

Metasurface Holograms

Presented by:



OSA Holography and Diffractive Optics Technical Group Webinar Event

25 January 2018, 19:00 EST



Holography and
Diffractive Optics
Technical Group

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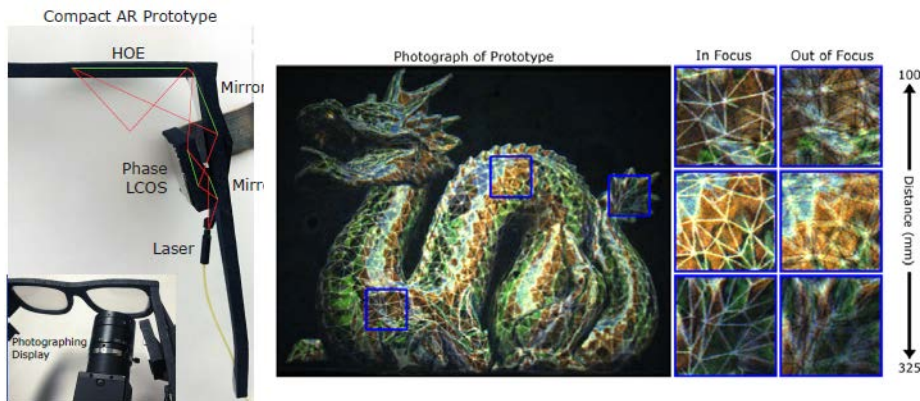
Holography and Diffractive Optics (FH)

- Our TG focuses on the design and implementation of holographic and diffractive-optic devices and systems for scientific, commercial, and other applications.
- The scope of our TG covers holography technology, holographic materials, digital holographic imaging, holographic microscope, computer-generated holograms, diffractive optical elements, holographic nano-fabrication methods, diffractive-optic micro-manipulation, spatial light modulators for phase modulation, 3D display using spatial light modulators, holographic measurements of 3D structures, phase unwrapping, and other related topics.

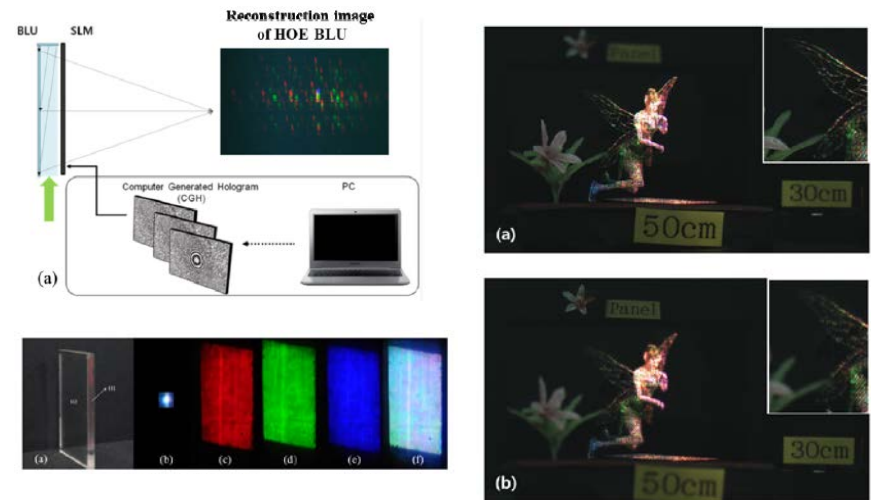
Holography and Diffractive Optics (FH)

- Research progress on holographic material / CGH / SLM is one of the key issues of holographic displays (for AR / VR), closely related to display industry

Holographic 3D Near-to-Eye Display for AR, Microsoft 2017



Holographic 3D Display, Samsung 2017



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NonImaging Optical Design (FN)

Optical Fabrication and Testing (FM)

Polarization (FP)

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Thin Films (FT)

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Photonics and Opto-Electronics

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Early Career Professionals

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Professional Development

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International Day of Light

Holography and Diffractive Optics (FH)

This group focuses on the design and implementation of holographic and diffractive-optic devices and systems for scientific, commercial, and other applications. The development of new approaches as well as materials needed to construct both analog and digital holograms are studied. The scope of this technical group covers holography technology, holographic materials, digital holographic imaging, holographic microscope, computer-generated holograms, diffractive optical elements, holographic nano-fabrication methods, diffractive-optic micro-manipulation, spatial light modulators for phase modulation, 3D display using spatial light modulators, holographic measurements of 3D structures, phase unwrapping, and other related topics.

GROUP LEADERSHIP		UPCOMING MEETINGS	RECENTLY PUBLISHED
Name	Affiliation	Title	
Yunlong Sheng	Universite Laval	Chair	
Pascal Picart	LAUM CNRS Université du Maine	Vice Chair	

Announcements

If you are a member of the Holography and Diffractive Optics Technical Group and have ideas for activities that may be of interest to your fellow members, please share them with the group's leader, [Yunlong Sheng](#).

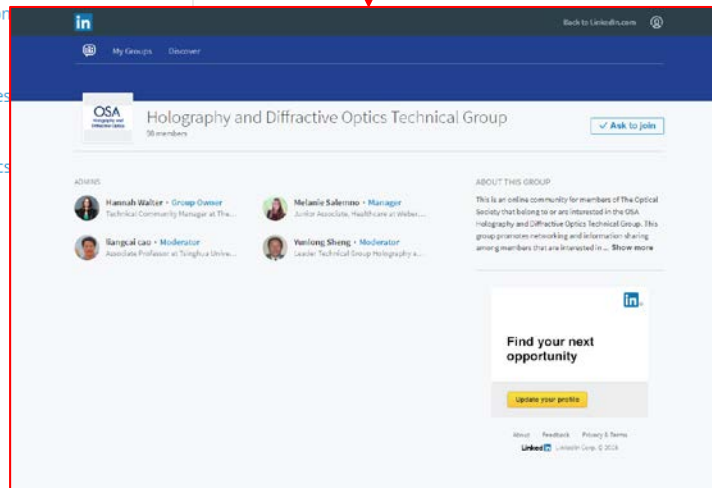
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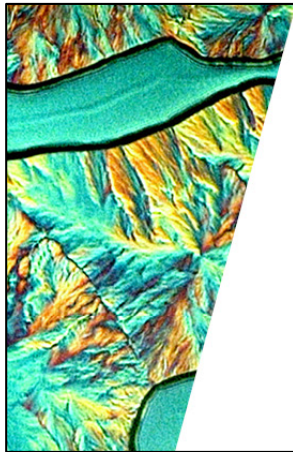
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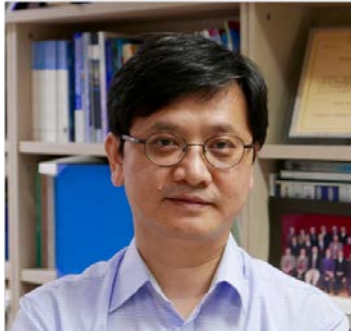




METASURFACE HOLOGRAMS WEBINAR

25 January 2018 • 19:00 EST

OSA Holography and
Diffractive Optics
Technical Group



Professor Byoung-ho Lee
Seoul National University, Korea

- Ph.D. in EECS from UC Berkeley (1993)
- Fellow of *OSA*, *SPIE*, and *IEEE*
- Member of *Korean Academy of Science and Technology* and *National Academy of Engineering of Korea*
- Former Director-at-Large of OSA, former chair of the Member and Education Services Council of OSA
- Former chair of Fabrication, Design & Instrumentation Technical Division
- Vice-President of Optical Society of Korea and the Korean Information Display Society
- Research fields: digital holography and 3D imaging, plasmonics and metasurfaces



Metasurface Holograms

ByoungHo Lee

byoungho@snu.ac.kr

*School of Electrical and Computer Engineering
Seoul National University, Republic of Korea*

January 25, 2018

I. Introduction

II. Metasurface holograms

III. Metalens and prospect of metasurface

IV. Conclusion

I. Introduction

II. Metasurface holograms

III. Metalens and prospect of metasurface

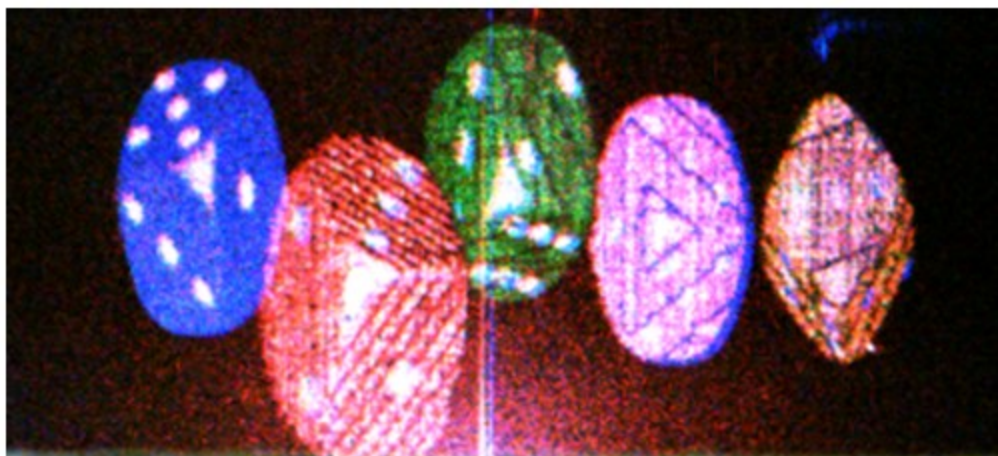
IV. Conclusion

Holography

- Holography reconstructs **wavefront** of light with **amplitude** and **phase** information the same as the original wave or as desired.



G. Li et al, *Opt. Lett.* **41**, 11 (2016).

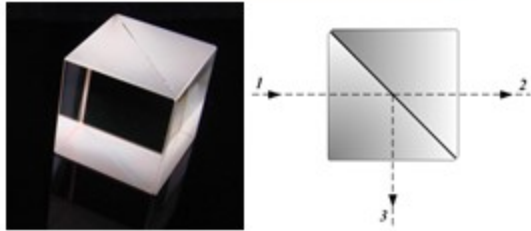


Examples: reconstructed red and full-color holograms using holographic optical element in our lab.

Optic components for light control

Directivity

Beam splitter

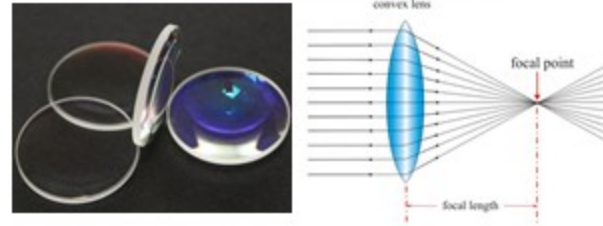


Mirror

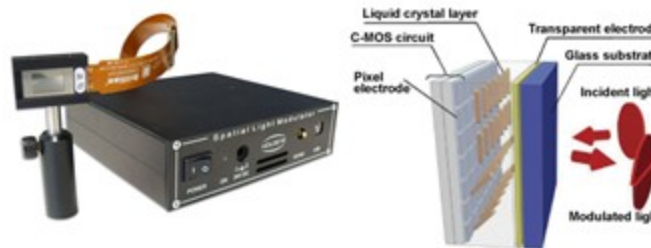


Phase or amplitude

Lens

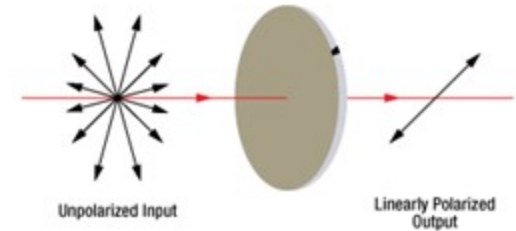


Spatial light modulator

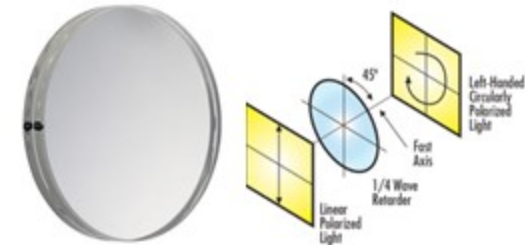


Polarization

Polarizer



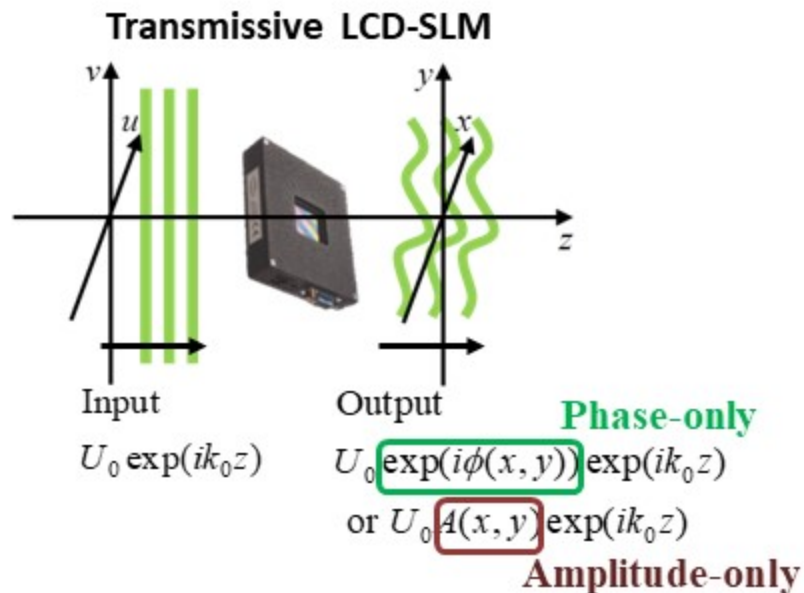
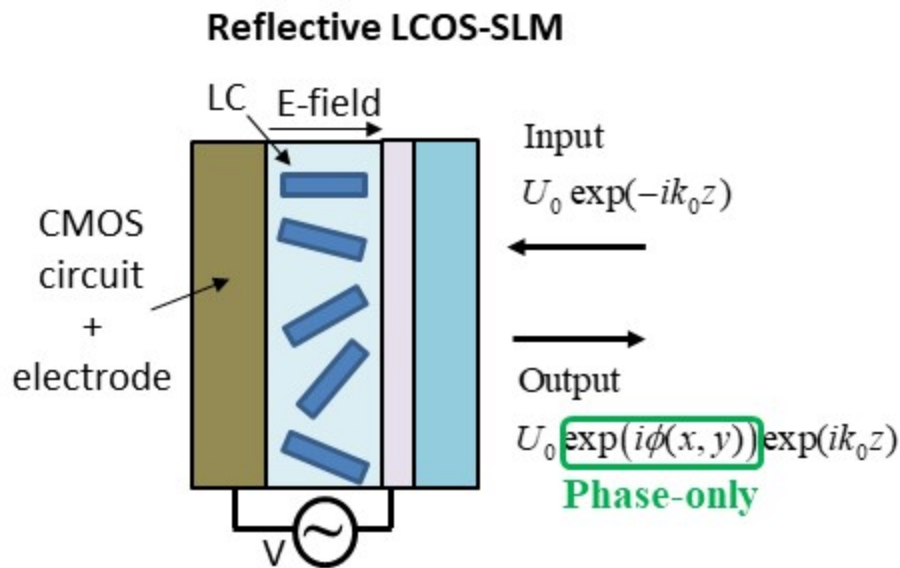
Waveplate



- Some limitations: large form factor, low resolution (SLM), narrow bandwidth

Spatial modulation of light

- Wavefront modulations for holography
- ➔ Spatial light modulators (SLMs) for **phase-only** / **amplitude-only** modulations



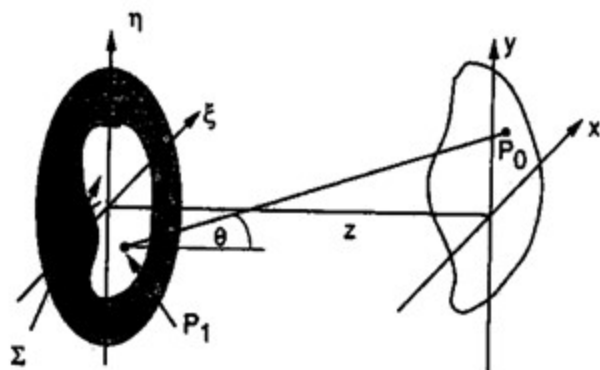
Limitations

- Phase **or** amplitude only modulation
- Large pixel pitch / grating period ($> \lambda$): low resolution, high diffraction order, narrow viewing angle in 3D hologram

Why metasurfaces?
➔ **Improved characteristics**
+
Miniaturization

Hologram generation: Fresnel wavefront calculation

- Fresnel diffraction theory: approximated scalar wave diffraction



$$\text{Assuming } z^3 \gg \frac{\pi}{4\lambda} [(x-\xi)^2 + (y-\eta)^2]_{\max}^2$$

→ Fresnel approximation

Fresnel diffraction integral

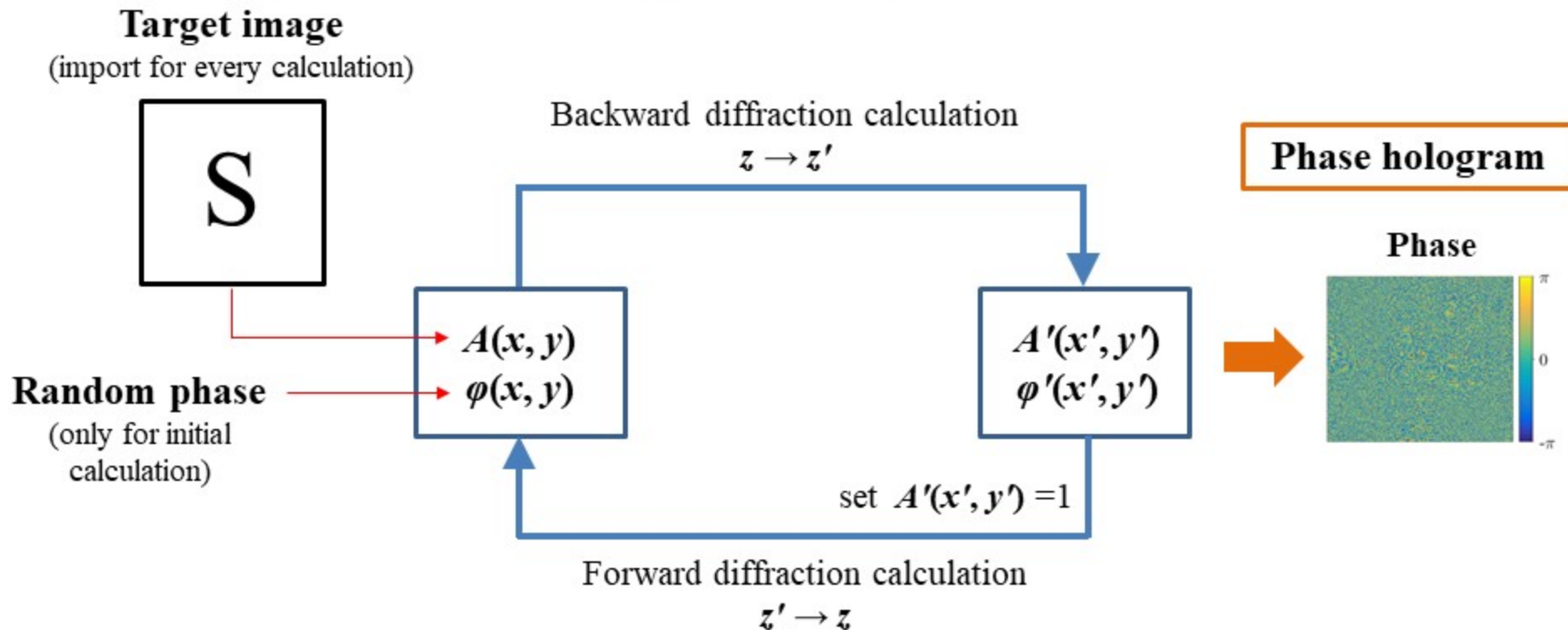
$$U(x, y) = \frac{e^{ikz}}{i\lambda z} \int \int_{-\infty}^{\infty} U(\xi, \eta) \exp \left\{ i \frac{k}{2z} [(x-\xi)^2 + (y-\eta)^2] \right\} d\xi d\eta$$

$$U(x, y) = \int \int_{-\infty}^{\infty} U(\xi, \eta) h(x-\xi, y-\eta) d\xi d\eta \quad h(x, y) = \frac{e^{ikz}}{i\lambda z} \exp \left[\frac{ik}{2z} (x^2 + y^2) \right]$$

$$U(x, y) = \frac{e^{ikz}}{i\lambda z} e^{i \frac{k}{2z} (x^2 + y^2)} \mathcal{F} \left\{ U(\xi, \eta) e^{i \frac{k}{2z} (\xi^2 + \eta^2)} \right\} \Bigg|_{f_x = x/\lambda z, f_y = y/\lambda z}$$

J. W. Goodman, Introduction to Fourier Optics, 4th ed., Freeman, 2017.

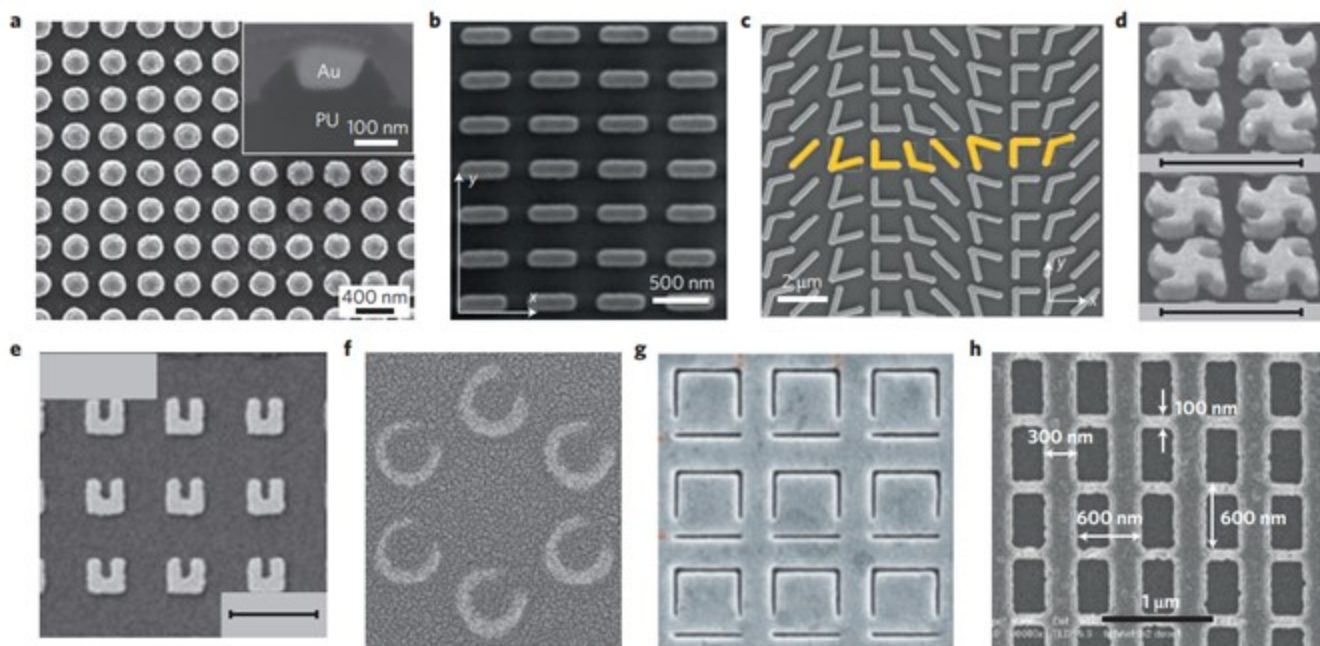
Hologram generation: Phase-only holograms (with GS algorithm)



- Phase-only holograms require proper approximations due to the absence of amplitude information (In general, both amplitude and phase information are required.)
- The Gerchberg-Saxton (GS) algorithm is a method to increase the image quality of phase-only holograms by sacrificing the phase information at the image plane.

Metasurfaces: periodically arranged meta-atoms

- Artificially fabricated ultrathin-film structure for controlling light



N. Meinzer et al., *Nat. Photon.* **8**, 891 (2014).

- Applications: **flat optic devices**, super-lensing, optical magnetism, nonlinear effects etc.

I. Introduction

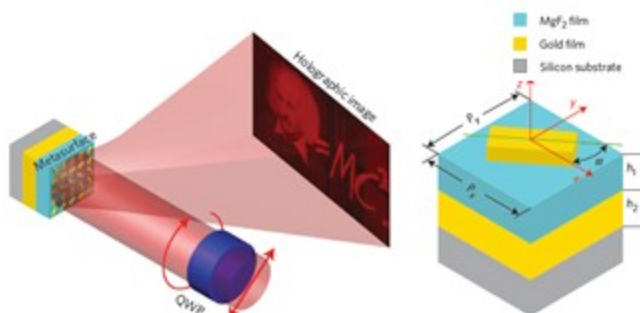
II. Metasurface holograms

III. Metalens and prospect of metasurface

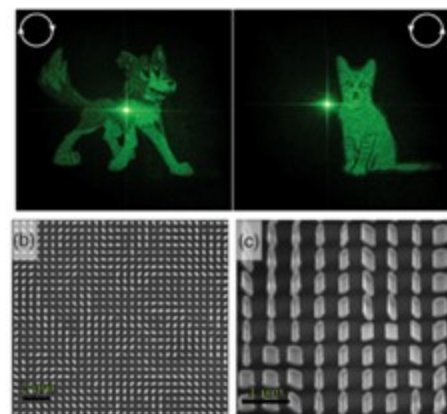
IV. Conclusion



X. Ni et al., *Nat. Comm.* 4, 2807 (2013).



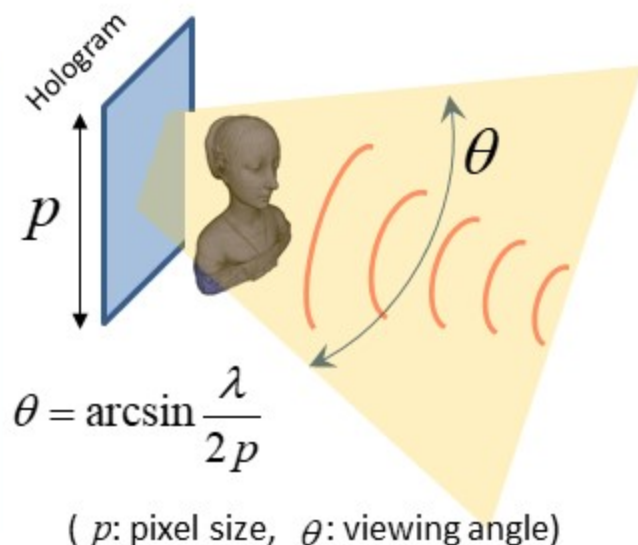
G. Zheng et al., *Nat. Nanotech.* 10, 308–312 (2015).



J. P. B. Muller et al., *Phys. Rev. Lett.* 118, 113901 (2017).

➤ Why metasurface holograms?

- Subwavelength resolutions (~ 350 nm)
- Large viewing angle and higher information density
- Suppressed high-order diffraction cross-talk
- Ultracompact (thickness $\leq \lambda$)

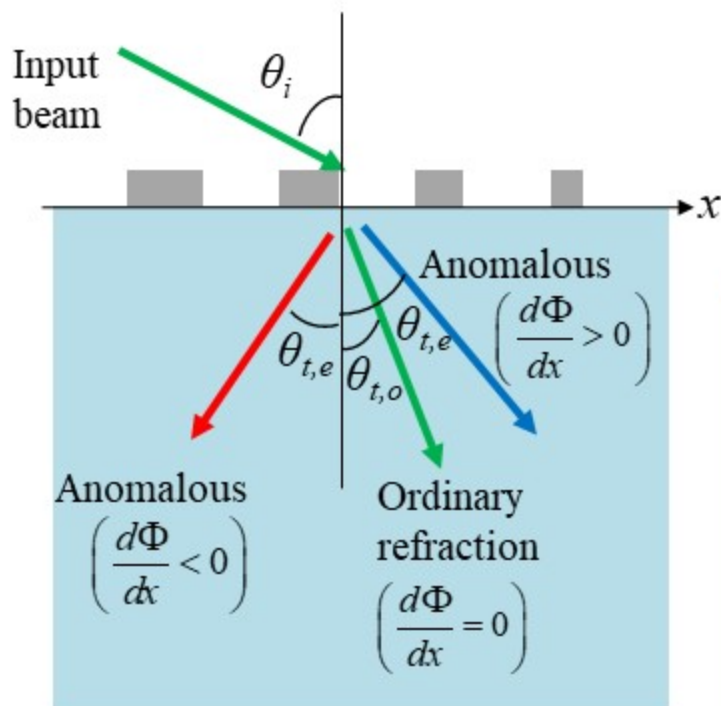




Abrupt phase modulation of wave by metasurface

0. Breaking the Snell's law: abrupt phase modulation by ultrathin metasurface

- Generalized Snell's law



N. Yu et al., *Science* 334, 6054 (2011).

$$\sin \theta_t = \frac{n_i}{n_t} \sin \theta_i + \frac{1}{n_t k_0} \frac{d\Phi}{dx}$$

Abrupt phase gradient

- Abrupt phase gradient can be imparted by ultrathin metasurface.
- Generalized Snell's law enables **anomalous** refraction of light.
- Spatial arrangement of proper meta-atoms
→ phase modulating metasurface

Half-wave plate: absolute conversion of optical handedness

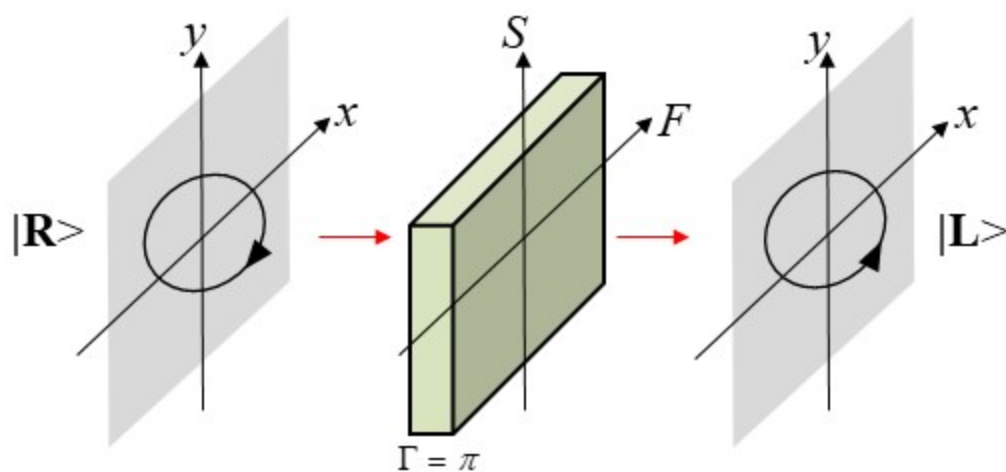
$$\Gamma = \pi$$

$$\mathbf{T} = \begin{pmatrix} 1 & 0 \\ 0 & \exp(-j\Gamma) \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$E_t = \mathbf{T} \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \exp(-j\Gamma) \end{pmatrix} \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} = \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix}$$

$$\sigma = 1 \quad \mathbf{RCP}$$

$$\sigma = -1 \quad \mathbf{LCP}$$

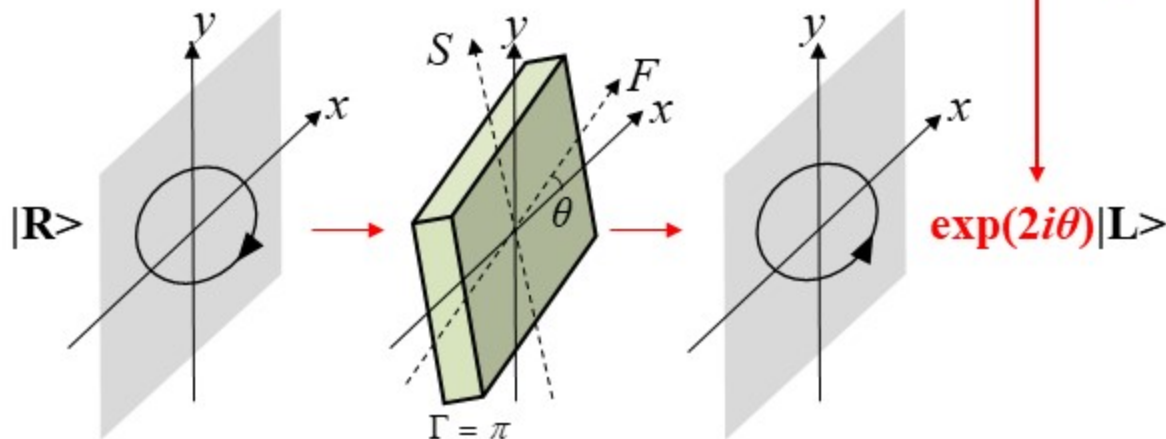


Rotated half-wave plate: **Pancharatnam-Berry (PB) phase** with conversion of handedness

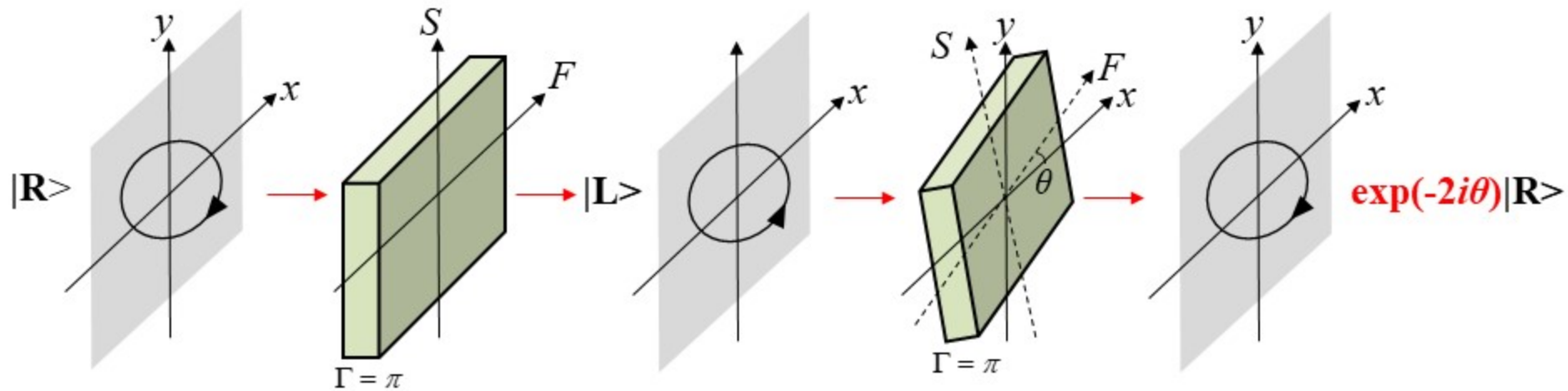
$$\begin{aligned} \Gamma &= \pi \\ \mathbf{T} &= R(-\theta) \begin{pmatrix} 1 & 0 \\ 0 & \exp(-j\Gamma) \end{pmatrix} R(\theta) \\ &= \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \\ &= \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix} \end{aligned}$$

$$E_t = T \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} = \exp(i2\sigma\theta) \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix}$$

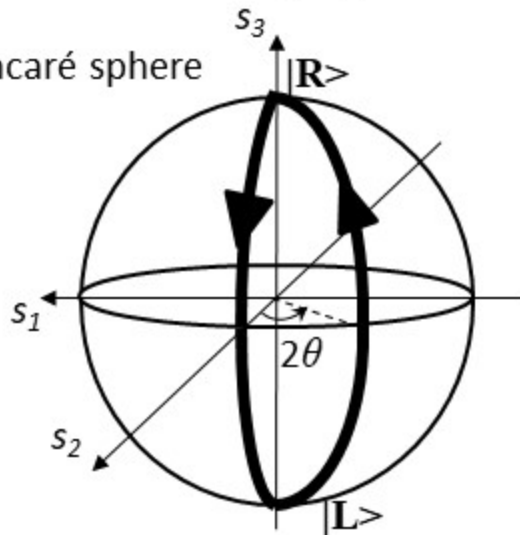
PB phase



Rotated half-wave plate: Pancharatnam-Berry (PB) phase



- Poincaré sphere



$$\phi_g = \frac{\Omega}{2} = \frac{1}{2} \int_0^{2\theta} \int_0^\pi \sin \phi \, d\phi \, d\theta = 2\theta$$

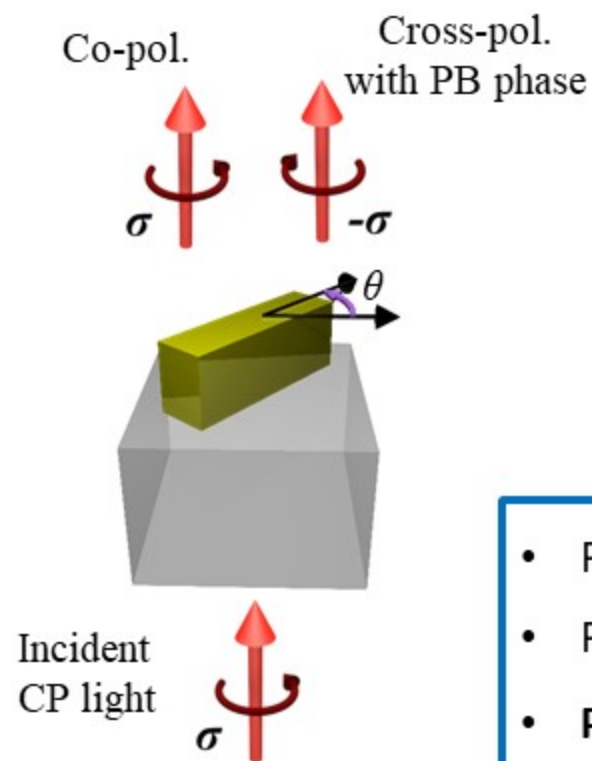
- Phase difference between input RCP light and output RCP light is determined by a rotation angle.
- The traversed surface areas = Ω
- PB phase (geometric phase) = $\phi_g = \Omega/2 = 2\theta$

N. Shitrit et al., *Opt. Lett.* 38, 21 (2013).

Methods of abrupt phase modulation by metasurfaces

1. Anisotropic phase modulation: Pancharatnam-Berry (PB) phase

- A meta-atom acting as a miniature elliptical polarizer



$$\begin{aligned}
 E_t &= R(-\theta) \begin{pmatrix} t_L & 0 \\ 0 & t_S \end{pmatrix} R(\theta) \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} \\
 &= \frac{t_L + t_S}{2} \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} + \frac{(t_L - t_S)}{2} \exp(2i\sigma\theta) \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix} \\
 &= E_{t,\sigma} + E_{t,-\sigma}
 \end{aligned}$$

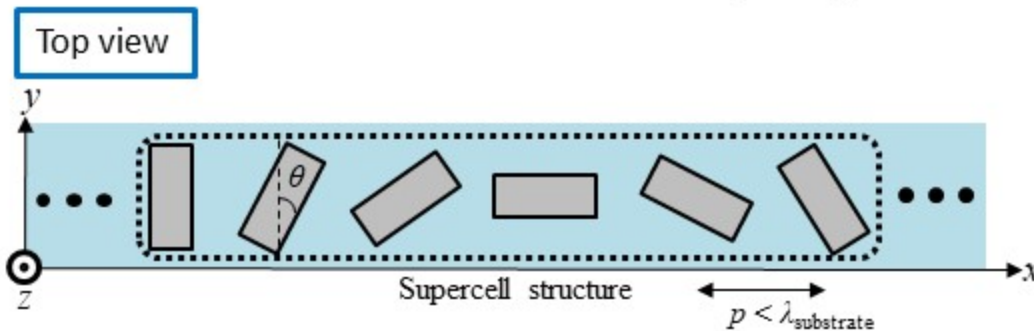
Large anisotropy is required for high cross-pol. modulation efficiency.

- Phase modulation using **circular polarization basis**
- Phase control of **cross-pol.** by rotation (θ)
- **PB phase ($2\sigma\theta$)** can cover full 2π phase by rotation of 180° .
- Wavelength-independent method: **broadband**

Methods of abrupt phase modulation by metasurfaces

1. Anisotropic phase modulation: Pancharatnam-Berry (PB) phase

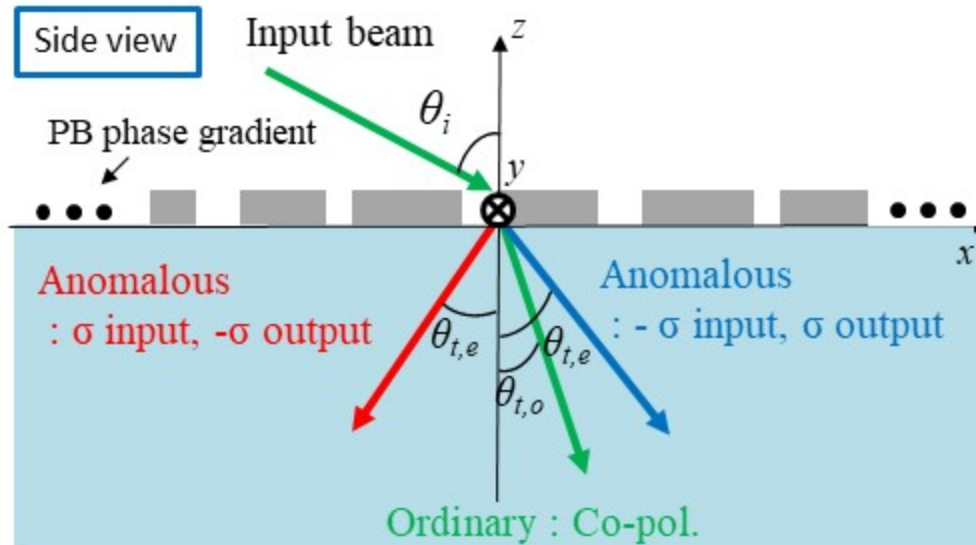
➤ Anomalous refraction by PB phase



$$\Phi = 2\sigma\theta$$

$$\sin \theta_t = \frac{n_i}{n_t} \sin \theta_i + \sigma \frac{\lambda_0}{\pi} \frac{d\theta}{dx}$$

$$\sigma = \begin{cases} 1 & \text{(RCP)} \\ -1 & \text{(LCP)} \end{cases}$$

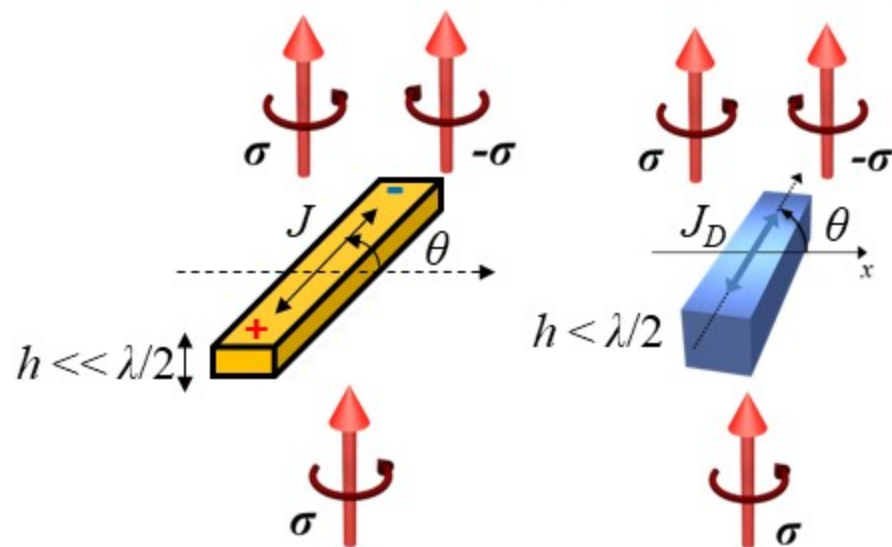


- PB phase gradient is opposite for two different input circular polarizations.
- Two different anomalous refractions occur **according to input handedness**.
- Regardless of input handedness, co-pol. light is not affected by phase gradient.

Methods of abrupt phase modulation by metasurfaces

1. Anisotropic phase modulation: Pancharatnam-Berry (PB) phase

➤ Implementation: **metallic / high-index nanorod**



$$E_t \approx R(-\theta) \begin{pmatrix} t & 0 \\ 0 & 0 \end{pmatrix} R(\theta) \begin{pmatrix} 1 \\ i\sigma \end{pmatrix}$$

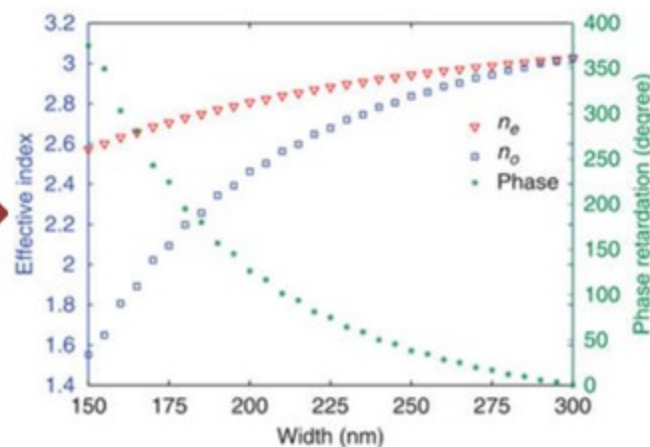
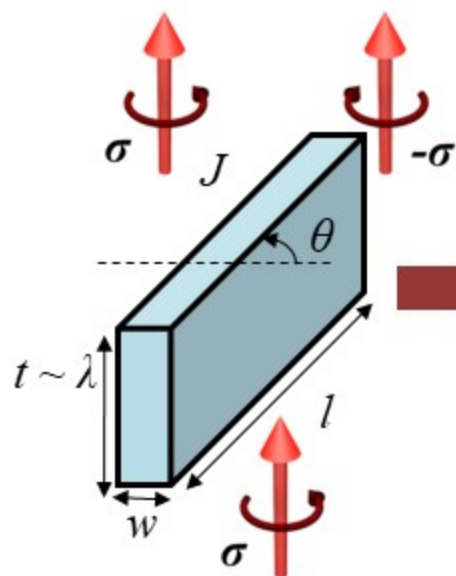
$$\approx \frac{t}{2} \left\{ \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} + e^{2i\sigma\theta} \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix} \right\}$$

- Plasmonic / high-index nanorod: large scattering cross-section
- Nearly acting as a miniature linear polarizer
- Ultra-thin thickness of nanorods $< \lambda/2$
- Electric dipole is induced along the long axis. $\rightarrow t_L = t$ and $t_S \approx 0$

Methods of abrupt phase modulation by metasurfaces

1. Anisotropic phase modulation: Pancharatnam-Berry (PB) phase

➤ Implementation: **low loss moderate-index nanofin (birefringent waveguide)**



$$E_t = \frac{t_L + t_S}{2} \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} + \frac{t_L - t_S}{2} e^{2i\sigma\theta} \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix}$$

$$t_L \approx a \exp\left(i \frac{2\pi}{\lambda_0} n_e t\right) \text{ and } t_S \approx b \exp\left(i \frac{2\pi}{\lambda_0} n_o t\right)$$

(a, b : real constants)

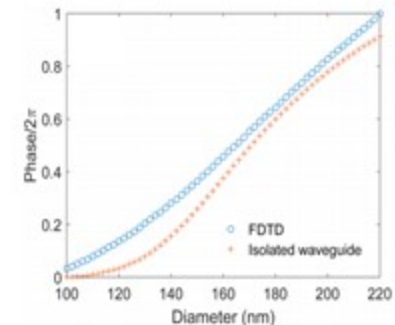
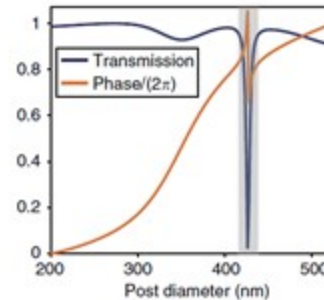
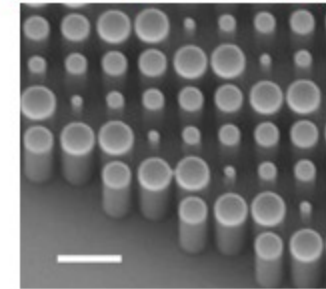
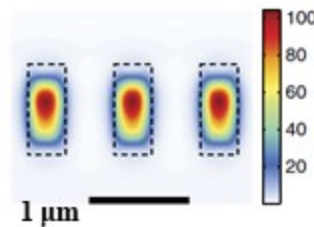
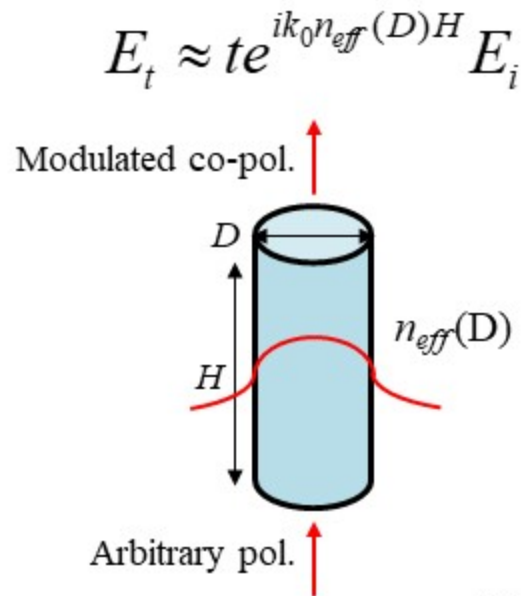
M. Khorasaninejad et al., *Nat. Commun.* 5, 5386 (2014).

- Dielectric nanofin : short birefringent **dual mode** (n_e, n_o) waveguide
- Higher-order modes are suppressed by subwavelength thickness and normal incidence.
- Acting as a miniature elliptical polarizer
- Nearly λ -thick nanofins required for enough phase retardation (difference between t_e and t_o)
- Fabry-Perot resonances can be avoided by proper selection of height.

Methods of abrupt phase modulation by metasurfaces

2. Isotropic phase modulation: control of **local meta-atom dimensions**

➤ Implementation: **low loss moderate-index dielectric nano-cylinder (isotropic waveguide)**



A. Arbabi et al., *Nat. Commun.* 6, 7069 (2015). M. Khorasaninejad et al., *Nano Lett.* 16, 11 (2016).

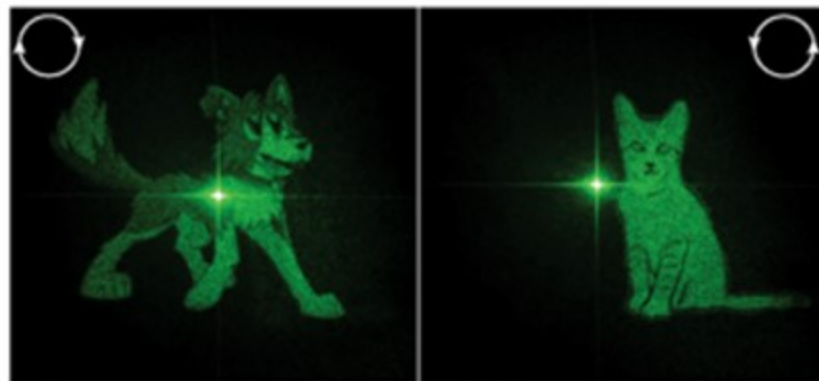
- Dimension of isotropic nano-cylinder is controlled for full 2π phase modulation.
- Dimension determines effective mode index of the fundamental waveguide mode.
- Subwavelength period \rightarrow Input light is mostly coupled to waveguide mode.
- Multi-wavelength operation is possible by avoiding FP resonant height condition.

Recent demonstrations of metasurface holograms

Anisotropic
PB phase
meta-atom
based

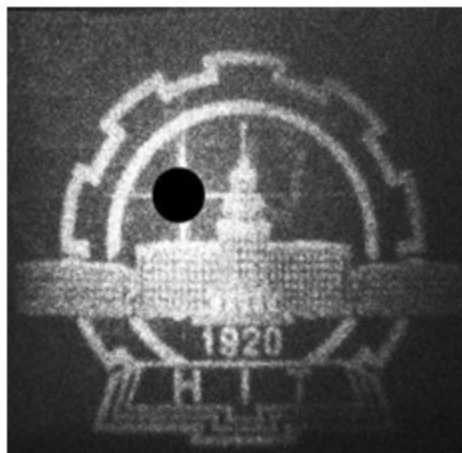


G. Zheng et al., *Nat. Nanotech.* 10, 308–312 (2015).



J. P. B. Muller et al., *Phys. Rev. Lett.* 118, 113901 (2017).

Isotropic
dielectric
meta-atom
based



W. Zhao et al., *Sci. Rep.* 6, 30613 (2016).



X. Su et al., *RSC Adv.* 7, 26371 (2017).

Why is complex modulation in holography important?

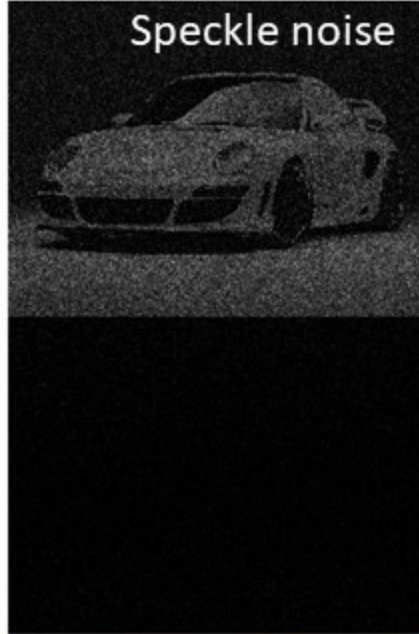
Complex



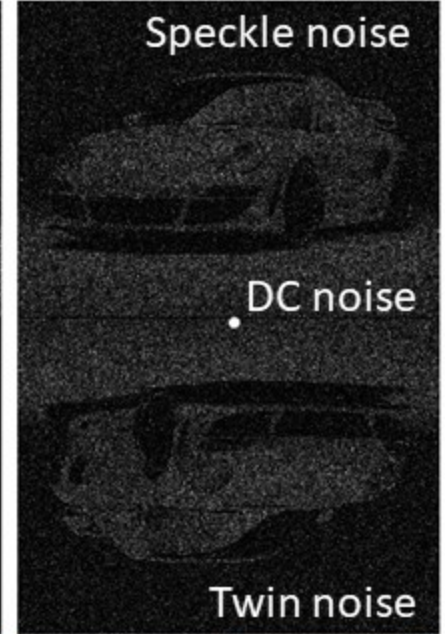
Amplitude only



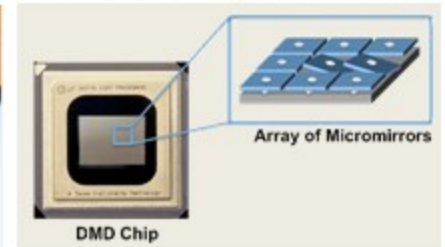
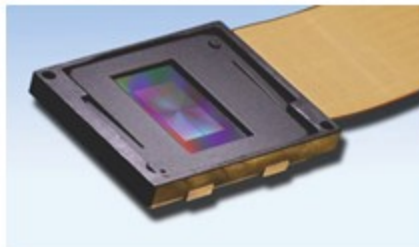
Phase only



Amplitude only
(binary)



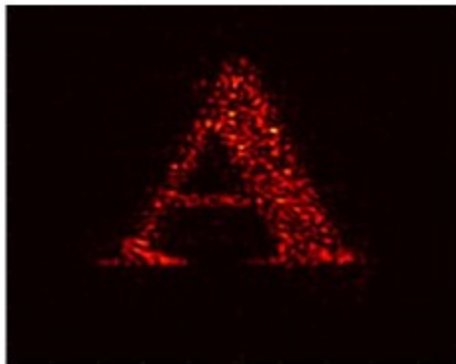
?



Metasurface for complex modulation and holograms

Comparison

Dipole modeling metasurface hologram simulations



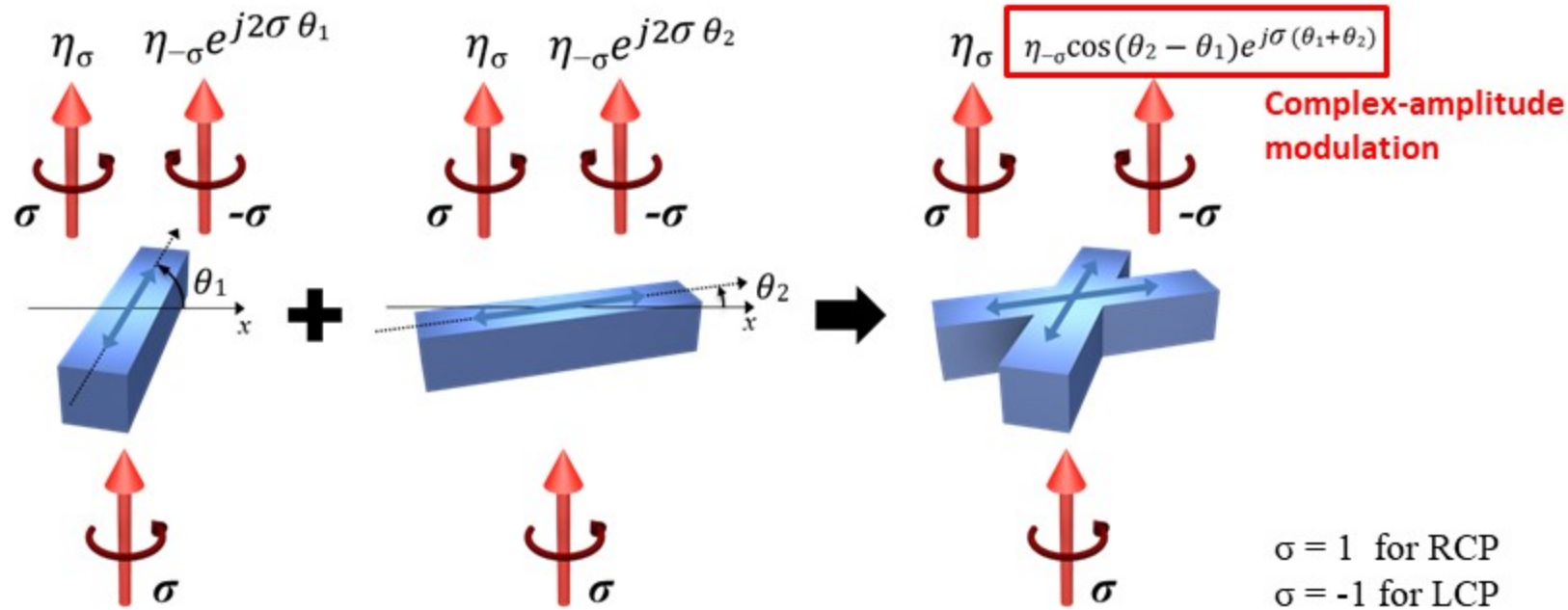
Phase-only hologram



Complex hologram

- Conventional methods using phase-only SLM or amplitude-only SLM
 - Small viewing angle ($\sim 5^\circ$) and restricted image size
 - Filtering of twin image is required.
- Complex modulation by metasurface can overcome these limits.

X-shaped silicon meta-atoms based on the PB phase



- Two crossed silicon nanorods \rightarrow crossed nearly-independent electric dipoles

$$E_{-\sigma} \propto e^{j2\sigma\theta_1} + e^{j2\sigma\theta_2} = 2 \cos(\theta_2 - \theta_1) e^{j\sigma(\theta_1 + \theta_2)}$$

Complex-amplitude of cross-polarized transmission

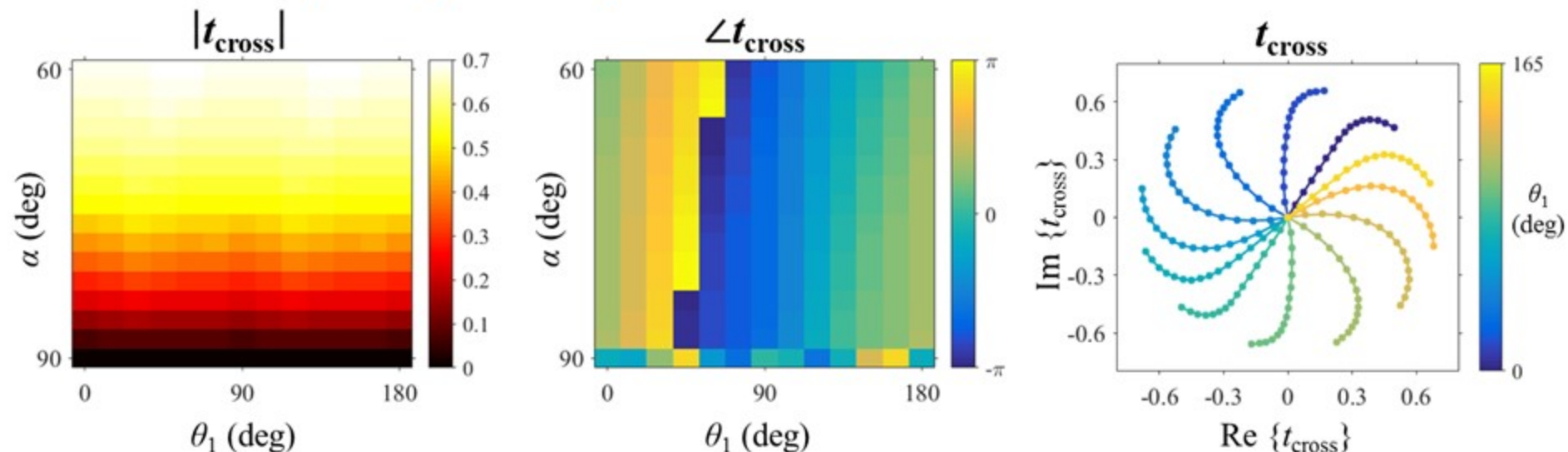


$$A = \cos(\theta_2 - \theta_1) \quad \text{(amplitude)}$$

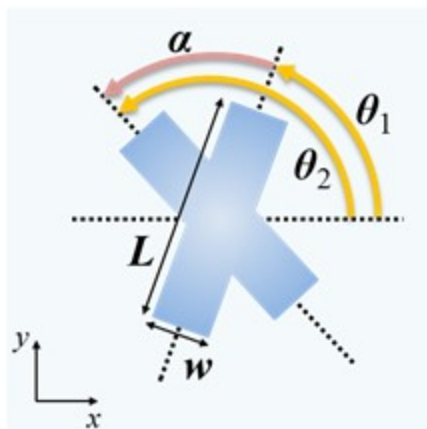
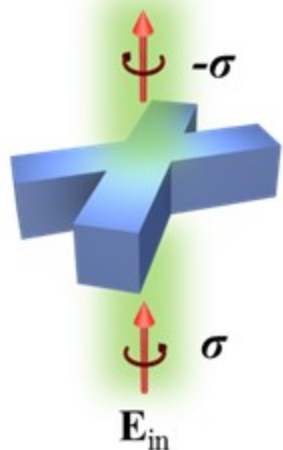
$$\phi = \theta_2 + \theta_1 \quad \text{(phase)}$$

- Full ranges of both amplitude and phase can be described.

Modulation capability of X-shaped meta-atoms

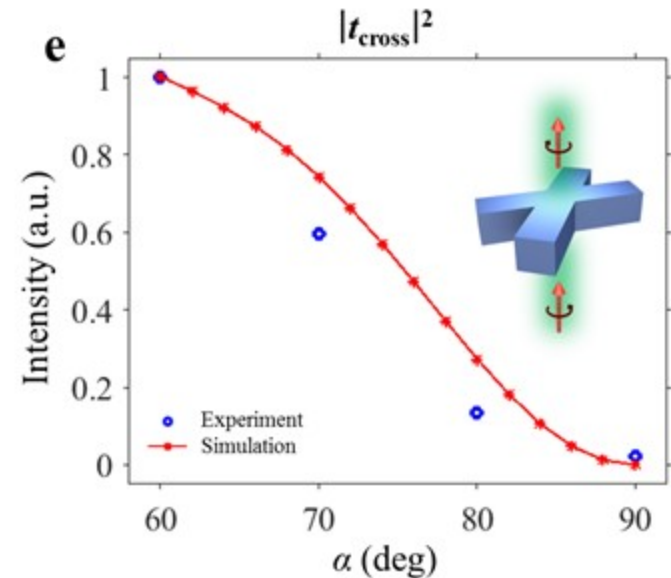
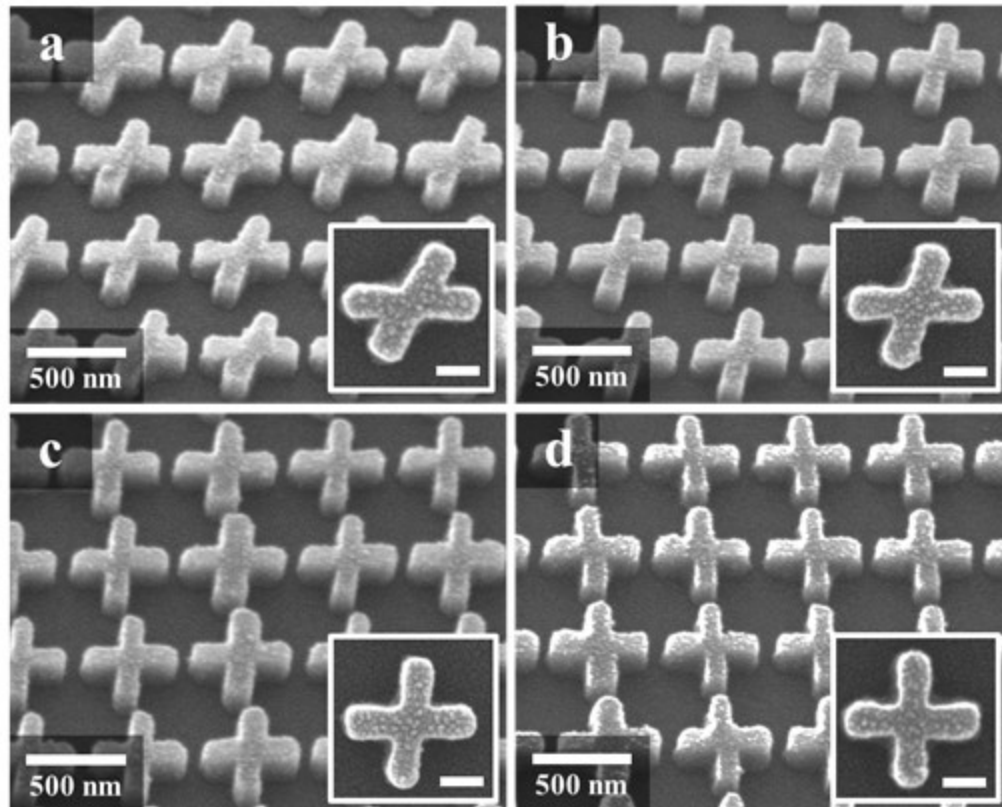


$$\mathbf{E}_{\text{out}} = t_{\text{cross}} \mathbf{E}_{\text{in}}$$



- t_{cross} is calculated for various disparity angles (α) and orientation angles (θ_1).
- Disparity angle (α) controls the amplitude while orientation angle (θ_1) controls the phase.
- A full range of complex-amplitude is continuously covered on complex domain.

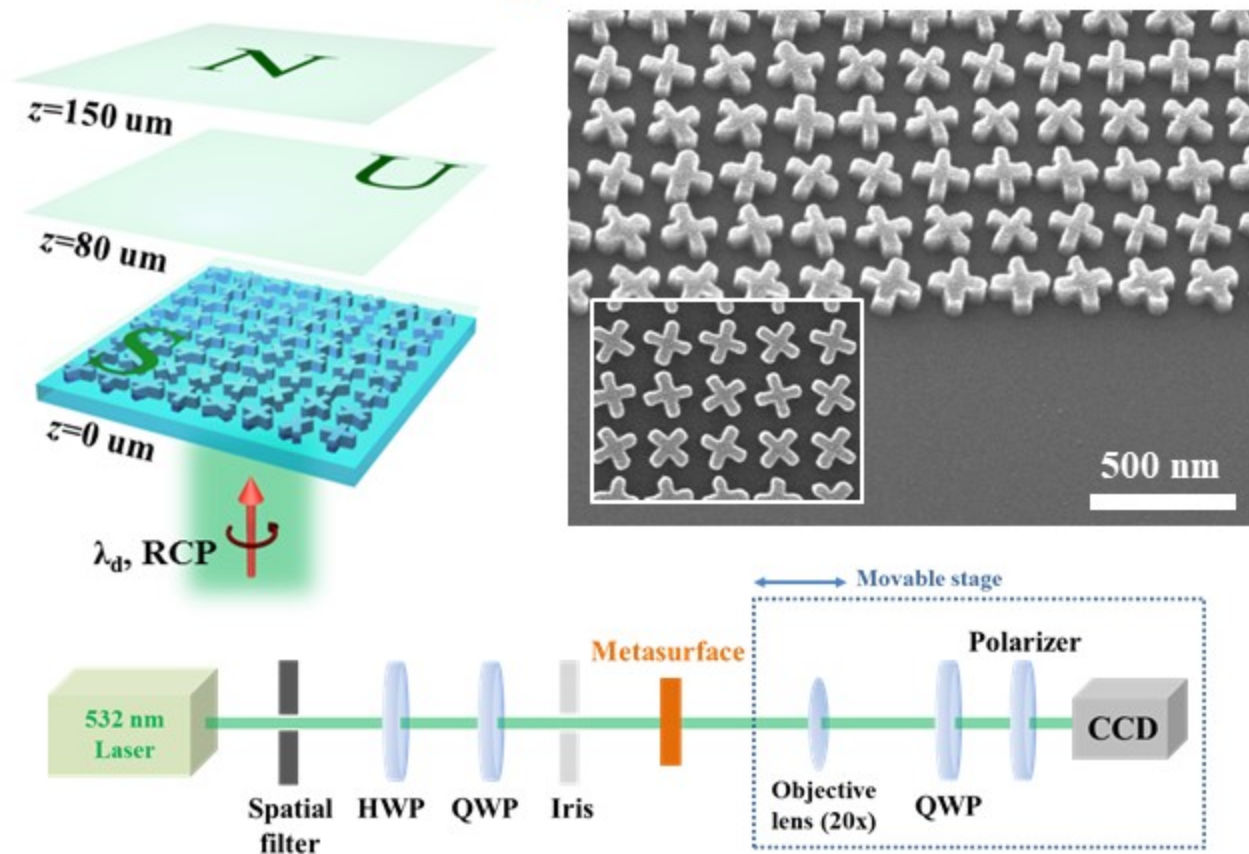
Experimentally measured cross-polarized transmissions (t_{cross})



- Standard electron beam lithography process was used. (co-work with prof. Junsuk Rho at POSTECH)
- Experimental results show a good agreement with the simulation results.
- As expected, disparity angle (α) controls the amplitude components of t_{cross} .



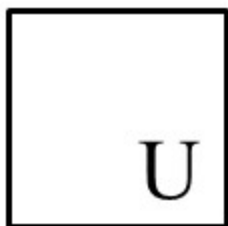
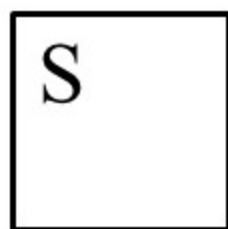
Complex-amplitude metasurface holograms



- As an example, hologram with letters of “SNU” is designed without any approximations.
- “S” is at metasurface plane ($z=0 \mu\text{m}$), and “U” and “N” are at $z=80$ and $150 \mu\text{m}$.
- Reconstructed images are measured by using CCD camera with an objective lens.

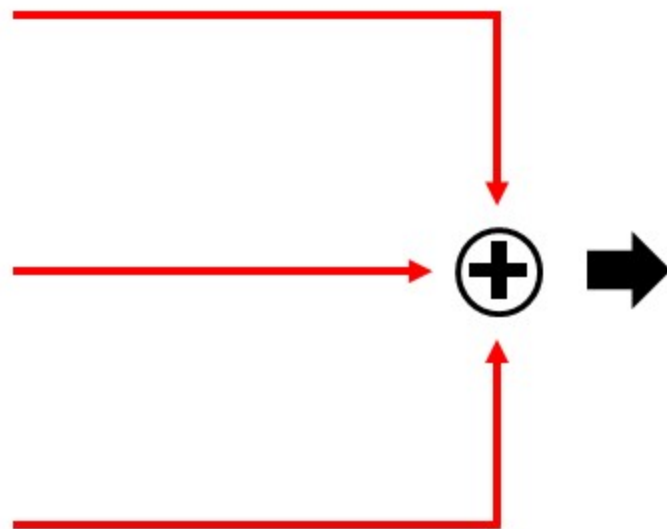
Complex holograms with various depth information

Target images



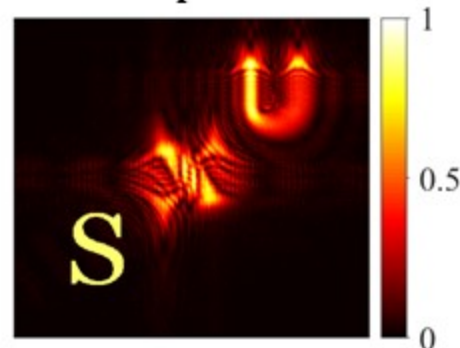
Backward diffraction calculation

$$z \rightarrow z'$$

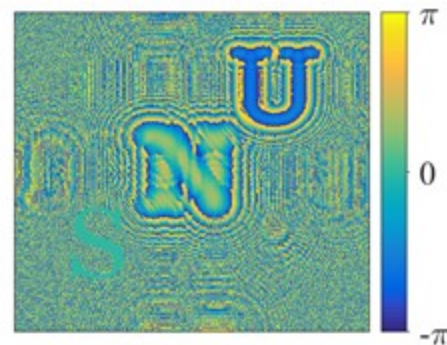


Complex hologram

Amplitude



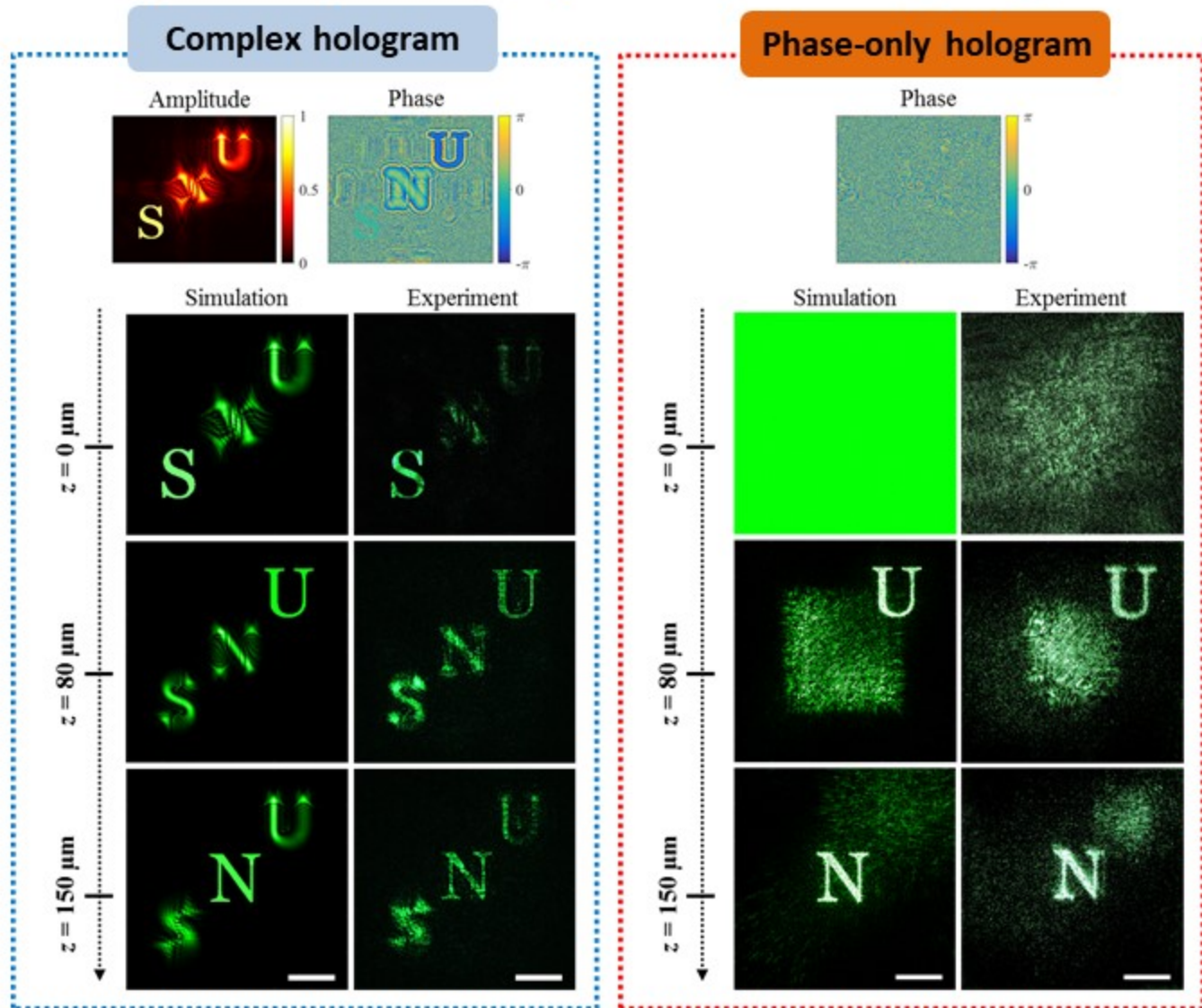
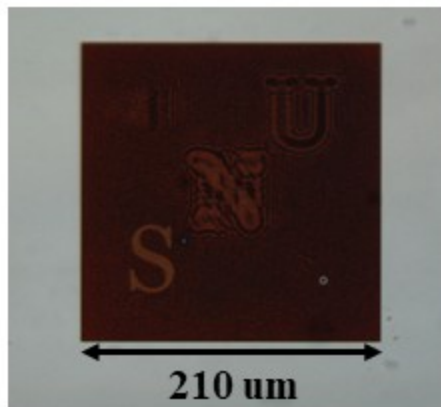
Phase



- Compared to phase-only holograms, a complex hologram does not need any approximations or scarifications of information on its calculation.
- Angular spectrum method was employed to generate holograms.
- Finally, the complex hologram for the three letters with different depths is calculated by the sum of the holograms of each image (S, N, U).



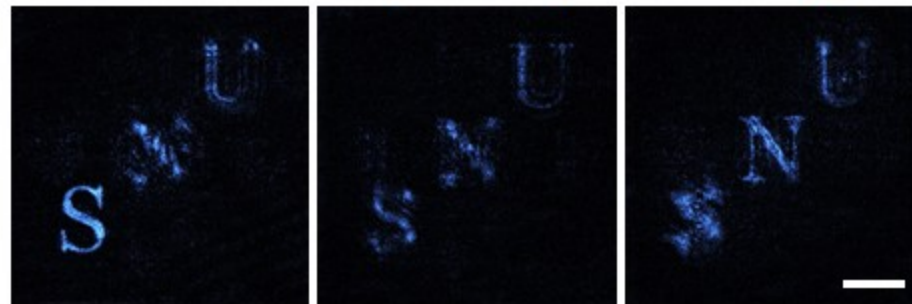
Experimental results: Complex-amplitude hologram vs. phase-only holograms





Experimental results: Broadband complex-amplitude hologram

$\lambda = 473 \text{ nm}$



$z_1 = 0 \mu\text{m}$

$z_2 = 90 \mu\text{m}$

$z_3 = 169 \mu\text{m}$

$\lambda = 660 \text{ nm}$



$z_1 = 0 \mu\text{m}$

$z_2 = 64 \mu\text{m}$

$z_3 = 121 \mu\text{m}$

- Due to the broadband property of the PB phase, X-shaped meta-atoms well operate in broadband visible region.
- Full complex-amplitude holograms show very clear images with a record-breaking signal-to-noise ratio (SNR = 211.3) where previously reported SNR for phase holograms are about 50.

G.-Y. Lee et al., "Complete amplitude and phase control of light using broadband holographic metasurface," *Nanoscale*, 2018.

I. Introduction

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IV. Conclusion

Optical applications & optic components

➤ Smartphone camera



Samsung Galaxy S8

➤ Augmented Reality (AR) device



Microsoft HoloLens

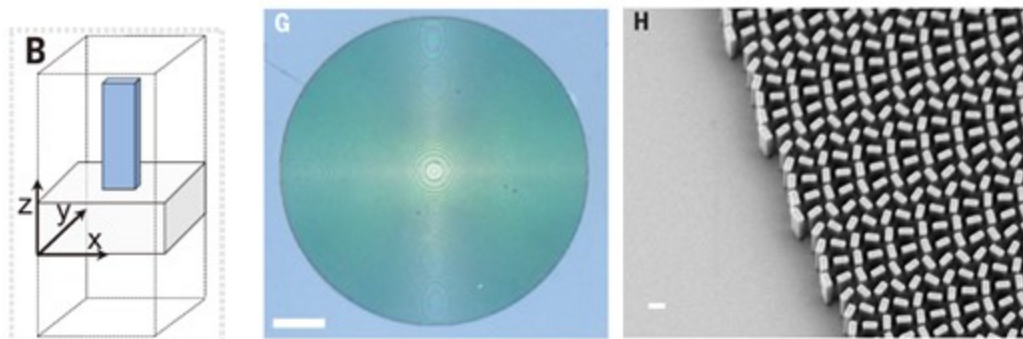
➤ Virtual Reality (VR) device



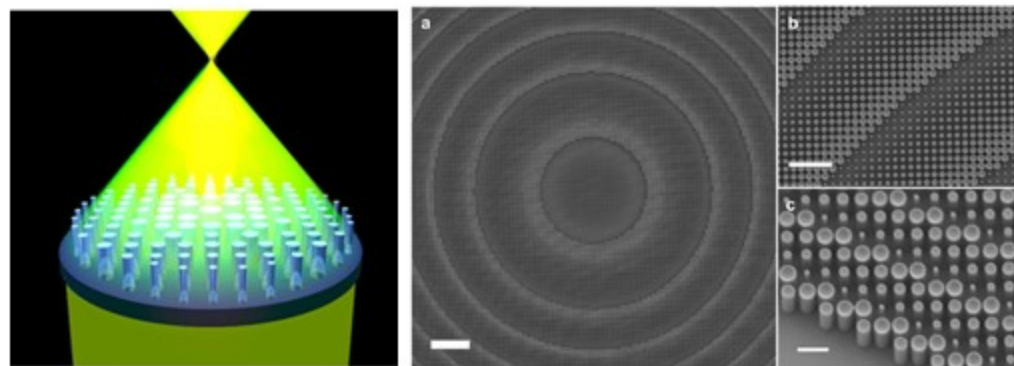
Samsung Gear VR

- There are various mobile optical applications using a series of optic components.
- Miniaturization of optic components is essential for advanced mobile applications.

Metasurface lens (metalens) : overview



M. Khorasaninejad et al., *Science* 352, 6290 (2016).



M. Khorasaninejad et al., *Nano Lett.* 16, 7229–7234 (2016).

- Metalens concept overcomes the limitations of conventional bulk-optic lenses including large form factor and low numerical aperture.
- Recent advances in metalens have shown high efficiency (>80 %) and high NA (NA=0.8) within ultrathin thickness (<500 nm).

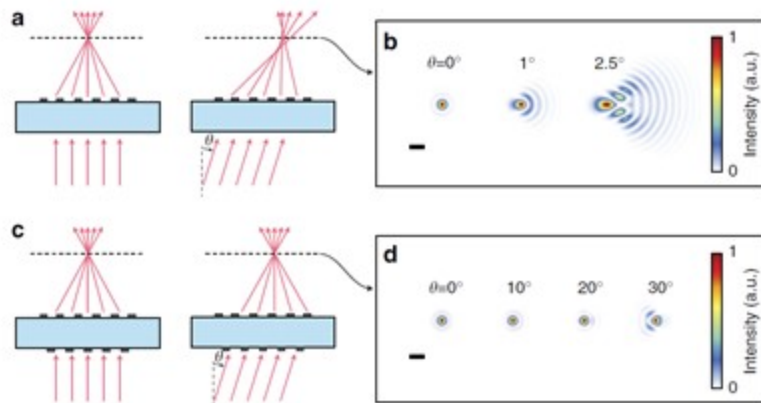
Study on metalens

Reference	Efficiency	Polarization	Wavelength (nm)
F. Aieta <i>et al.</i> ¹	~1%	Linear	1550
X. Ni <i>et al.</i> ²	~10%	Linear	476-676
A. Arbabi <i>et al.</i> ³	~62-74%	Insensitive	1450-1550
S. Vo <i>et al.</i> ⁴	~70%	Insensitive	850
D. Lin <i>et al.</i> ⁵	Not reported	Circular	550
M. Khorasanineja <i>et al.</i> ⁶	Above 50%	Circular	405, 532, 660
M. Khorasanineja <i>et al.</i> ⁷	~30%, ~70%, ~90%,	Insensitive	405, 532, 660

1. F. Aieta *et al.*, *Nano Lett.* 12, 4932–4936 (2012).
2. X. Ni *et al.*, *Light Sci. Appl.* 2, e72 (2013).
3. A. Arbabi *et al.*, *Nat. Commun.* 6, 7069 (2015).
4. S. Vo *et al.*, *IEEE Photon. Technol. Lett.* 26, 1375–1378 (2014).
5. D. Lin *et al.*, *Science* 345, 298–302 (2014).
6. M. Khorasaninejad *et al.*, *Science* 352, 6290 (2016).
7. M. Khorasaninejad *et al.*, *Nano Lett.* 16, 7229–7234 (2016).

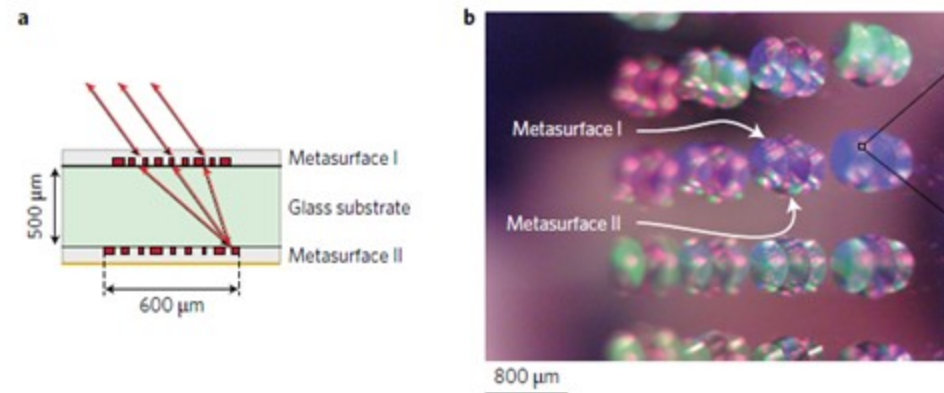
Metalens applications

➤ Metalens free from off-axis aberration



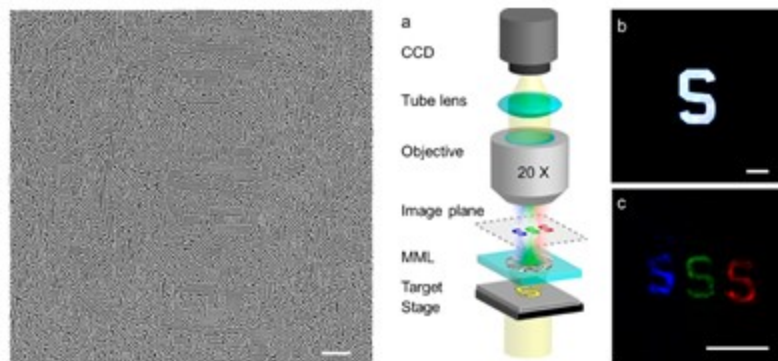
A. Arbabi et al., *Nat. Commun.* 7, 13682 (2016)

➤ Retroreflector



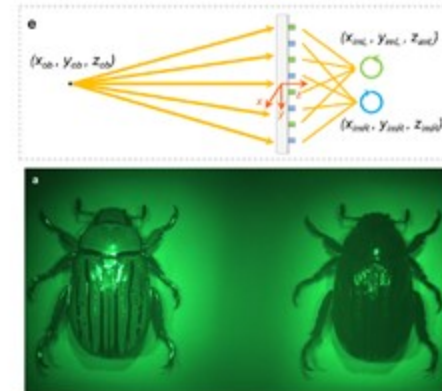
A. Arbabi et al., *Nat. Photonics.* 11, 415-420 (2017).

➤ Spectral imaging



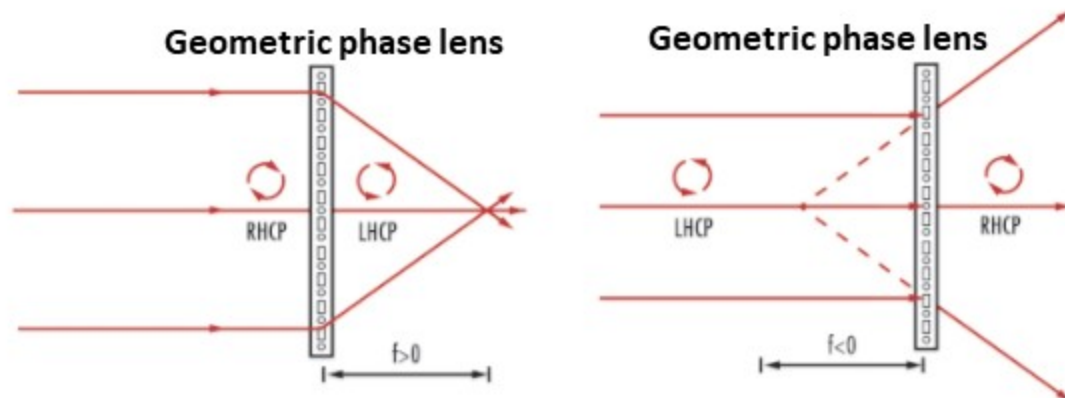
D. Lin et al., *Nano Lett.* 16, 7671-7676 (2016)

➤ Chirality-sensitive imaging

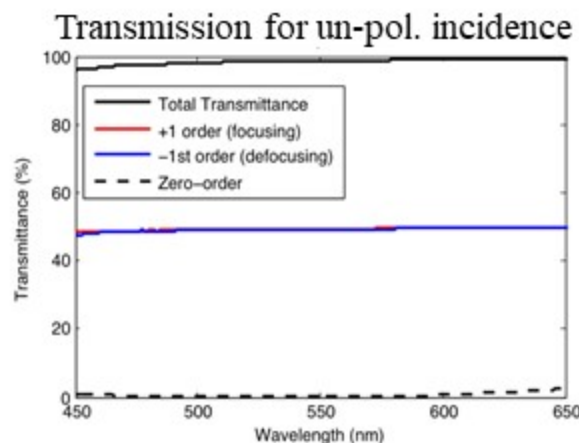


M. Khorasaninejad et al., *Nano Lett.* 16, 4595-4600 (2016).

➤ Geometric phase lens (PB phase lens) using liquid crystals (commercialized)



<http://www.imagineoptix.com/technology/geometric-phase-lens/>



Rotated half-wave plate

$$E_t = T \begin{pmatrix} 1 \\ i\sigma \end{pmatrix} = \exp(i2\sigma\theta) \begin{pmatrix} 1 \\ -i\sigma \end{pmatrix}$$

Geometric phase (PB phase)

Generalized Snell's law

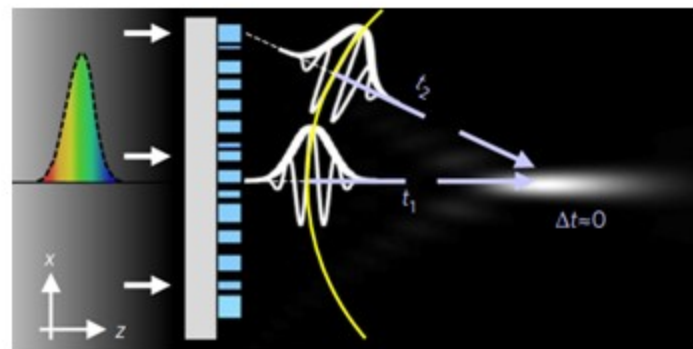
$$\sin \theta_t = \frac{n_i}{n_t} \sin \theta_i + \sigma \frac{\lambda_0}{\pi} \frac{d\theta}{dx}$$

$$\sigma = \begin{cases} 1 & \text{(RCP)} \\ -1 & \text{(LCP)} \end{cases}$$

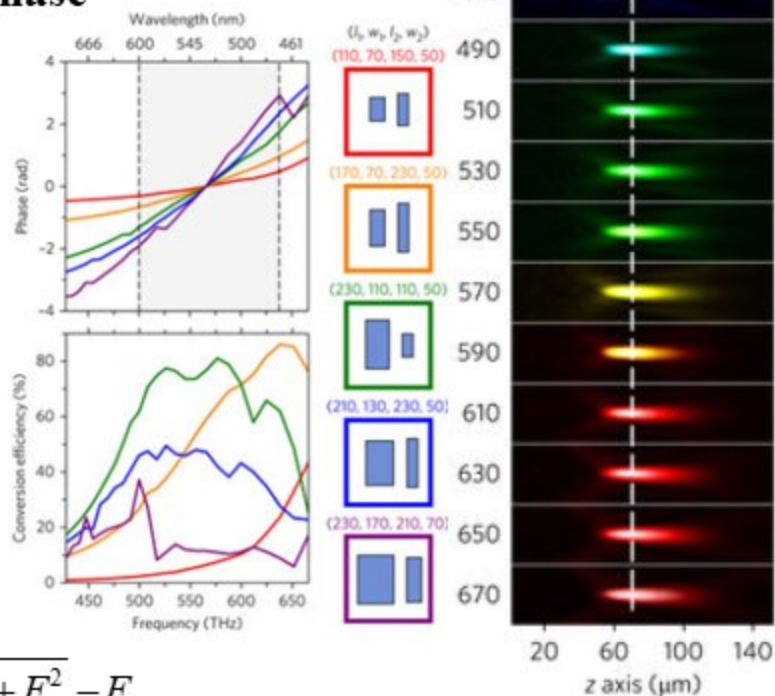
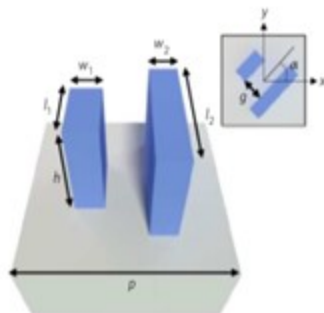
- **Dual-polarity** focusing according to optical handedness
- Extremely high transmission
- Commercially available now

Future of metallens: achromatic thin lens

➤ Broadband achromatic metallens based on PB phase



W. T. Chen, et al., Nat. nanotech. 1, 2018

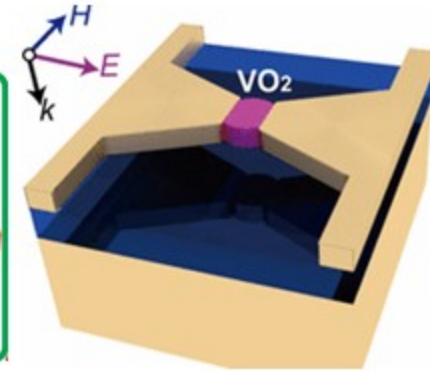
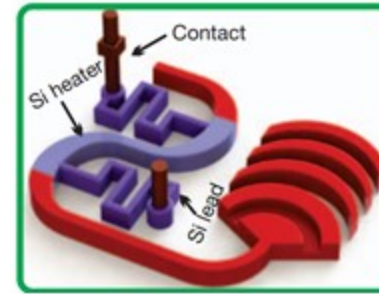
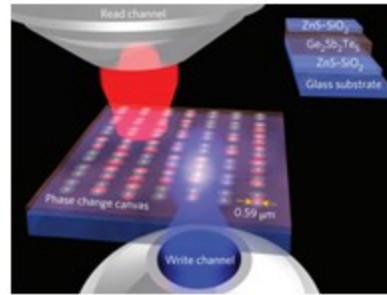
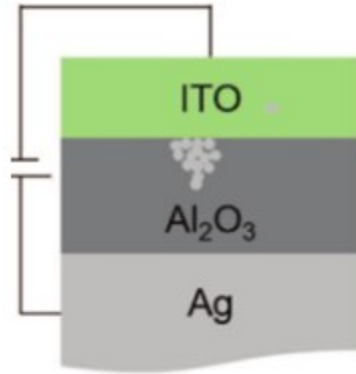


$$\varphi(r, \omega) = -\frac{\omega}{c} \left(\sqrt{r^2 + F^2} - F \right) \approx \varphi(r, \omega_0) + \left. \frac{\partial \varphi(r, \omega)}{\partial \omega} \right|_{\omega=\omega_0} (\omega - \omega_0)$$

$$\Rightarrow \varphi(r, \omega) \approx \exp(in_{\text{eff}}(r, \omega)k_0h) \exp(2i\theta(r)) \quad \& \quad \frac{\partial \varphi(r, \omega)}{\partial \omega} \approx -\frac{\sqrt{r^2 + F^2} - F}{c}$$

- Coupled dielectric nanofins (\rightarrow truncated short TiO_2 waveguide)
- Investigation of $n_{\text{eff}}(r, \omega)$ and transmittance of nanofin for aberration compensation
- Group delay is designed to be nearly non-dispersive for mitigating aberration.
- Visible wavelengths (470-670 nm)
- NA ~ 0.2 (@ $f = 63 \mu\text{m}$) and focusing efficiency up to 20%

Future of metasurfaces: active light modulation



Tuning mechanism	Field-effect ionic conductance	Optical pulse	Thermo-optic	Electro-thermal
Materials	Indium tin oxide	$\text{Ge}_3\text{Sb}_2\text{Te}_6$	Si	VO_2
Wavelength	Visible	Visible / NIR	Telecom	NIR
Tunability	78 %	-	-	33 %
Speed	0.02 s	500 ms	-	1.27 ms
Reference	K. Thyagarajan et al., <i>Adv. Mater.</i> 1701044 (2017).	Q. Wang et al., <i>Nat. Photon.</i> 10(1), 60 (2016).	J. Sun et al., <i>Nature</i> 493(7431), 195 (2013).	Z. Zhu et al., <i>Nano Lett.</i> 17(8), 4881 (2017).

I. Introduction

II. Metasurface holograms

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IV. Conclusion

- **Metasurface hologram** shows a good potential to improve conventional holograms.
- Using X-shaped meta-atoms, we first demonstrated **broadband complex-amplitude holograms** with silicon base in the visible range.
- Metasurface lens (**metalens**) has attracted much interest to miniaturize optical applications including augmented reality (AR) and virtual reality (VR) devices.
- **Aberration-controlled metalens** and active modulation of light with **tunable metasurfaces** are important topics of study.



THANK YOU

ByoungHo Lee

byoungho@snu.ac.kr

*School of Electrical and Computer Engineering
Seoul National University, Republic of Korea*