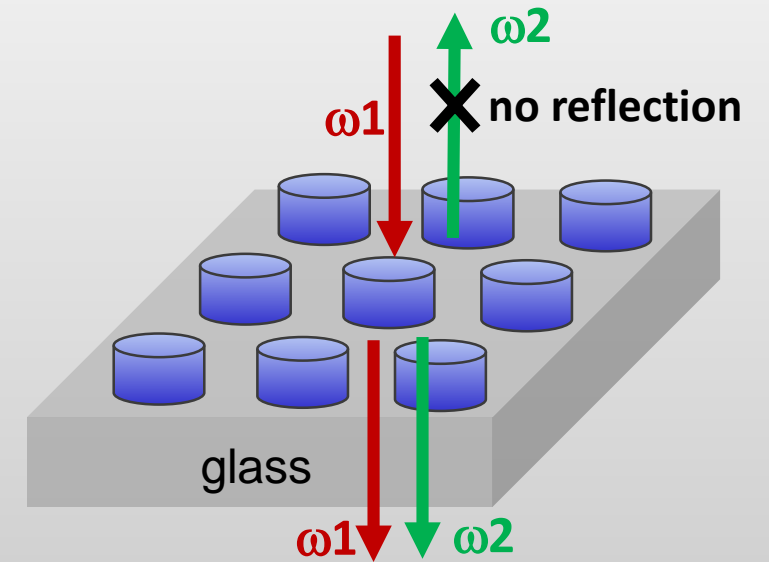
Perspective: Flat Nonlinear Photonic Devices

BaTiO₃ thin film for nonlinear metasurfaces



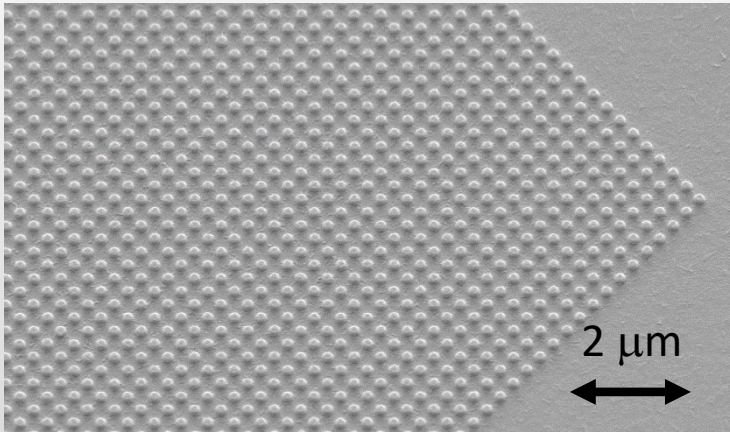
Pulse Laser Deposition of M. Trassin and M. Fiebig, Material sciences, ETHZ

Artificial nanostructured materials with exotic electromagnetic responses

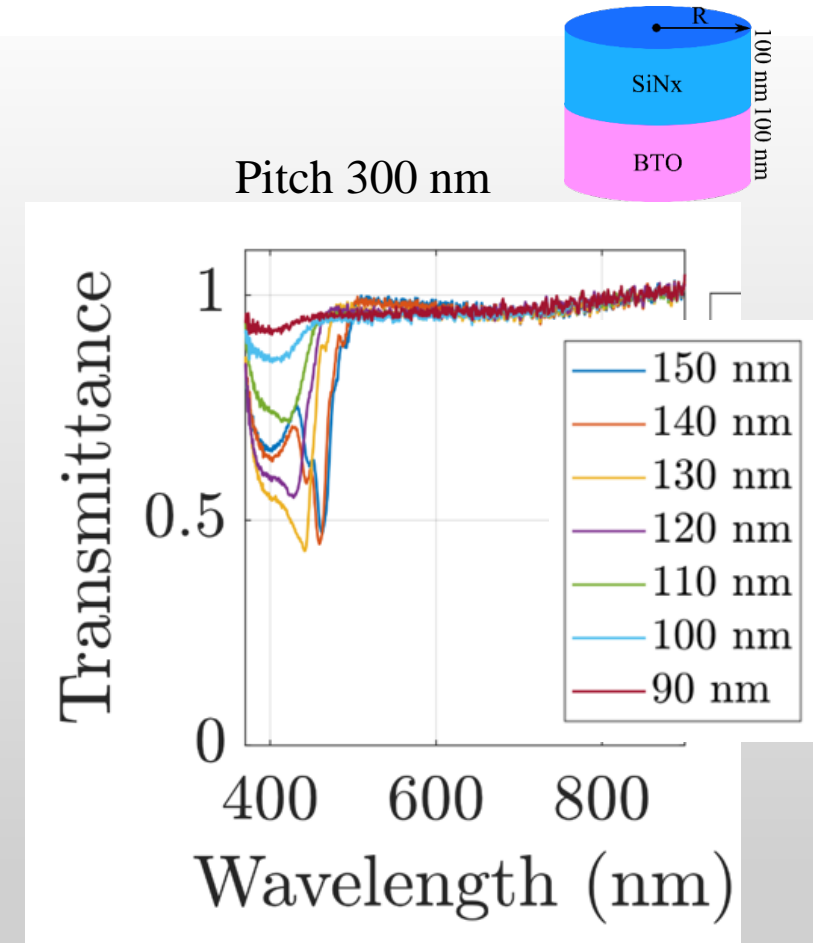
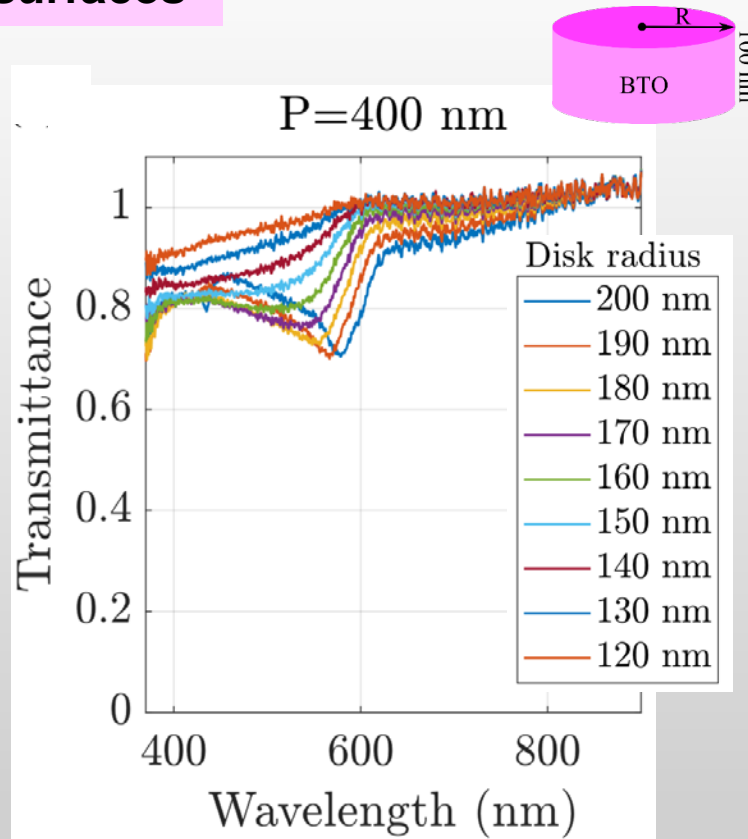


Perspective: Flat Nonlinear Photonic Devices

BaTiO₃ thin film for nonlinear metasurfaces



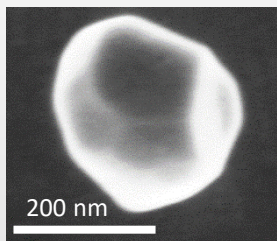
Pulse Laser Deposition of M. Trassin and M. Fiebig, Material sciences, ETHZ



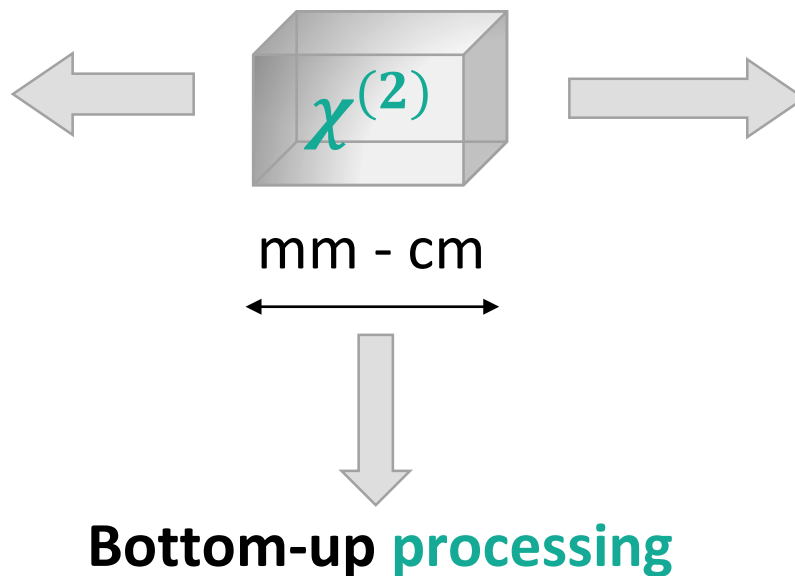
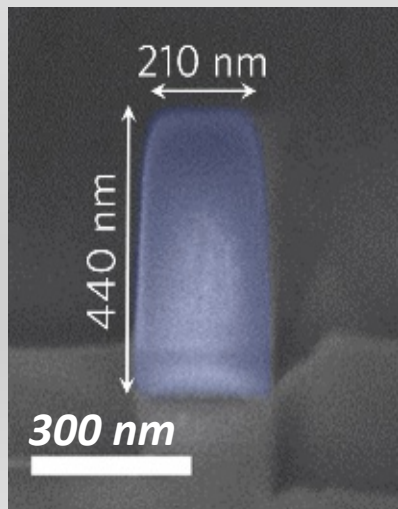
Fabrication and etching challenges of oxides thin films

$\chi^{(2)}$ materials at small scale

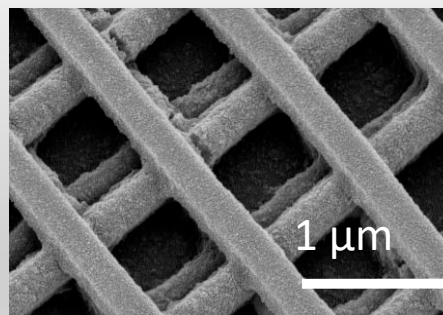
$\chi^{(2)}$ Building blocks



20 - 100s nm

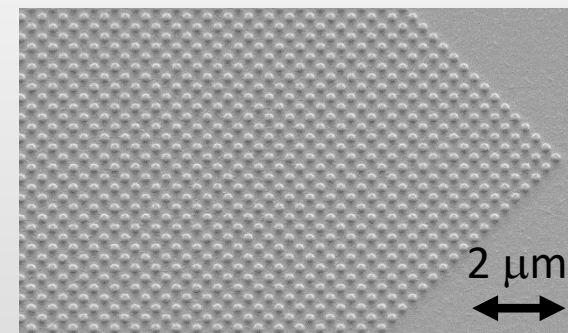


Wide surface area photonic crystal



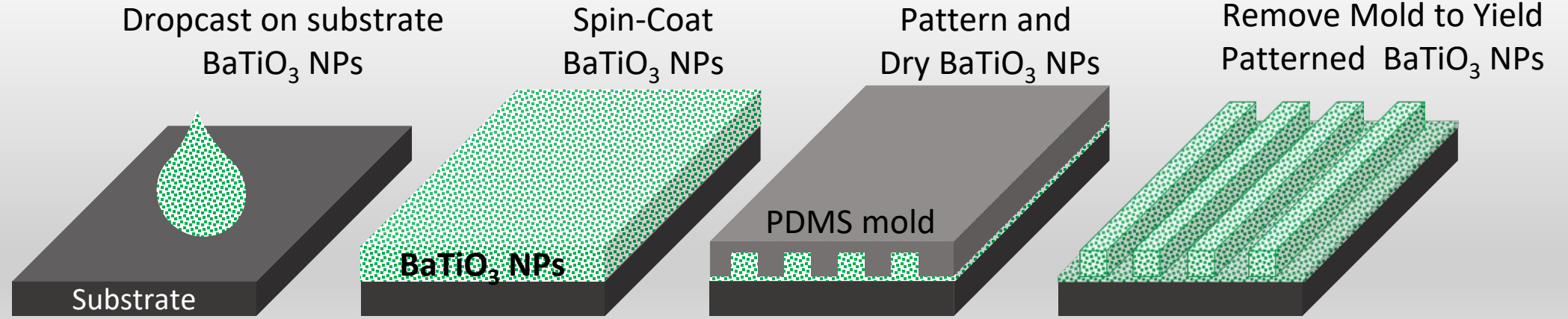
Top-down fabrication

Flat photonic structures

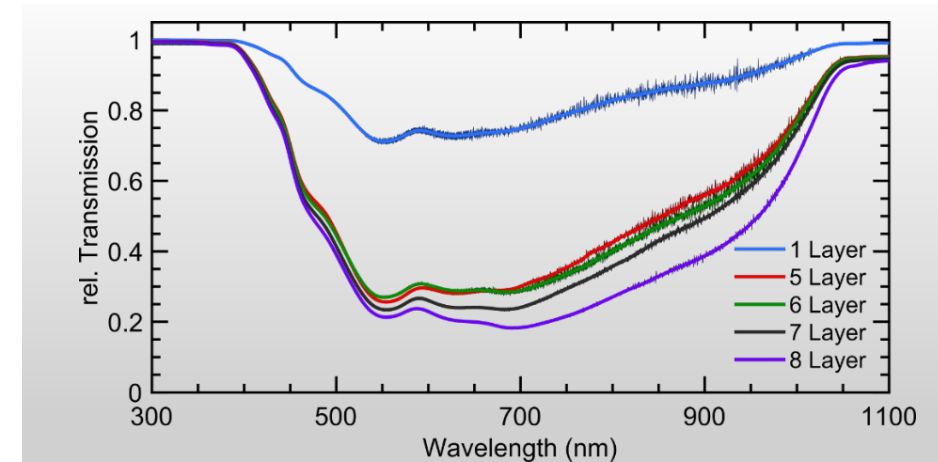
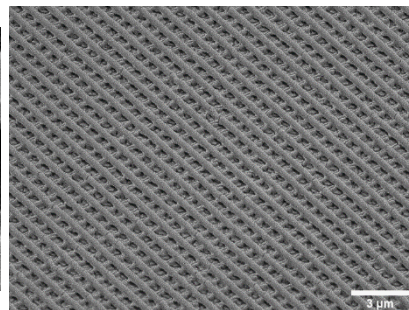
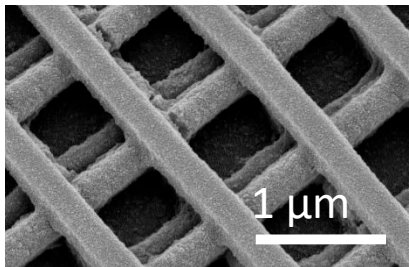


Bottom-up assembly for large surface area

Colloidal Solution Processing and nanoimprint lithography



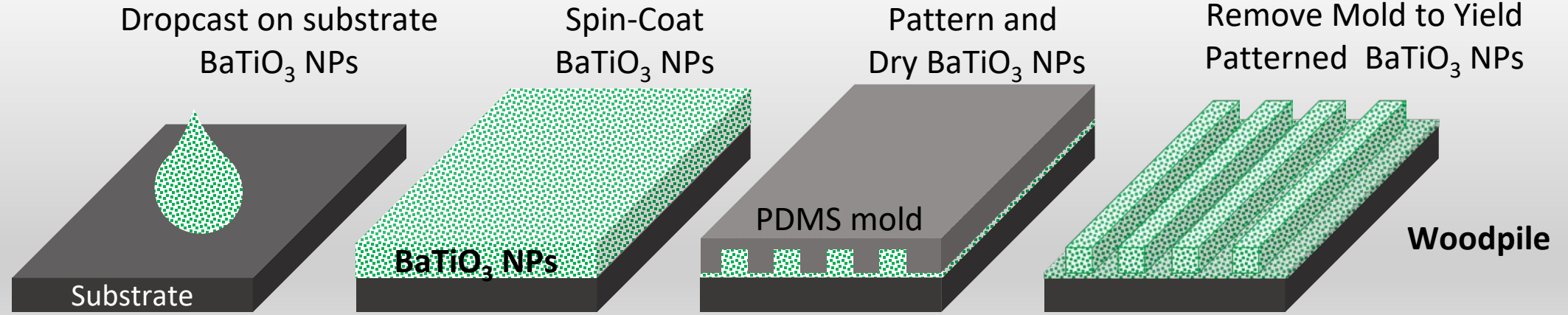
Woodpile up to 8 layers



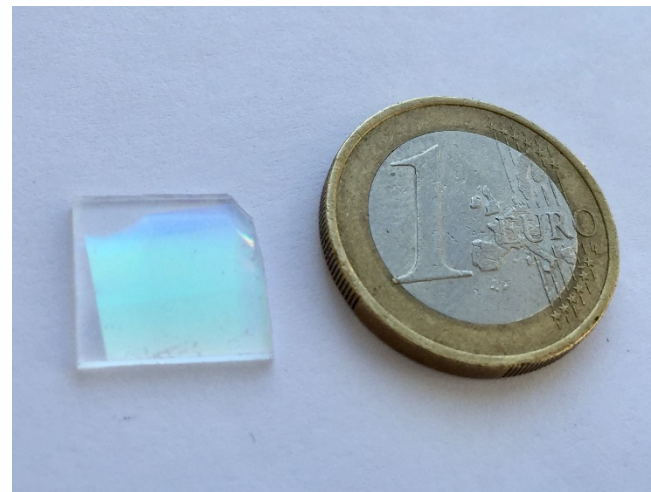
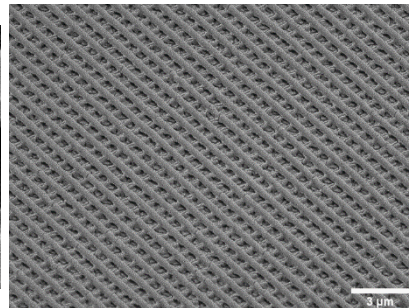
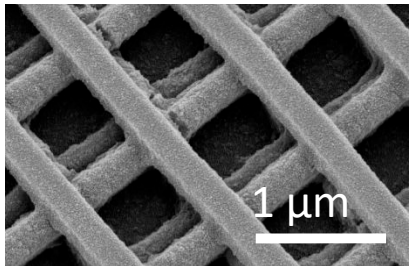
V. V. Vogler-Neuling, ... and R. Grange, Metamaterials
Marseille, France 08.28-31, Talk (2017).

Bottom-up assembly for large surface area

Colloidal Solution Processing and nanoimprint lithography



Woodpile up to 8 layers



Large surface area 2 x 2 cm

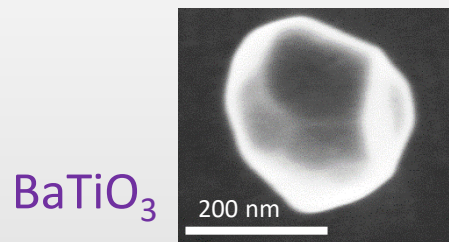
Disposable sensors

No etching

V. V. Vogler-Neuling, ... and R. Grange, *Metamaterials* Marseille, France 08.28-31, Talk (2017).

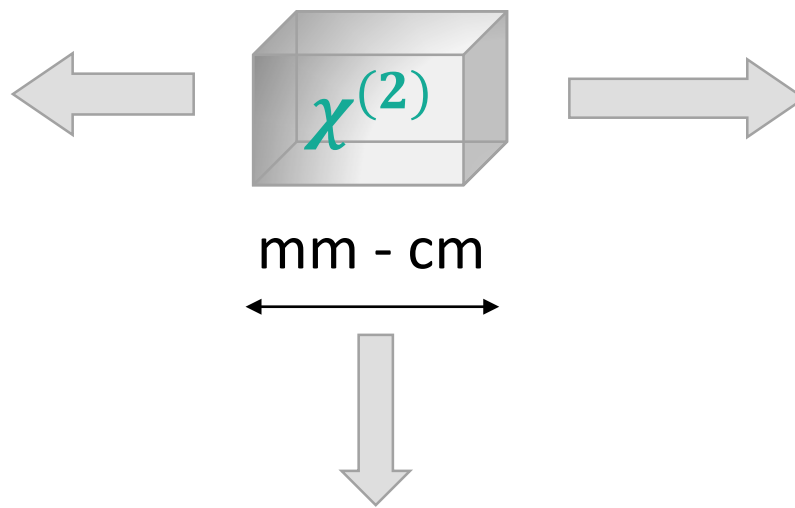
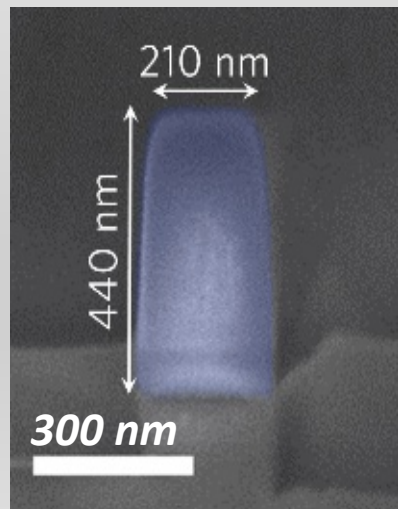
$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks



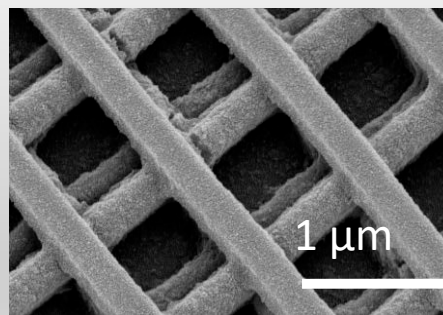
20 - 100s nm

AlGaAs anapole



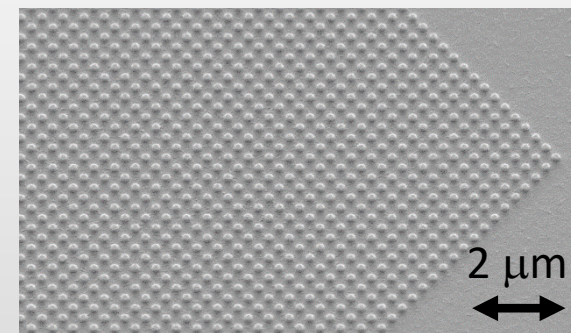
Bottom-up processing

Wide surface area photonic crystal

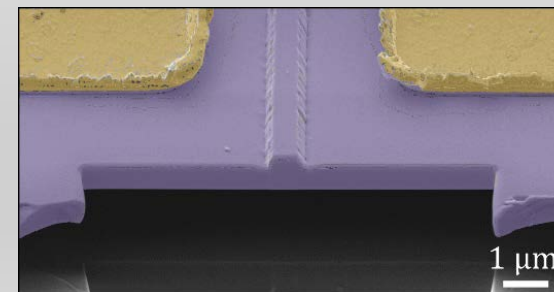


Top-down fabrication

Flat photonic structures

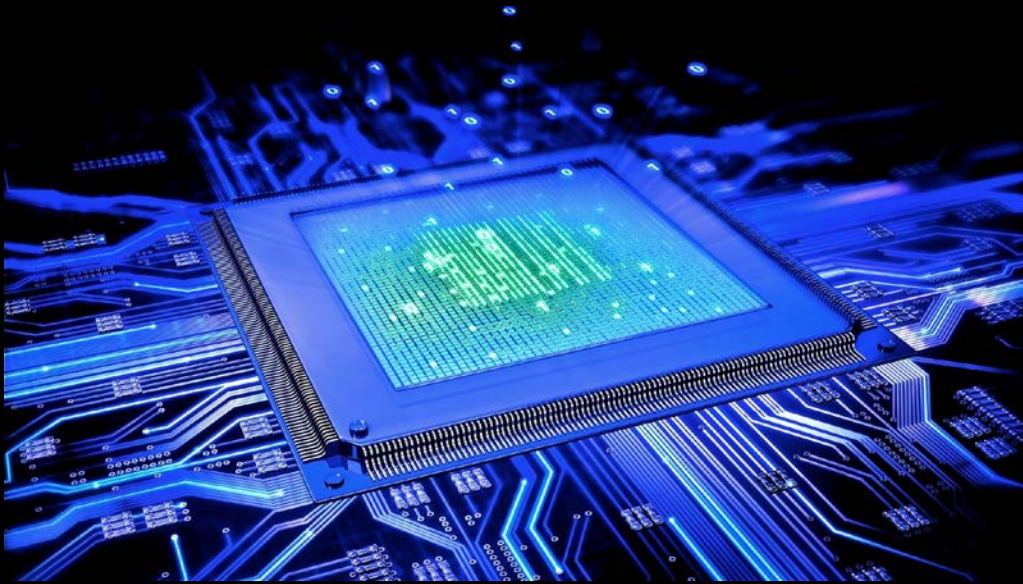


Integrated devices



Integrated Photonics

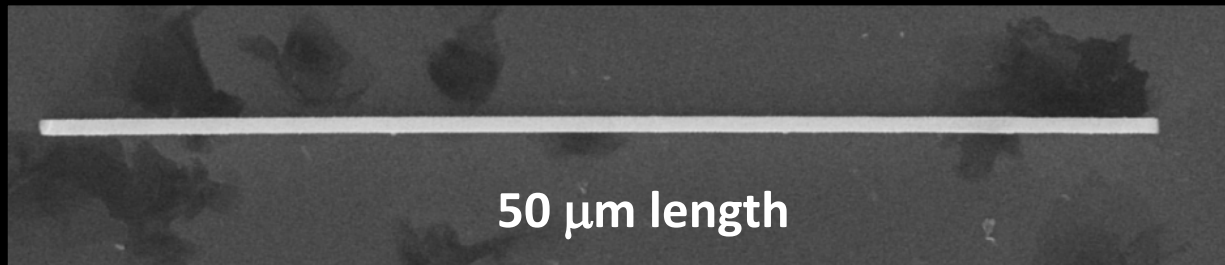
Less weight and more transmitted data



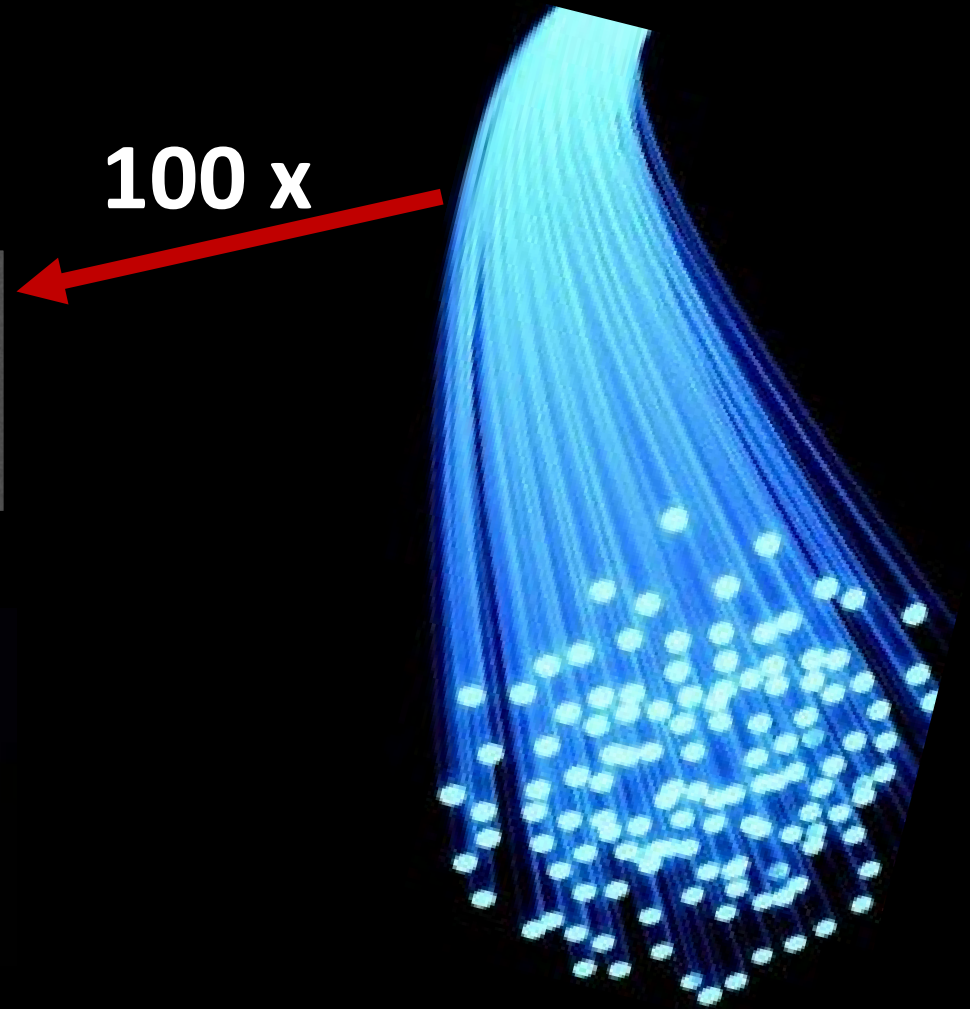
Broadly transparent, low losses, fast, compact, robust

LiNbO₃ nanowaveguides by ion beam enhanced etching

400 nm diameter



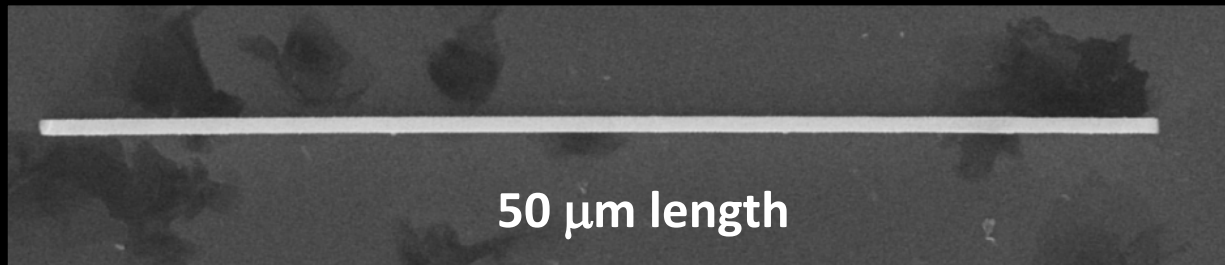
100 x



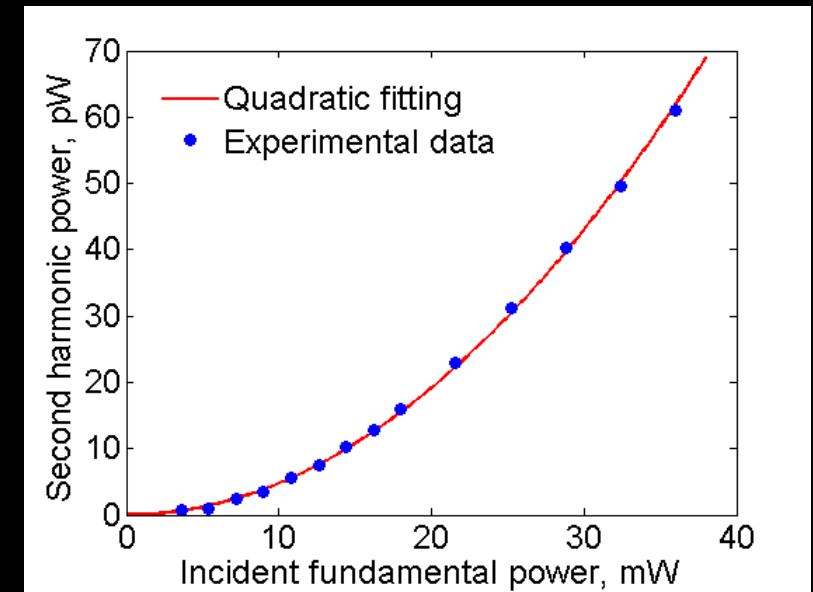
A. Sergeev, ... R. Grange, Opt. Exp. **21**, 19012, 2013.

LiNbO₃ nanowaveguides by ion beam enhanced etching

400 nm diameter

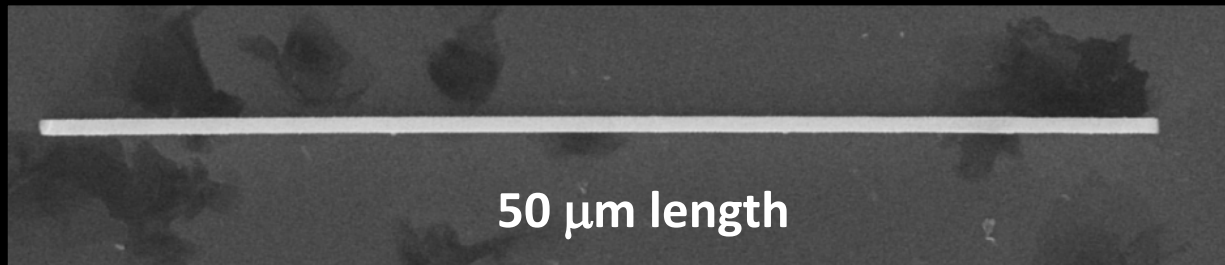


Signal dependence
on the incident laser power



LiNbO₃ nanowaveguides by ion beam enhanced etching

400 nm diameter



How to achieve phase-matching?

In waveguides

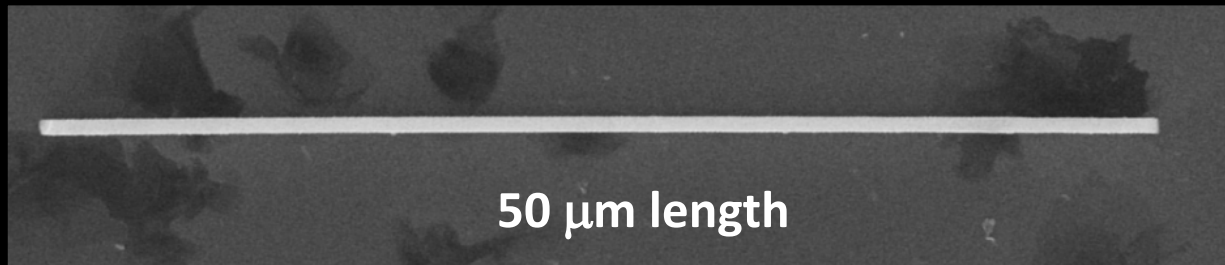
modal phase-matching



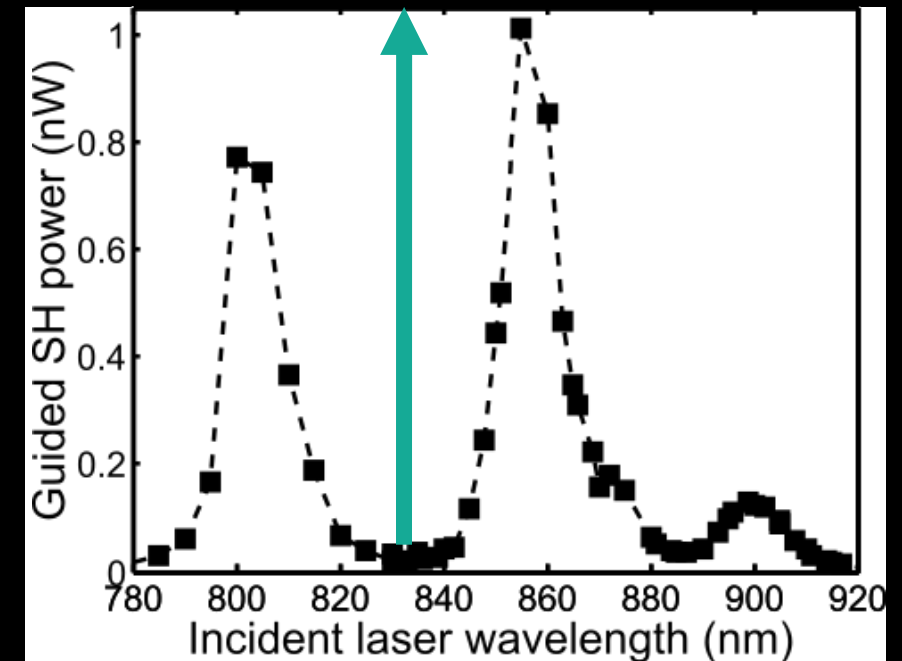
Sergeyev, A.; et al. ACS Photonics 2015, 2 (6), 687–691.

LiNbO₃ nanowaveguides by ion beam enhanced etching

400 nm diameter



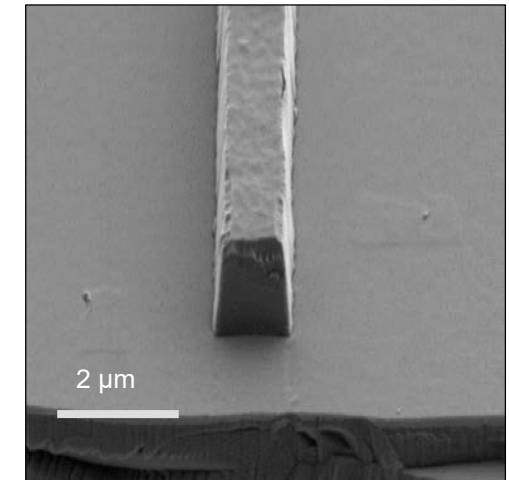
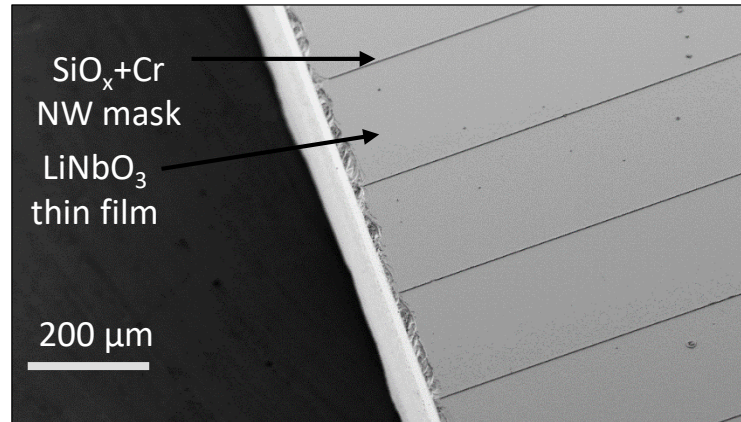
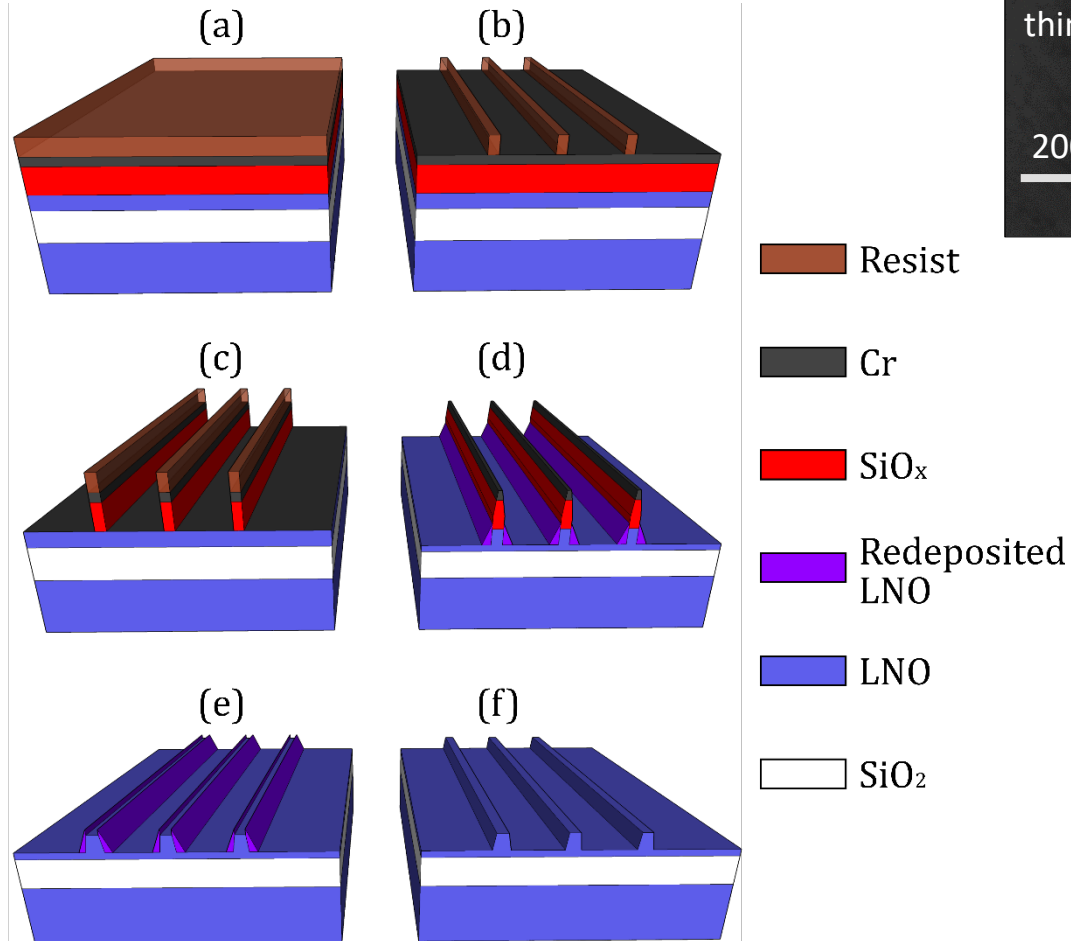
50 times increase with modal phase matching



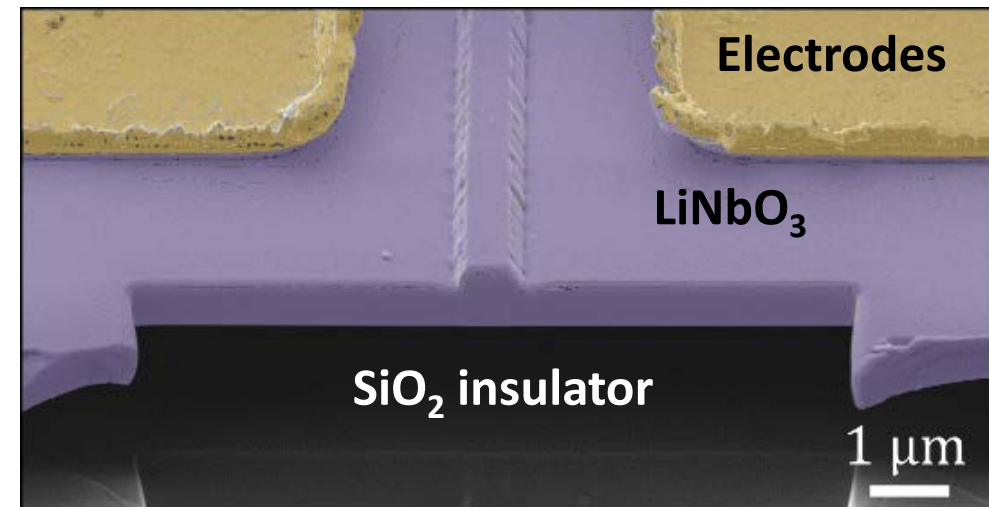
Sergeyev, A.; et al. ACS Photonics 2015, 2 (6), 687–691.

LiNbO₃ on insulator LNOI

Process flow: Plasma etching
(BRNC & First clean rooms facilities)



LNOI ridge waveguides



Escalé, M. R.; Pohl, D.; Sergeyev, A.; Grange, R. Optics Letters, 43, 1515, 2018

Extreme electro-optic tuning in lithium niobate

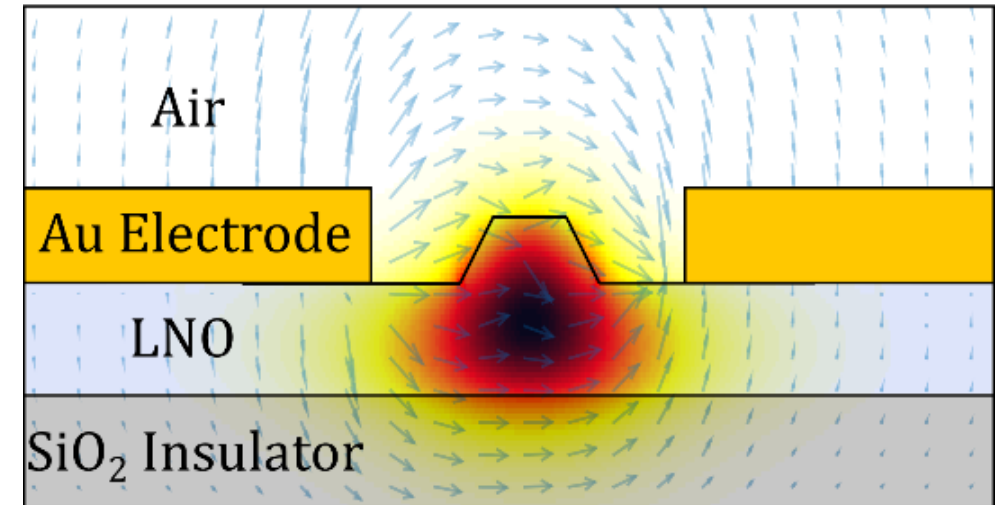
Pockel's effect

Change in the refractive index
linearly proportional to the electric field

$$\begin{pmatrix} \Delta(1/n^2)_1 \\ \Delta(1/n^2)_2 \\ \Delta(1/n^2)_3 \\ \Delta(1/n^2)_4 \\ \Delta(1/n^2)_5 \\ \Delta(1/n^2)_6 \end{pmatrix} = \begin{pmatrix} 0 & -3.4 & 8.6 \\ 0 & 3.4 & 8.6 \\ 0 & 0 & 30.8 \\ 0 & 28 & 0 \\ 28 & 0 & 0 \\ -3.4 & 0 & 0 \end{pmatrix} \cdot \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix}$$

Electro-optic
tensor of LiNbO_3

Electro-optic tuning capabilities

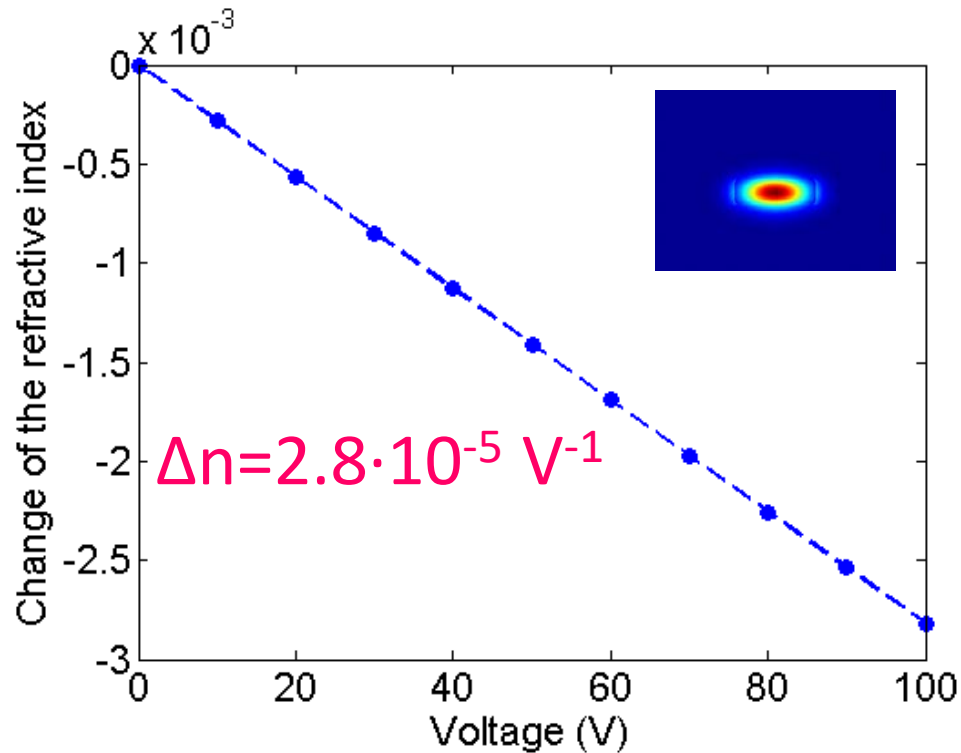


Electric field has to be
parallel to the z-crystal axis

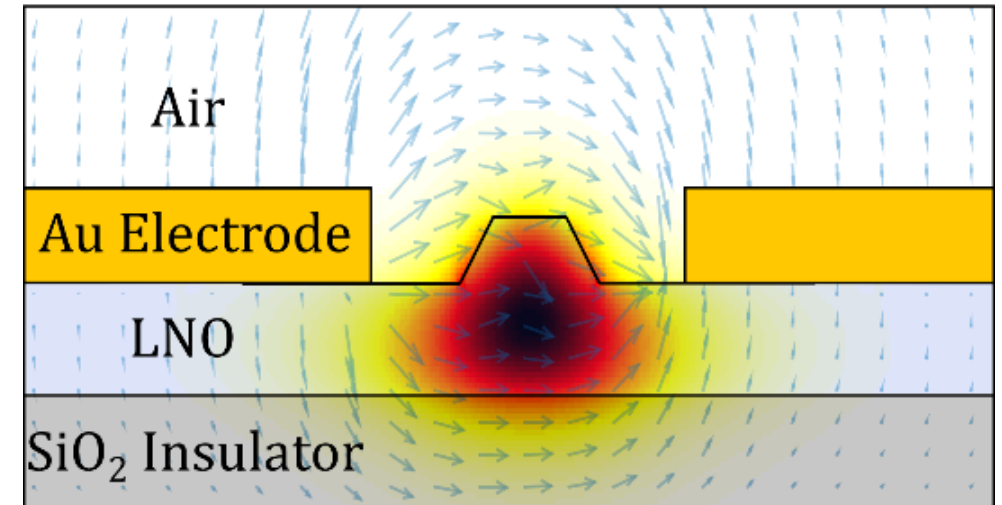
Extreme electro-optic tuning in lithium niobate

Pockel's effect

Theoretical change of the refractive index of TE00



Electro-optic tuning capabilities

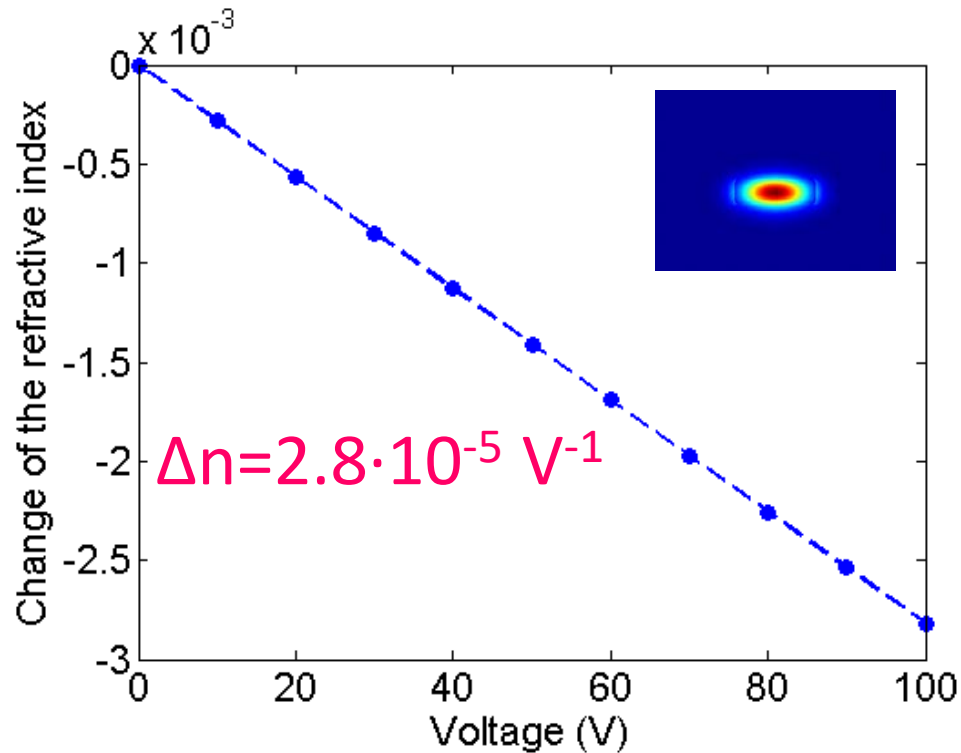


Electric field has to be parallel to the z-crystal axis

Extreme electro-optic tuning in lithium niobate

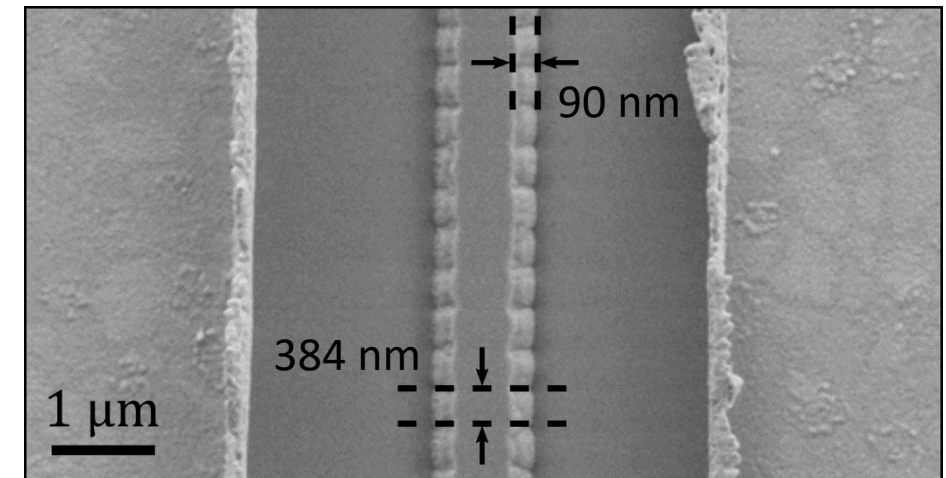
Pockel's effect

Theoretical change of the refractive index of TE00



Distributed Bragg reflector

Multilayers of alternating materials with varying n , each layer causes a partial reflection



Period : 384 nm

Corrugation depth : 90 nm

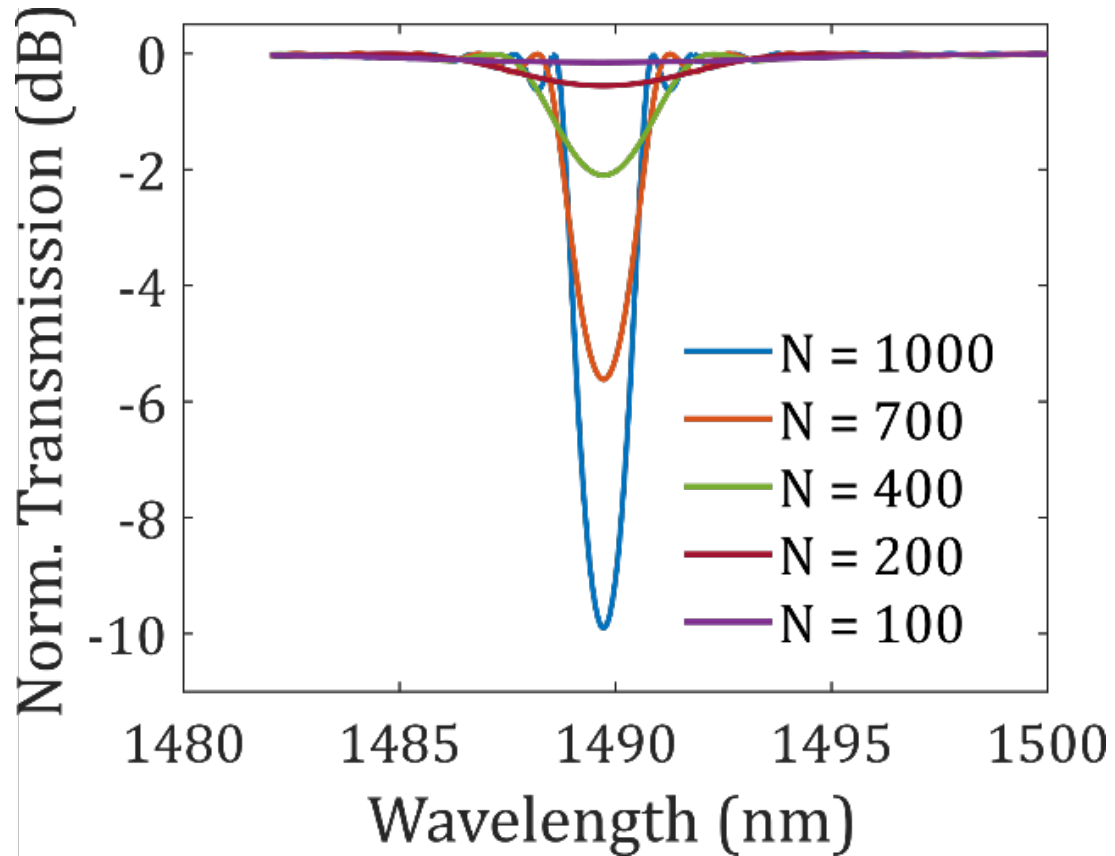
Propagation losses:

$$\alpha_{\text{TE}} = 8.07 \text{ dB/cm}$$

$$\alpha_{\text{TM}} = 3 \text{ dB/cm}$$

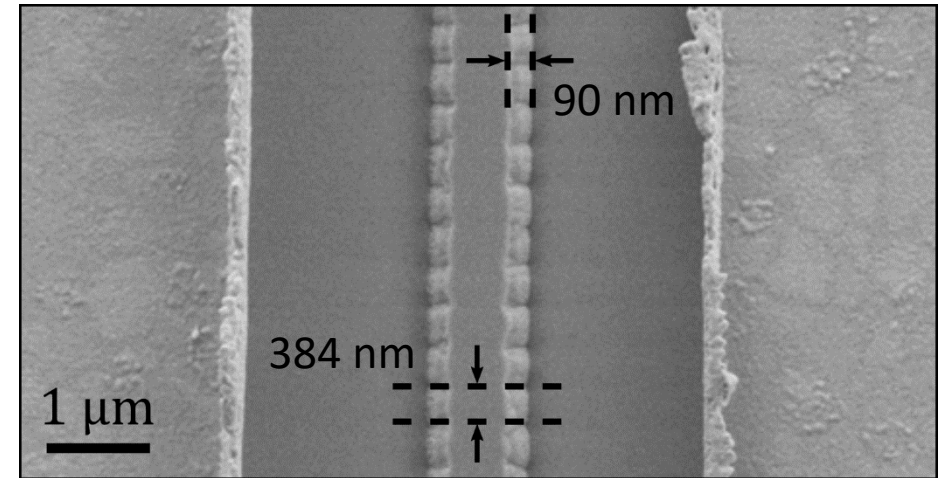
Extreme electro-optic tuning in lithium niobate

Stop Band: Transmission Spectra



Distributed Bragg reflector

Multilayers of alternating materials with varying n , each layer causes a partial reflection



Period : 384 nm

Corrugation depth : 90 nm

Propagation losses:

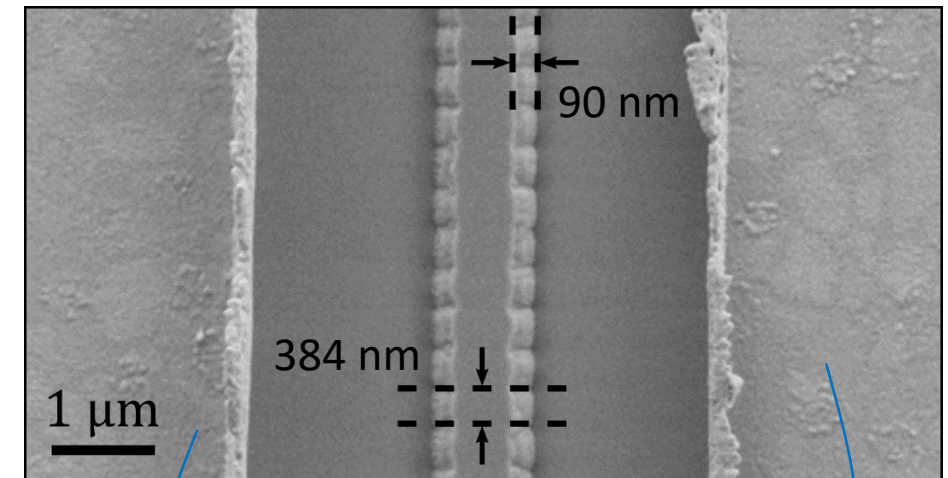
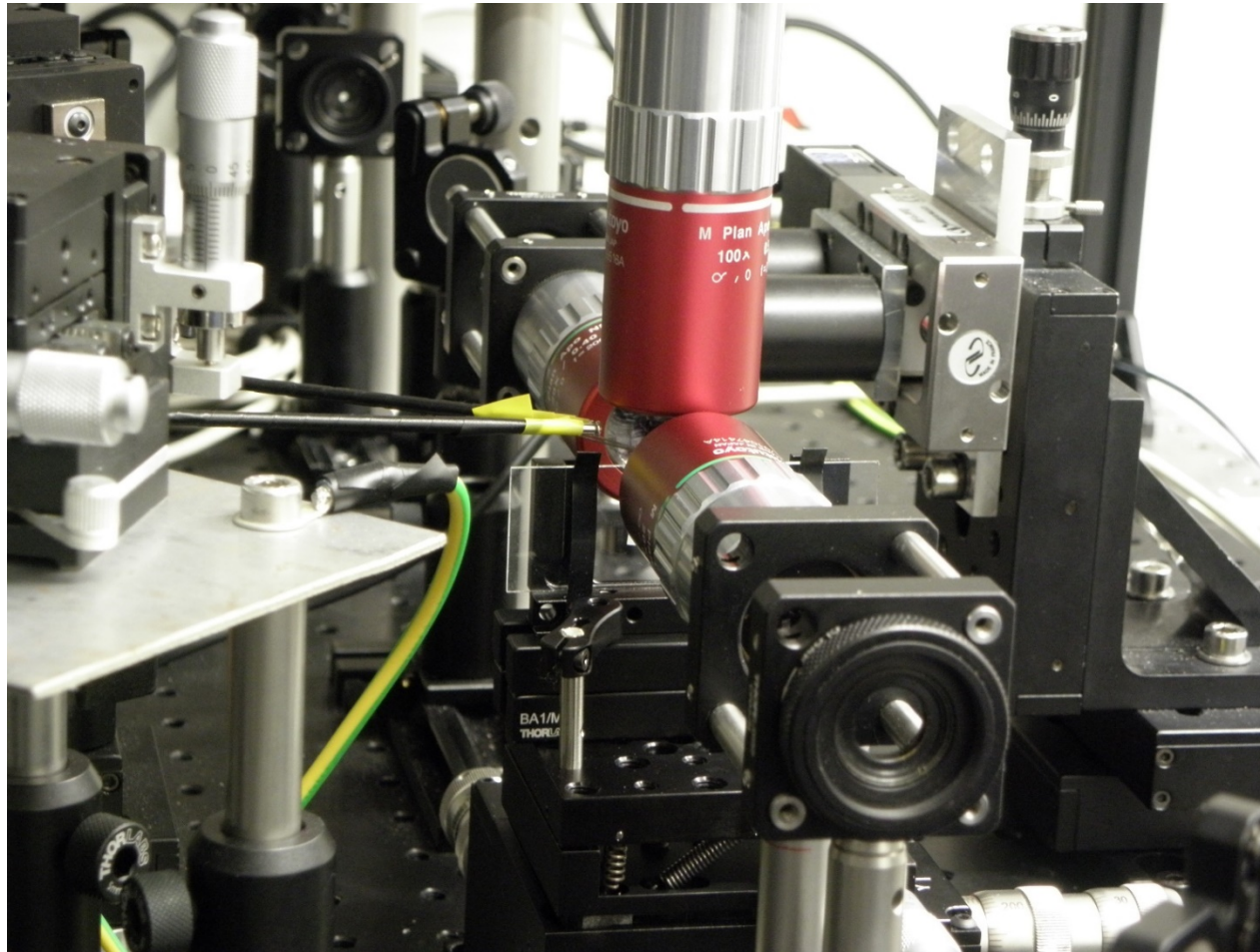
$$\alpha_{TE} = 8.07 \text{ dB/cm}$$

$$\alpha_{TM} = 3 \text{ dB/cm}$$

Extreme electro-optic tuning in lithium niobate

Distributed Bragg reflector

Multilayers of alternating materials with varying n , each layer causes a partial reflection

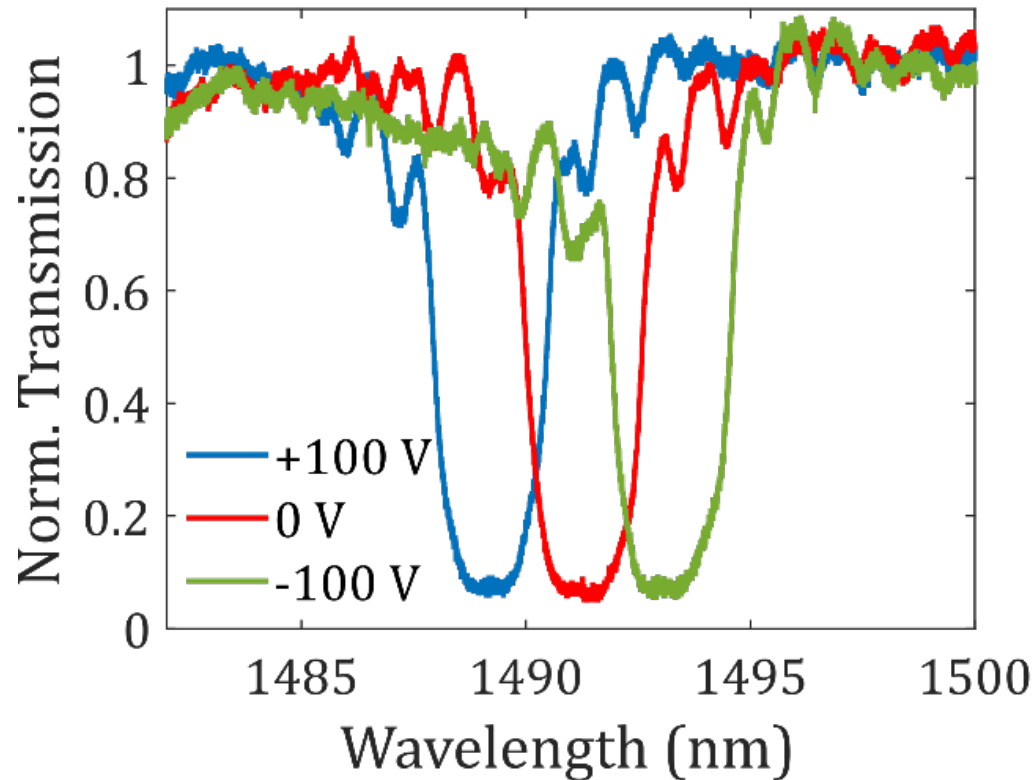


V

Escalé, M. R.; Pohl, D.; Sergeyev, A.; Grange, R. Optics Letters, 43, 1515, 2018

Extreme electro-optic tuning in lithium niobate

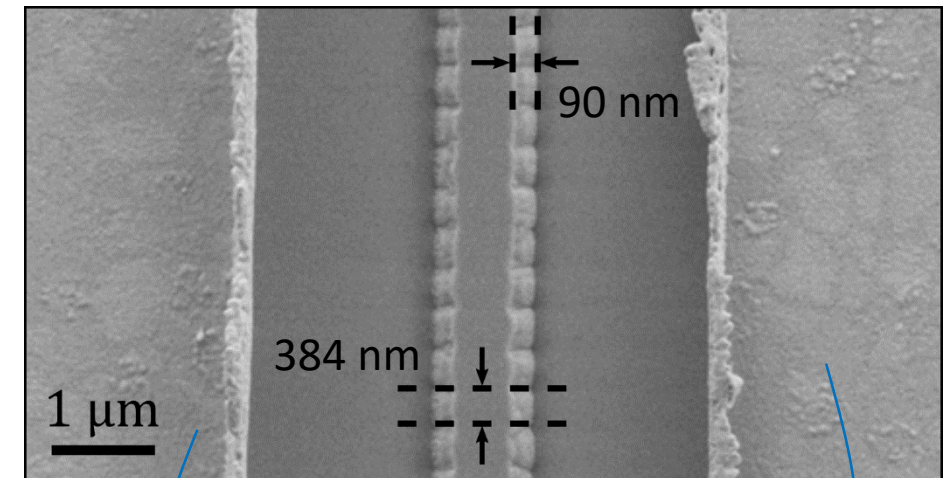
Wide tuning of the stop band



- R_{peak} 93.5 %
- 14 dB in transmission suppression

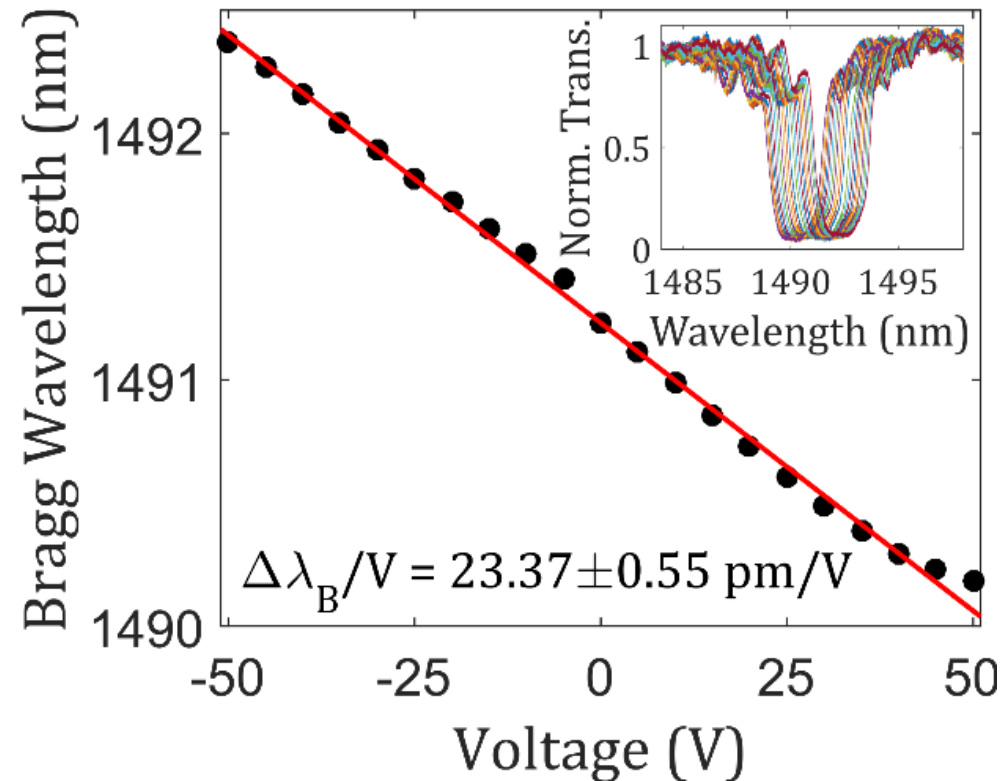
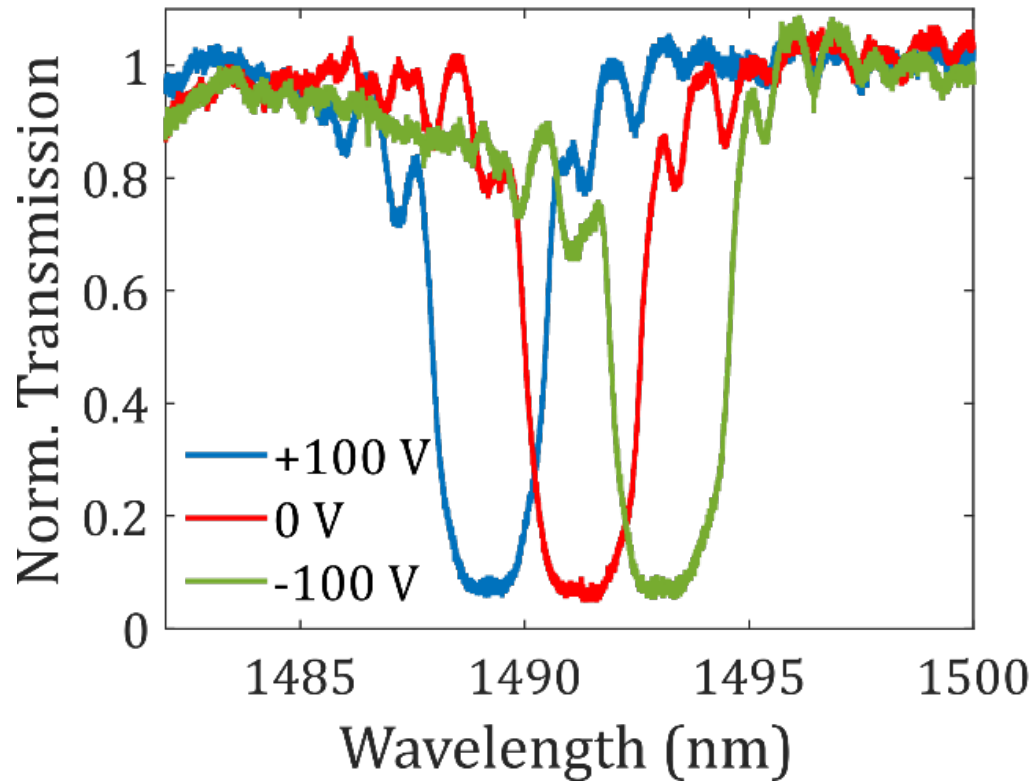
Distributed Bragg reflector

Multilayers of alternating materials with varying n , each layer causes a partial reflection



Extreme electro-optic tuning in lithium niobate

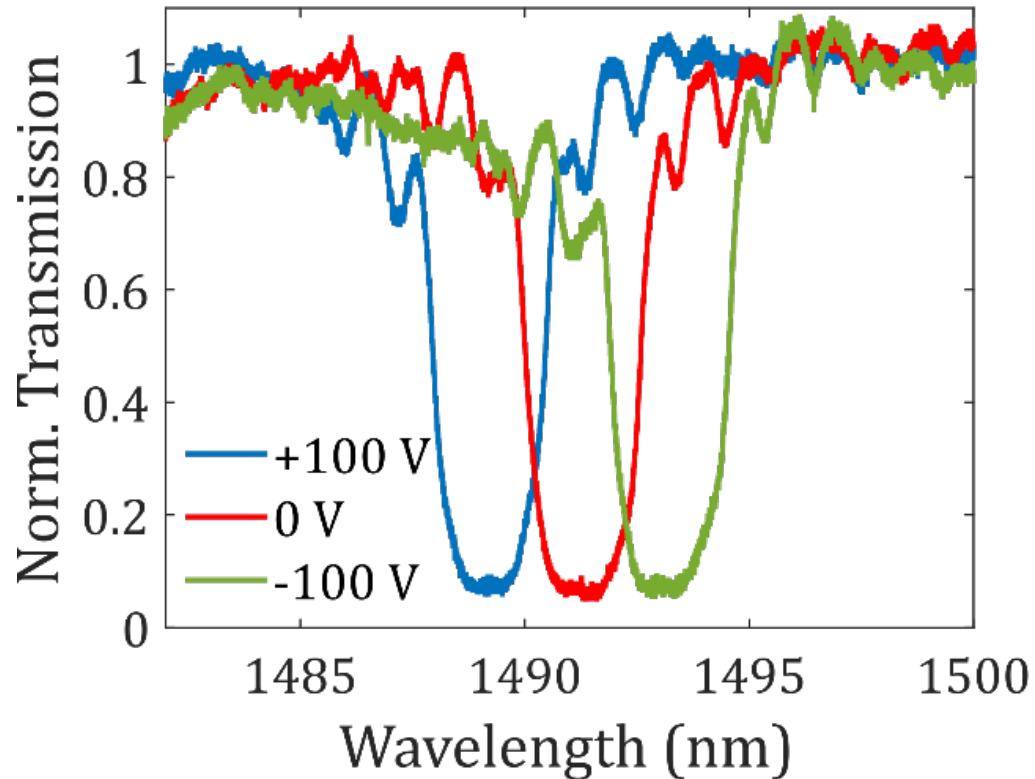
Wide tuning of the stop band



33 x more tuning than state-of-the-art devices !

Extreme electro-optic tuning in lithium niobate

Wide tuning of the stop band



Insertion losses : 12.9 dB

Extinction ratio 14dB

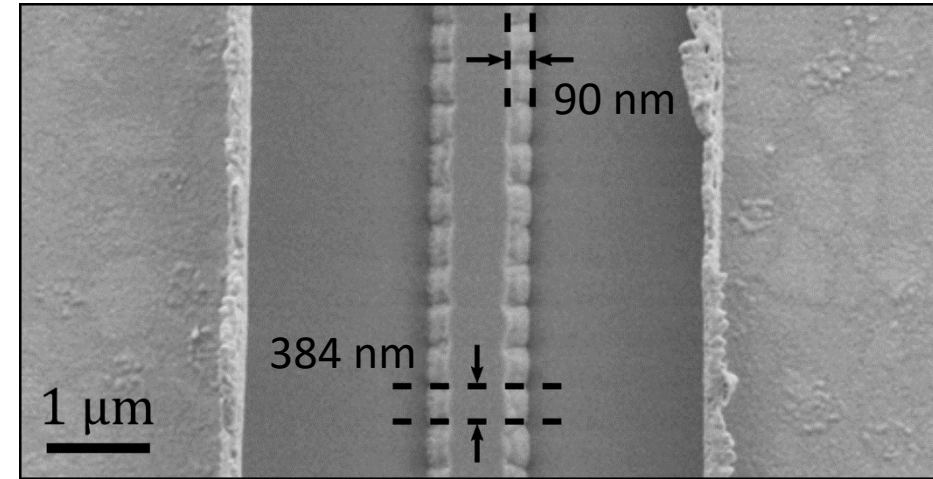
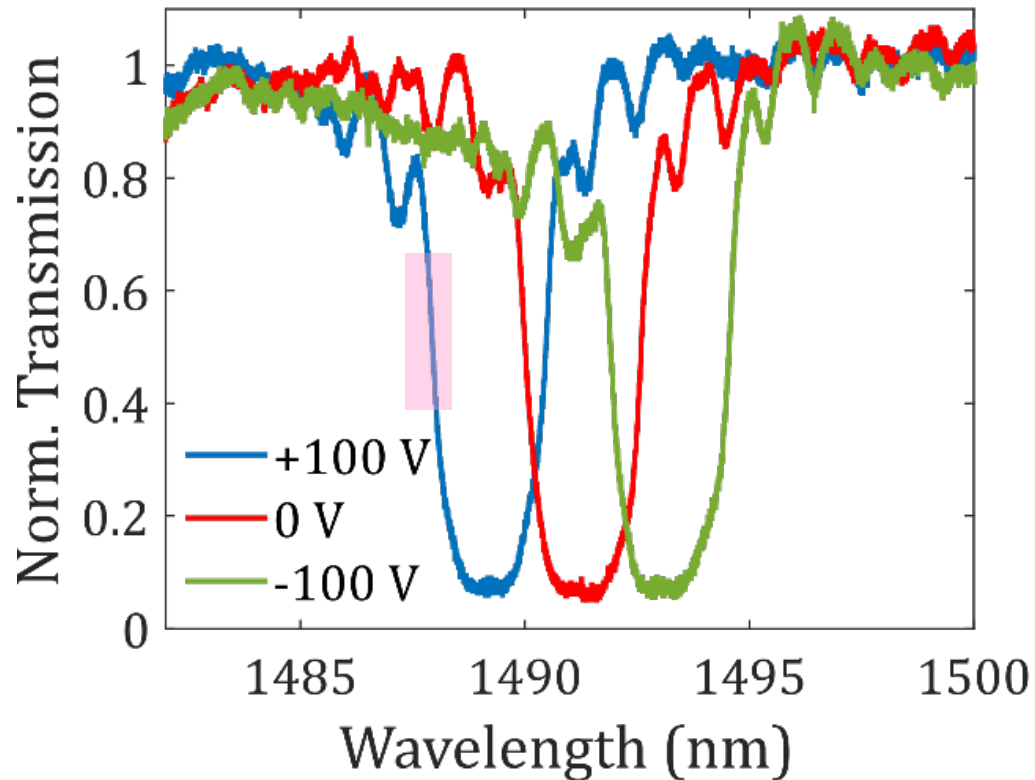
Footprint: 10×500 μm^2

Bandwidth 2nm

33 x more tuning than state-of-the-art devices !

Extreme electro-optic tuning in lithium niobate

Wide tuning of the stop band

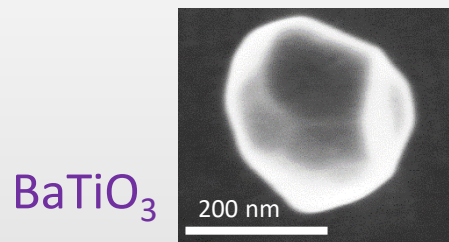


electrode gap is 4.5 μm

33 x more tuning than state-of-the-art devices !

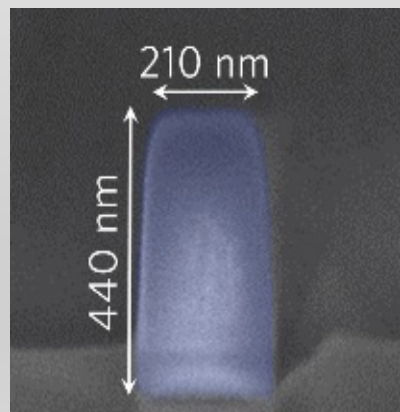
$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks

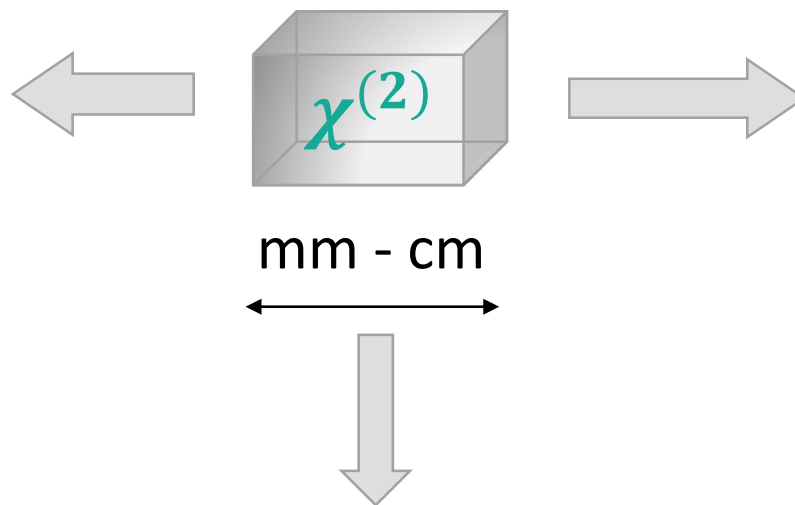


Timpu, et al. ACS Photonics, 4, 2017.

AlGaAs anapole

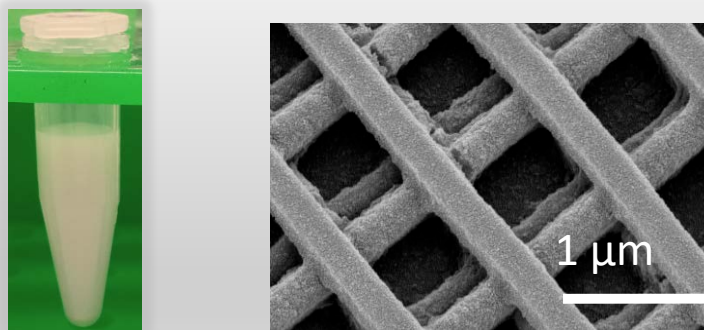


M. Timofeeva, et al. , Nano Letters, accepted May 17, 2018



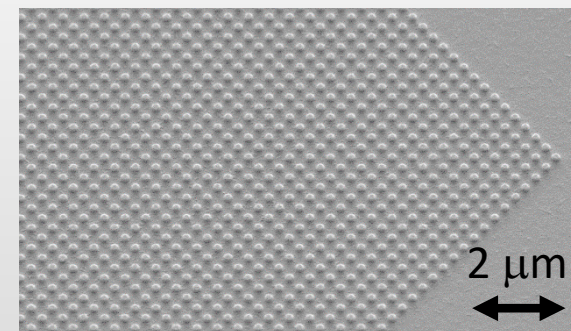
Bottom-up processing

Wide surface area photonic crystal

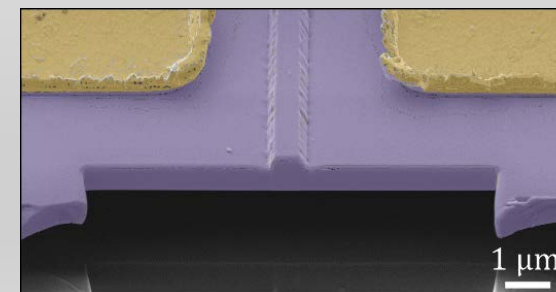


Top-down fabrication

Flat photonic structures



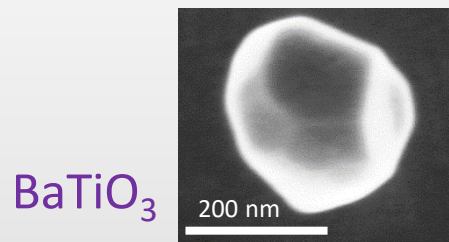
Integrated devices



Escalé, M et al. Opt. Lett., 43, 1515, 2018

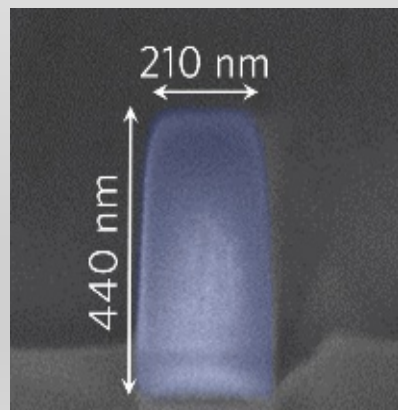
$\chi^{(2)}$ materials at small scale

$\chi^{(2)}$ Building blocks



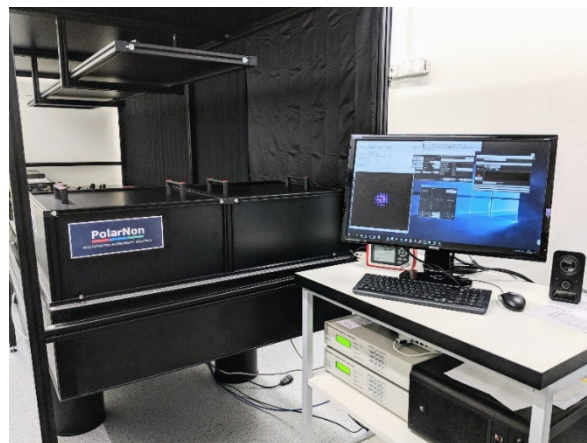
Timpu, et al. ACS Photonics, 4, 2017.
Renaut et al., Nano Lett., 2019

AlGaAs anapole



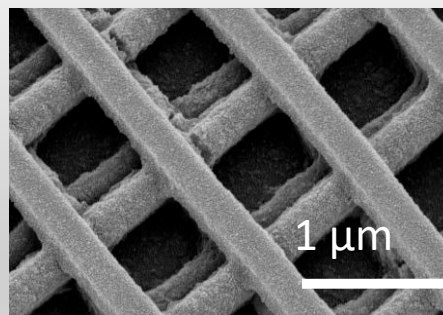
M. Timofeeva, et al. , Nano Letters, accepted May 17, 2018

Super-Polar microscope



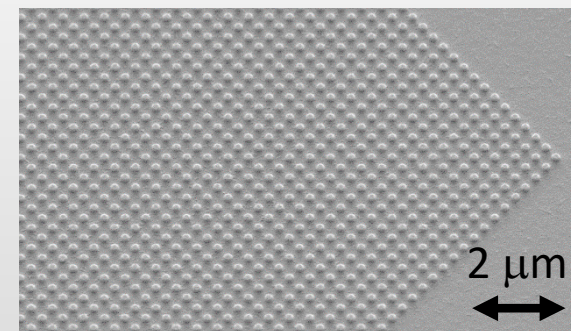
Bottom-up processing

Wide surface area photonic crystal

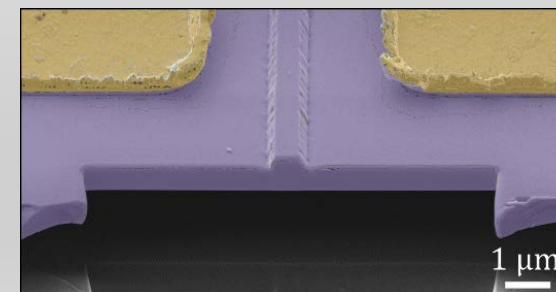


Top-down fabrication

Flat photonic structures



Integrated devices



Escalé, M et al. Opt. Lett., 43, 1515, 2018